

Constructing a Rosetta Stone for Plasma Heating and Particle Acceleration in Kinetic Plasma Physics

Gregory G. Howes
University of Iowa

Princeton Plasma Physics Laboratory Heliophysics Seminar
Virtually anywhere
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Collaborators

Andrew J. McCubbin
Sarah A. Horvath
Collin Brown
Arya Afshari
Peter Montag
James Juno
Jennifer L. Verniero
James W. R. Schroeder
Kristopher G. Klein
Jason M. TenBarge
Tak Chu Li

U of Iowa
Princeton/Iowa
UC Berkeley
Wheaton Coll
U of Arizona
Princeton U
Dartmouth Coll

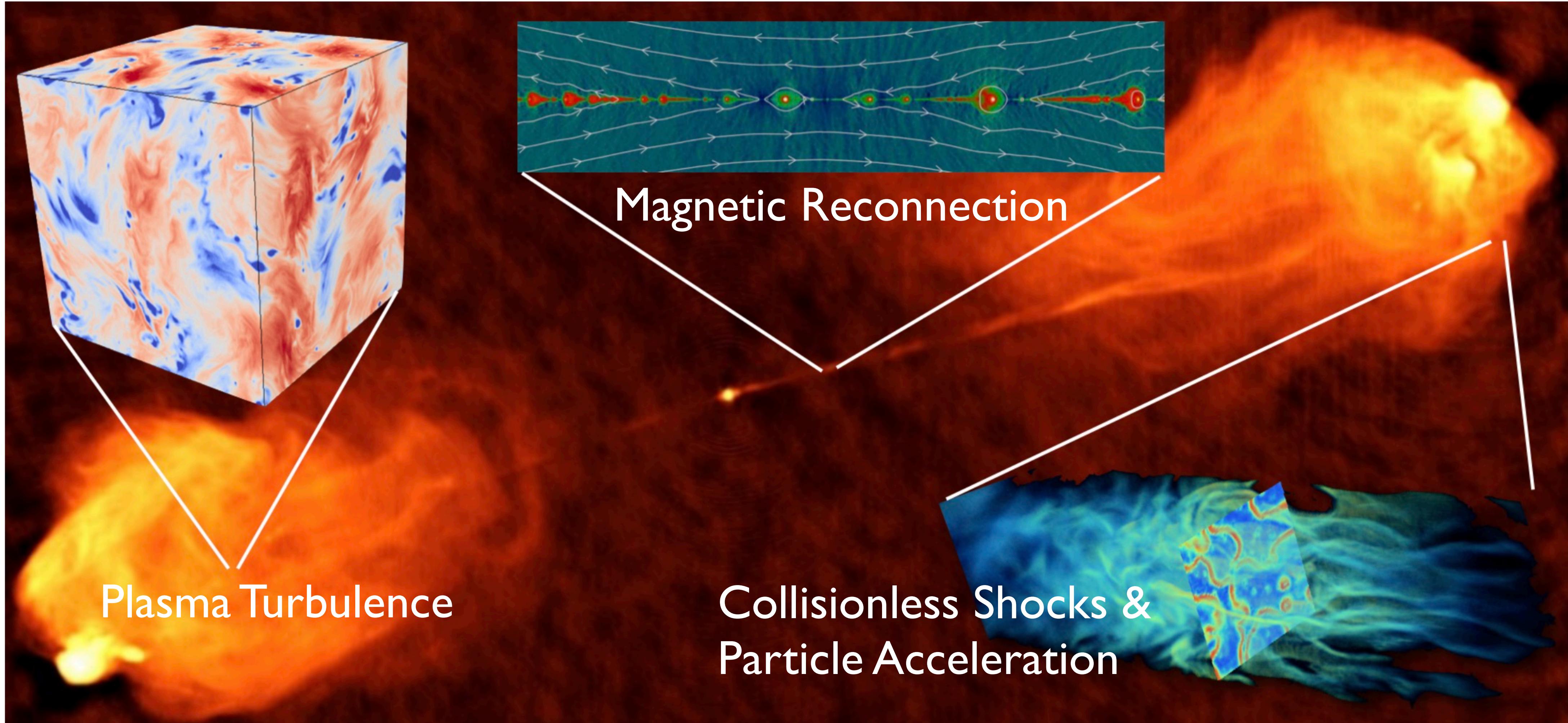
Christopher H. K. Chen
Francesco Valentini
Steven Baek
Jasper Halekas
Damiano Caprioli
Anatoly Spitkovsky
Lynn Wilson III
Ryusuke Numata
Matt Kunz
Lev Arzamasskiy

Queen Mary U
U of Calabria
U of Iowa
U of Iowa
U Chicago
Princeton U
NASA Goddard
U of Hyogo, Japan
Princeton U
Princeton U

Outline

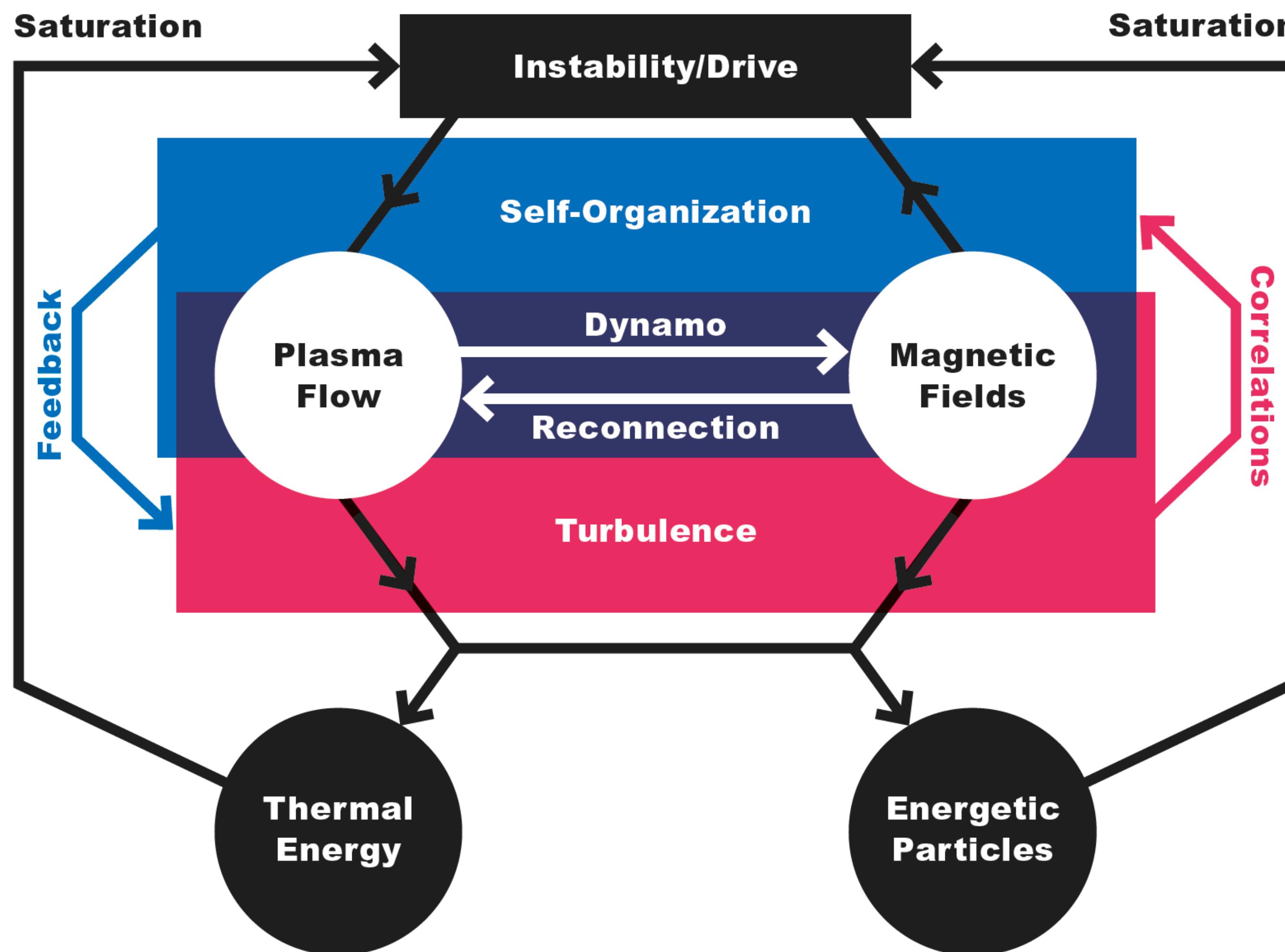
- Plasma Heating and Particle Acceleration in the Heliosphere
- Kinetic Theory of Particle Energization
 - Field-Particle Correlation Technique
- Three Applications of the Field-Particle Correlation Technique
 - Plasma Heating by Dissipation of Plasma Turbulence
 - Ion Energization in Collisionless Shocks
 - Electron Fermi Acceleration in Collisionless Magnetic Reconnection
- Constructing a “Rosetta Stone” for particle energization
- Conclusions

Fundamental Plasma Physics Processes



(Connecting Micro and Macro Scales: Acceleration, Reconnection, and Dissipation in Astrophysical Plasmas,
Kavli Institute for Theoretical Physics, UCSB 2019)

Energetics of the Plasma Universe



“Understanding the Energetics of the Plasma Universe”,
Plasma: At the Frontier of Scientific Discovery

DOE Report 2015

Understanding the **flow of energy** in astrophysical plasmas is a key overarching theme:

- Turbulence
- Instabilities
- Magnetic Reconnection
- Collisionless Shocks
- Particle Acceleration
- Magnetic Dynamo

Under the weakly collisional conditions of most space and astrophysical plasmas, **kinetic theory is essential to understand these processes**.

Developing a Predictive Capability

We need to understand which **processes** are responsible for particle energization

Example: What heats the Solar Corona to $T > 10^6$ K ?

- Magnetic Reconnection (nanoflares, loop opening, etc.) (Parker, 1991; Fisk et al. 1999)
- Dissipation of Plasma Turbulence (Hollweg, 1986; Cranmer et al. 2003, 2005, 2007)
- Other mechanisms? (velocity filtration, etc.) (Scudder, 1992)

But . . . even if we know the general process, what is the **specific mechanism?**

Example: Turbulence

- What processes govern the dissipation of turbulence and resulting plasma heating?

Proposed Dissipation Mechanisms in Turbulence



Weakly Collisional Plasma Turbulence

(1) **Resonant Wave-Particle Interactions** (Landau damping, transit-time damping, cyclotron damping)

(Barnes 1966; Coleman 1968; Denskat *et al.*, 1983; Isenberg & Hollweg 1983; Goldstein *et al.* 1994; Quataert 1998; Leamon *et al.*, 1998, 1999, 2000; Gary 1999; Quataert & Gruzinov, 1999; Isenberg *et al.* 2001; Hollweg & Isenberg 2002; Howes *et al.* 2008; Schekochihin *et al.* 2009; TenBarge & Howes 2013; Howes 2015; Li, Howes, Klein, & TenBarge 2016)

(2) **Nonresonant Wave-Particle Interactions** (stochastic ion heating, magnetic pumping)

(Berger, 1958; Johnson & Cheng, 2001; Chen *et al.* 2001; White *et al.*, 2002; Voitenko & Goosens, 2004; Bourouaine *et al.*, 2008; Chandran *et al.* 2010; Chandran 2010, Chandran *et al.* 2011; Bourouaine & Chandran 2013,Lichko:2017)

(3) **Dissipation in Current Sheets** (collisionless magnetic reconnection)

(Dmitruk *et al.* 2004; Markovskii & Vasquez 2011; Matthaeus & Velli 2011; Osman *et al.* 2011; Servidio 2011; Osman *et al.* 2012a,b; Wan *et al.* 2012; Karimabadi *et al.* 2013; Zhdankin *et al.* 2013; Osman *et al.* 2014a,b; Zhdankin *et al.* 2015a,b; Loureiro & Boldyrev, 2017a,b; Boldyrev & Loureiro 2017; Walker, Boldyrev, & Loureiro 2018)

We can use information about the flow of energy in **velocity space**
to distinguish between different mechanisms.

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Maxwell-Boltzmann Equations of Kinetic Plasma Theory

$$\frac{\partial f_s}{\partial t} + \mathbf{v} \cdot \nabla f_s + \frac{q_s}{m_s} \left[\mathbf{E} + \frac{\mathbf{v} \times \mathbf{B}}{c} \right] \cdot \frac{\partial \mathbf{f}_s}{\partial \mathbf{v}} = \left(\frac{\partial f_s}{\partial t} \right)_{\text{coll}}^0$$

Lorentz Term responsible for interactions between fields and particles

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \cdot \mathbf{E} = 4\pi \rho_q$$

$$\nabla \times \mathbf{B} = \frac{4\pi}{c} \mathbf{j} + \frac{1}{c} \frac{\partial \mathbf{E}}{\partial t}$$

$$\nabla \cdot \mathbf{B} = 0$$

Distribution Function:
3D-3V phase space

$$f_s(\mathbf{r}, \mathbf{v}, t)$$

Electromagnetic
Fields

$$\mathbf{E}(\mathbf{r}, t), \mathbf{B}(\mathbf{r}, t)$$

Particle Energization

Conserved Vlasov-Maxwell Energy

$$W = \int d^3\mathbf{r} \frac{|\mathbf{E}|^2 + |\mathbf{B}|^2}{8\pi} + \sum_s \int d^3\mathbf{r} \int d^3\mathbf{v} \frac{1}{2} m_s v^2 f_s$$

EM Field Energy

Particle Energy

$$W_s = \int d^3\mathbf{r} \int d^3\mathbf{v} \frac{1}{2} m_s v^2 f_s$$

We want to measure the change in particle energy ...

... using measurements of the change in the distribution function.

$$\frac{\partial W_s}{\partial t} = \int d^3\mathbf{r} \int d^3\mathbf{v} \frac{1}{2} m_s v^2 \frac{\partial f_s}{\partial t}$$

Vlasov Equation

$$\frac{\partial f_s}{\partial t} + \mathbf{v} \cdot \nabla f_s + \frac{q_s}{m_s} \left[\mathbf{E} + \frac{\mathbf{v} \times \mathbf{B}}{c} \right] \cdot \frac{\partial f_s}{\partial \mathbf{v}} = 0$$

Particle Energization

$$\frac{\partial W_s}{\partial t} = \int d^3\mathbf{r} \int d^3\mathbf{v} \frac{1}{2} m_s v^2 \left[-\mathbf{v} \cdot \nabla f_s - \frac{q_s}{m_s} \mathbf{E} \cdot \frac{\partial f_s}{\partial \mathbf{v}} - \frac{q_s}{m_s} \left(\frac{\mathbf{v}}{c} \times \mathbf{B} \right) \cdot \frac{\partial f_s}{\partial \mathbf{v}} \right]$$

Perfect Differential

Electric field does work on particles

Rate of Change of Particle Energy

$$\boxed{\frac{\partial W_s}{\partial t} = - \int d^3\mathbf{r} \int d^3\mathbf{v} \left[q_s \frac{v^2}{2} \mathbf{E}(\mathbf{r}, t) \cdot \frac{\partial f_s(\mathbf{r}, \mathbf{v}, t)}{\partial \mathbf{v}} \right]}$$

But this is integrated over velocity and space...

Not observationally accessible!

How Does Energy Flow in Phase Space?

Maximize use of 3V phase-space information to characterize energization

Define: **Phase-space energy density** $w_s(\mathbf{r}, \mathbf{v}, t) = \frac{m_s v^2}{2} f_s(\mathbf{r}, \mathbf{v}, t)$

$$\frac{\partial w_s(\mathbf{r}, \mathbf{v}, t)}{\partial t} = -\mathbf{v} \cdot \nabla w_s - q_s \frac{v^2}{2} \mathbf{E} \cdot \frac{\partial f_s}{\partial \mathbf{v}} - \frac{q_s v^2}{c} (\mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f_s}{\partial \mathbf{v}}$$

Advection of
particle energy

Pressure forces
in fluid theory

Work done by
electric field

Responsible for
net change in
particle energy

Work done by
magnetic field

Integrates over
velocity space
to zero

How Does Energy Flow in Phase Space?

$$\frac{\partial w_s(\mathbf{r}, \mathbf{v}, t)}{\partial t} = -\mathbf{v} \cdot \nabla w_s - \boxed{q_s \frac{v^2}{2} \mathbf{E} \cdot \frac{\partial f_s}{\partial \mathbf{v}}} - \frac{q_s}{c} \frac{v^2}{2} (\mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f_s}{\partial \mathbf{v}}$$

Note:

- 1) Advective (pressure) and magnetic terms can convert particle energy
e.g., Bulk flow kinetic energy to thermal kinetic energy
- 2) Electric field term is the only one that changes the net particle energy

Therefore we focus here on the work on particles done by $\mathbf{E}(\mathbf{r}, t)$

Other applications may require analyzing the other terms!

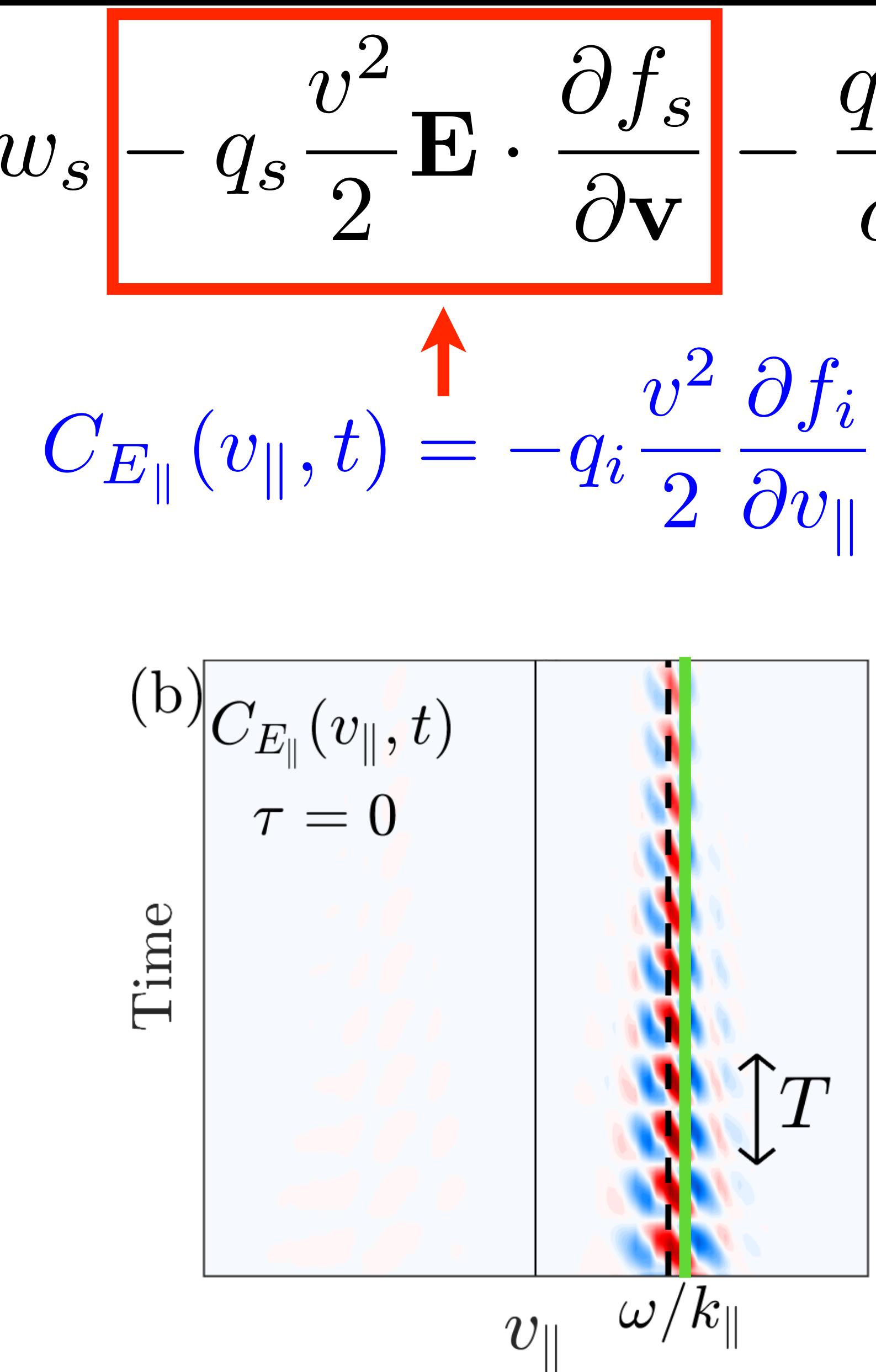
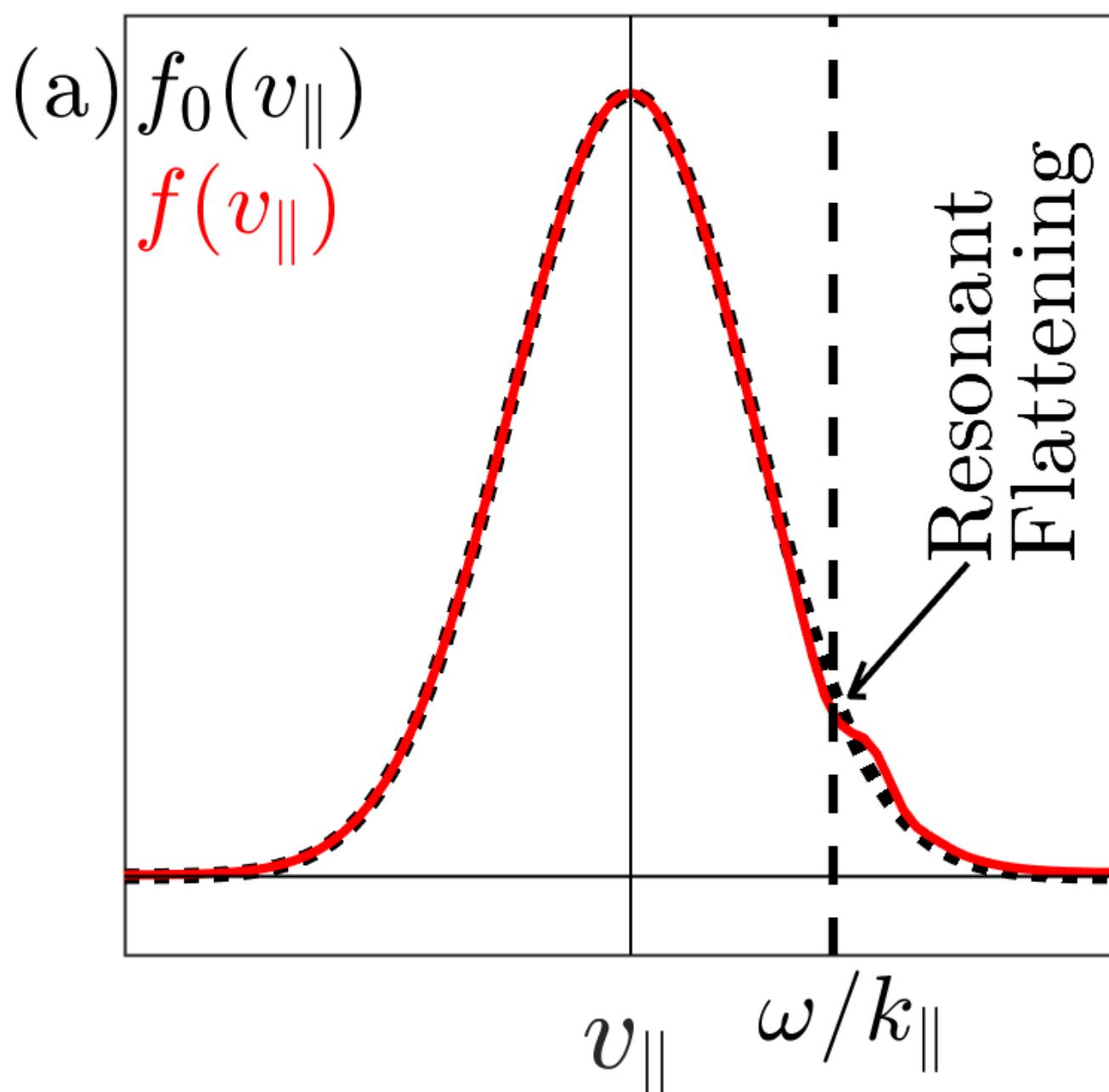
How Do We Exploit Velocity-Space Information?

$$\frac{\partial w_s(\mathbf{r}, \mathbf{v}, t)}{\partial t} = -\mathbf{v} \cdot \nabla w_s - q_s \frac{v^2}{2} \mathbf{E} \cdot \frac{\partial f_s}{\partial \mathbf{v}} - \frac{q_s v^2}{c} \frac{2}{2} (\mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f_s}{\partial \mathbf{v}}$$

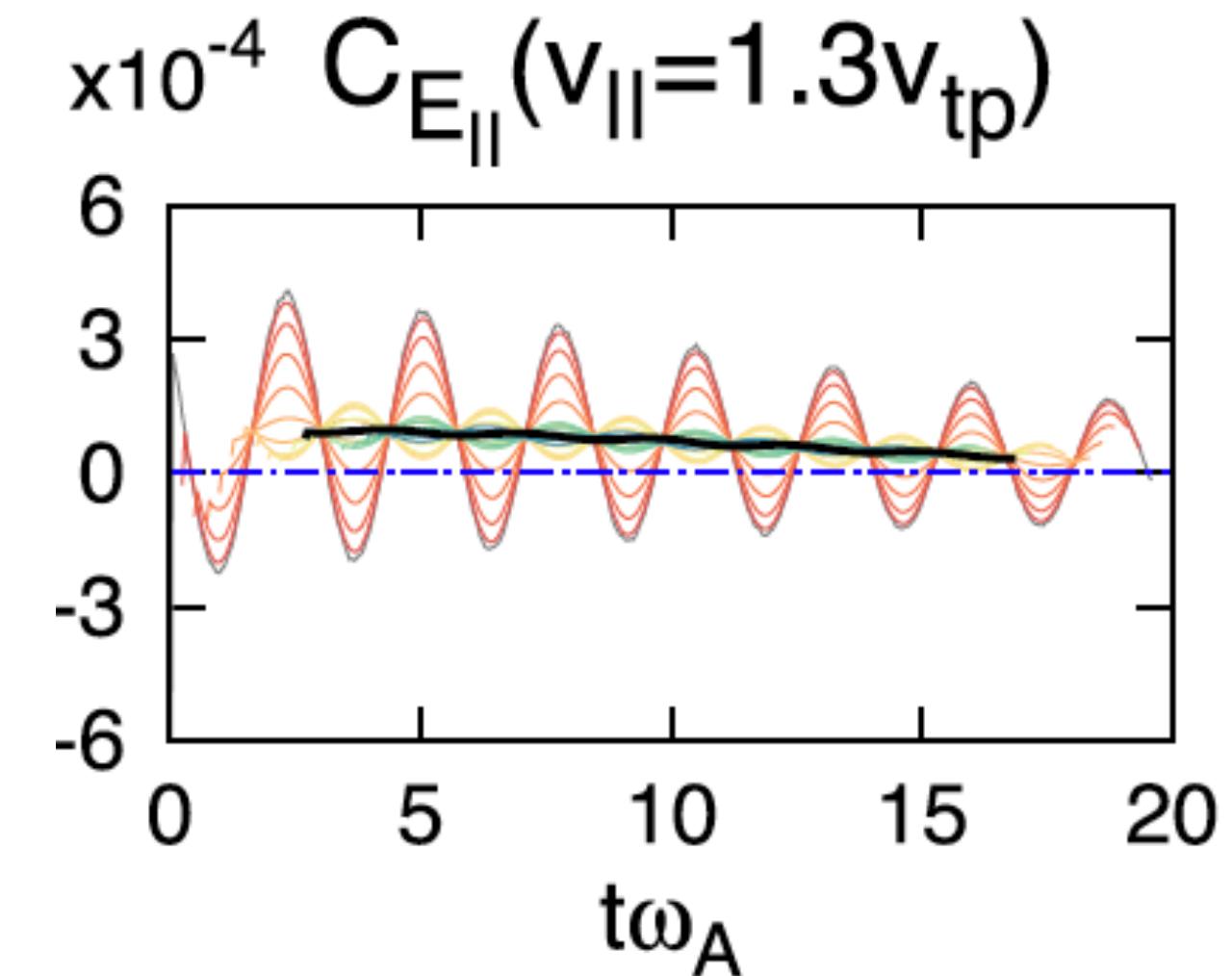
Two parts of signal:

Example:

Ion Landau Damping of a single Kinetic Alfvén Wave

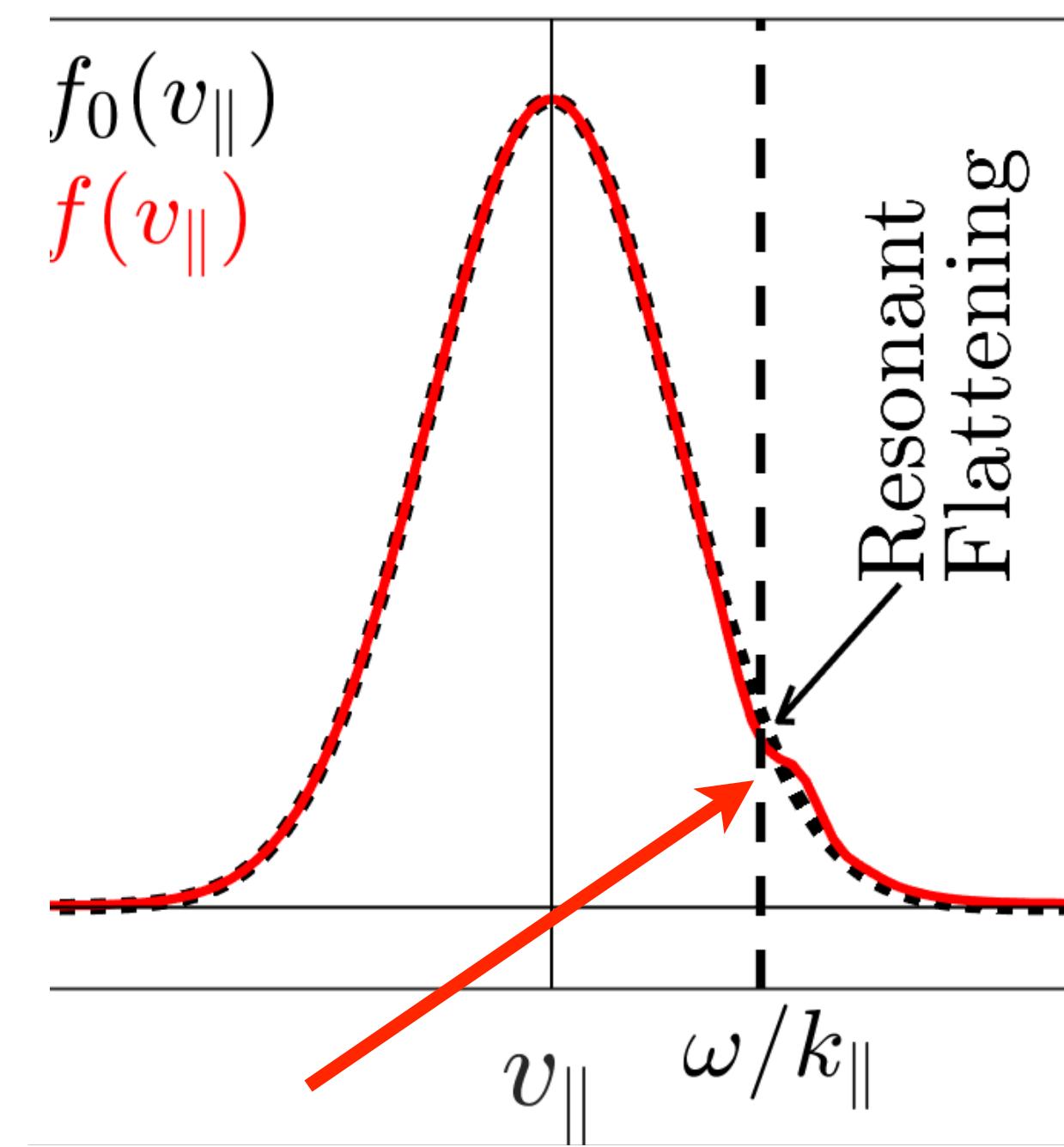
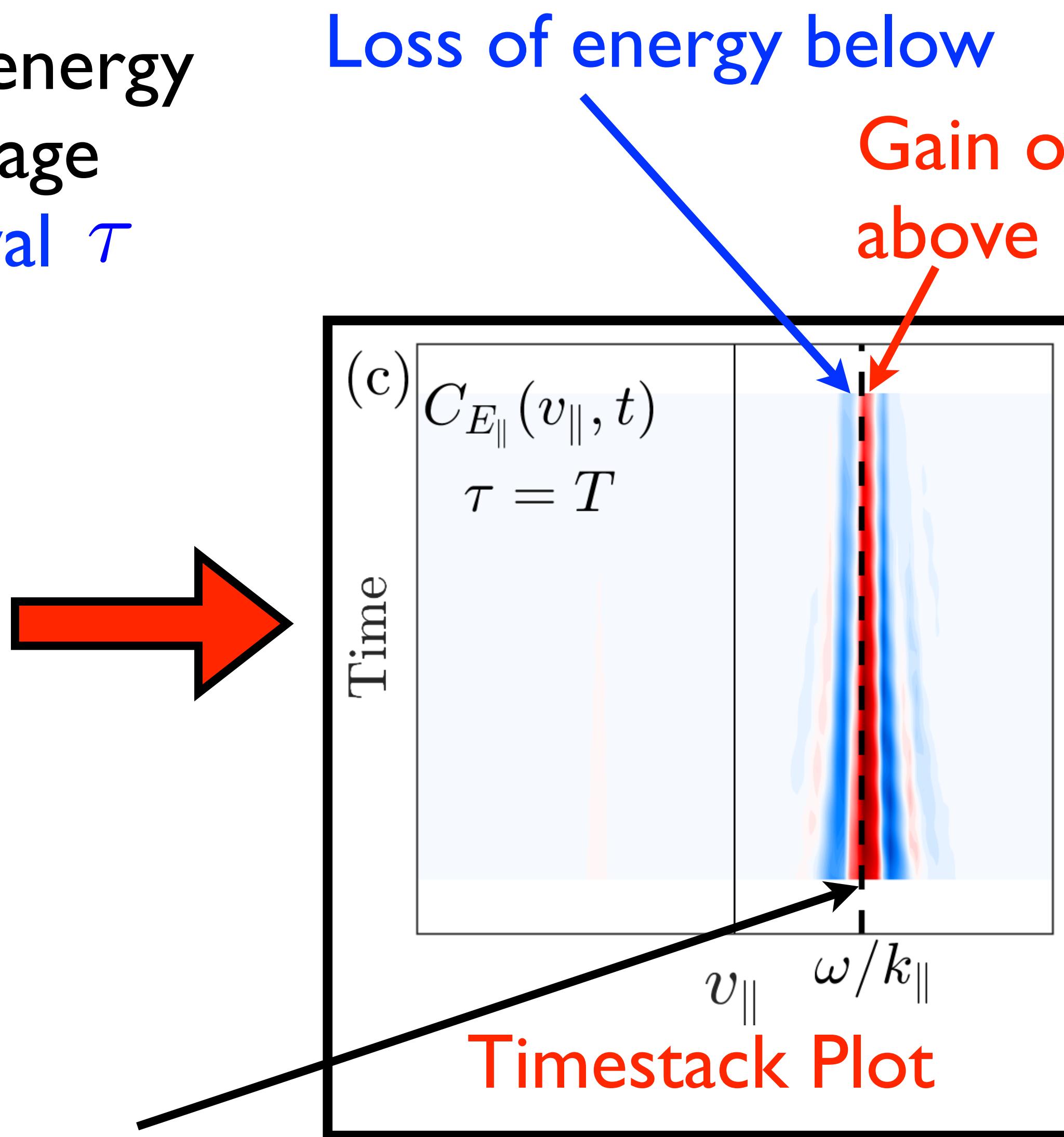
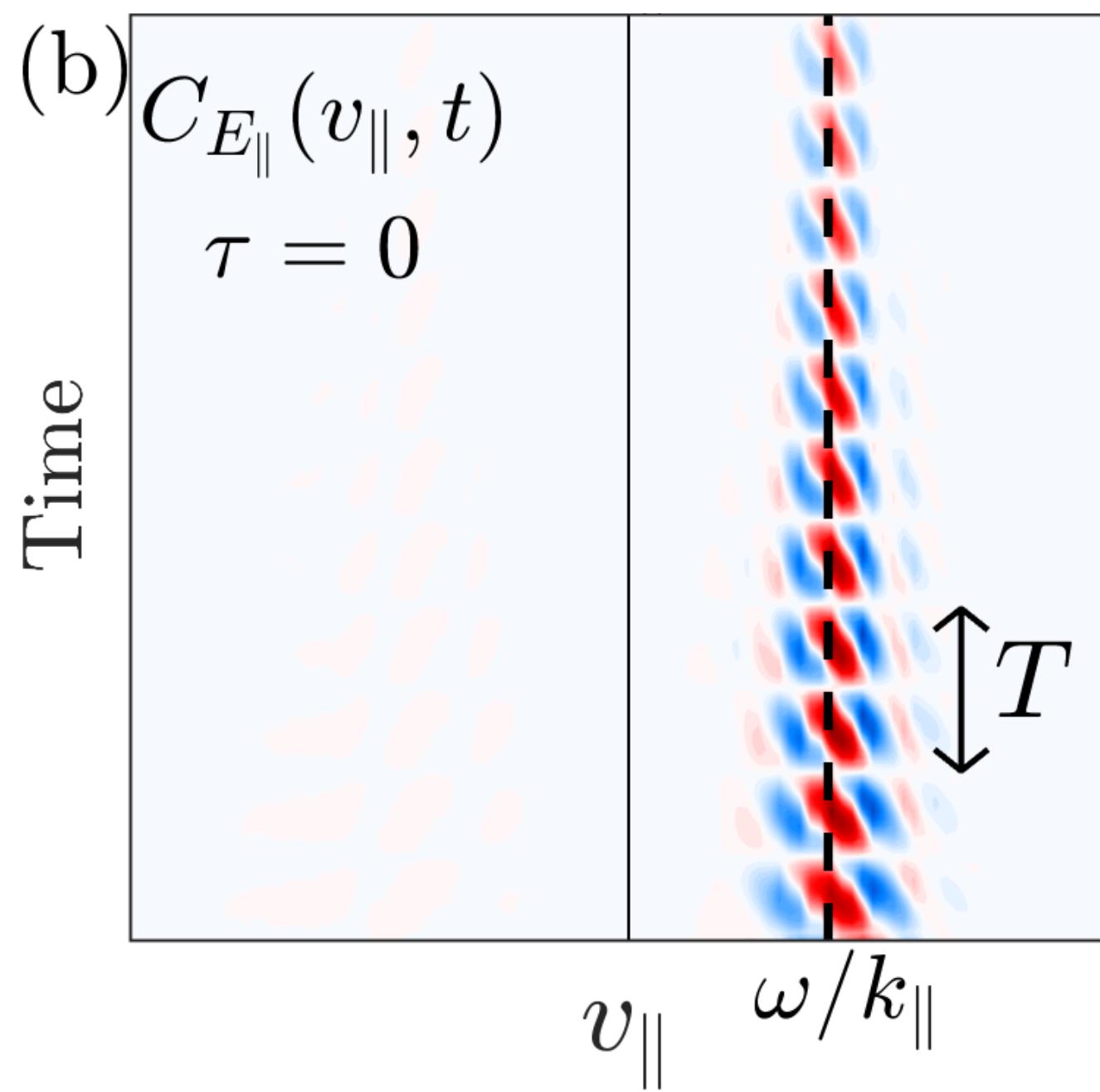


- 1) Conservative oscillatory energy transfer of undamped wave motion
- 2) Secular energy transfer of collisionless damping



How Do We Exploit Velocity-Space Information?

To eliminate oscillatory energy transfer, take a time average over a **correlation interval** τ

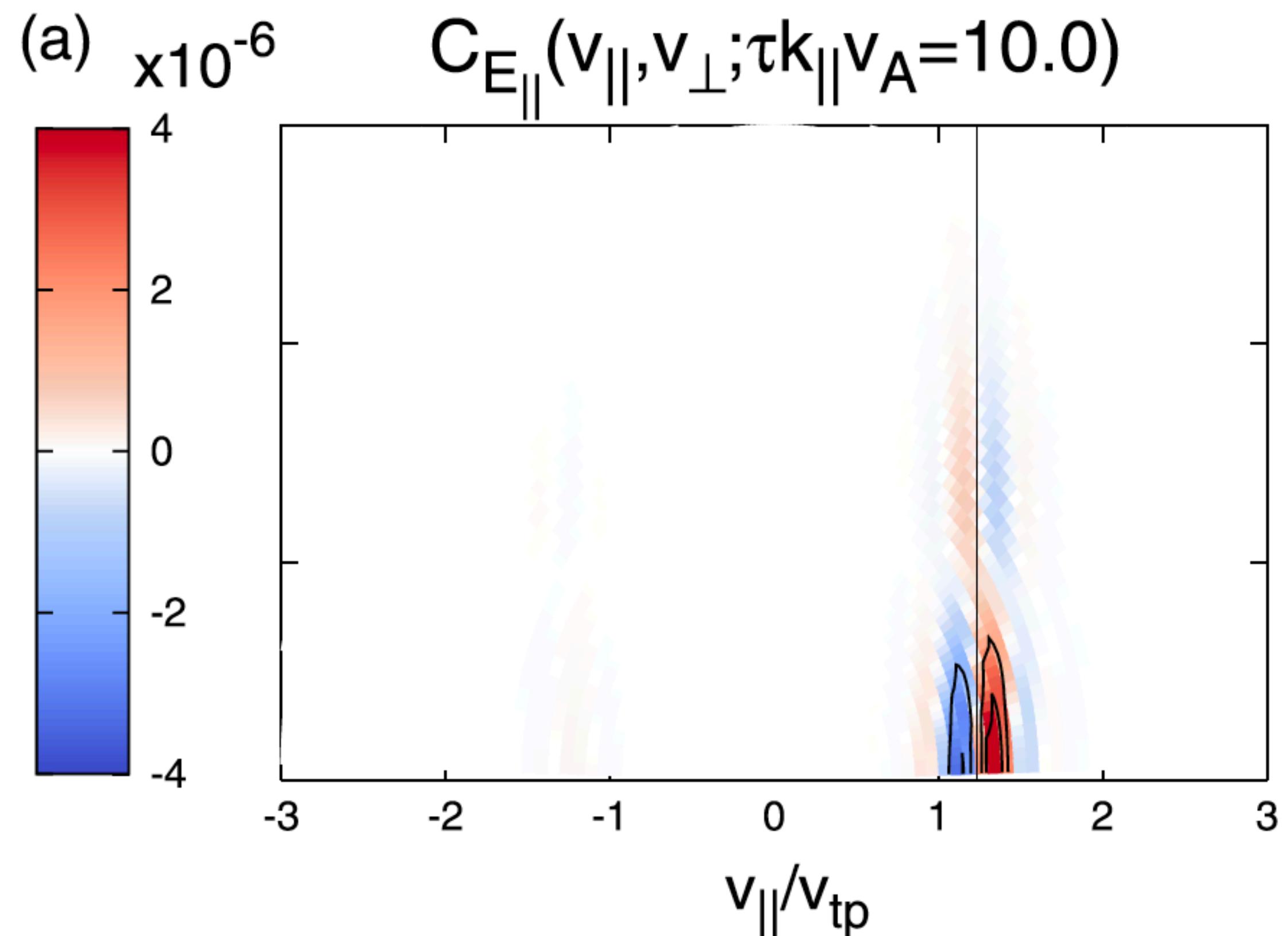


$$w_i = \frac{1}{2}mv^2 f_i(v\parallel, t)$$

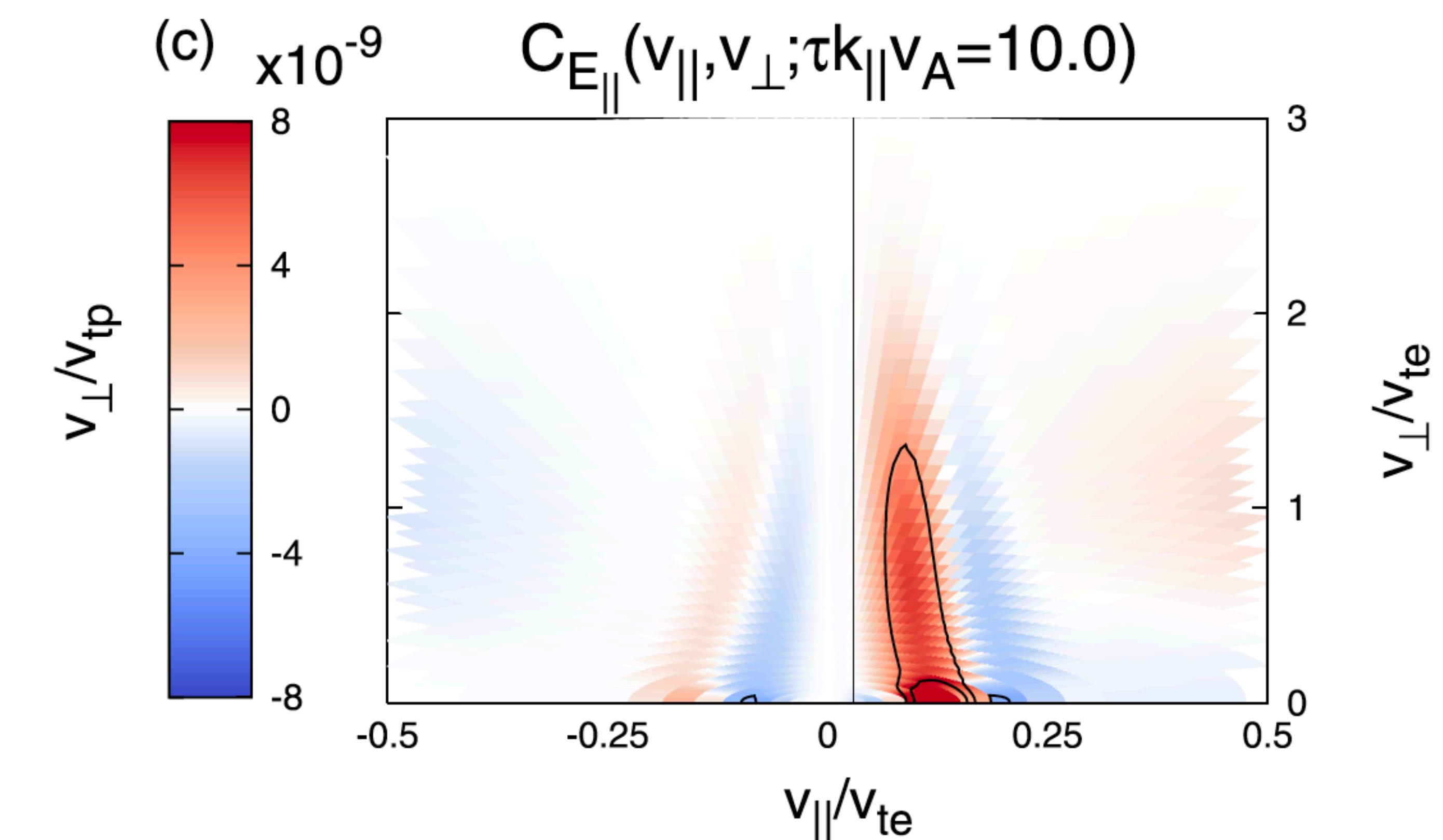
Gyrotropic Velocity-Space Signatures

Velocity-space signatures in gyrotropic velocity space (v_{\parallel}, v_{\perp}) (Howes, PoP, 2017)

Single KAW with: $\beta_i = 1$ $T_i/T_e = 1$ $k_{\perp}\rho_i = 1.3$



Velocity-space signature
of ion Landau damping



Velocity-space signature
of electron Landau damping

Field-Particle Correlation Technique

Phase-space energy density $w_s(\mathbf{r}, \mathbf{v}, t) = \frac{m_s v^2}{2} f_s(\mathbf{r}, \mathbf{v}, t)$

$$\frac{\partial w_s(\mathbf{r}, \mathbf{v}, t)}{\partial t} = -\mathbf{v} \cdot \nabla w_s - q_s \frac{v^2}{2} \mathbf{E} \cdot \frac{\partial f_s}{\partial \mathbf{v}} - \frac{q_s}{c} \frac{v^2}{2} (\mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f_s}{\partial \mathbf{v}}$$

At a single-point \mathbf{r}_0 ,
compute correlation of field \mathbf{E}
and particle $f_s(\mathbf{v})$ measurements
over correlation interval τ

Field-particle correlation

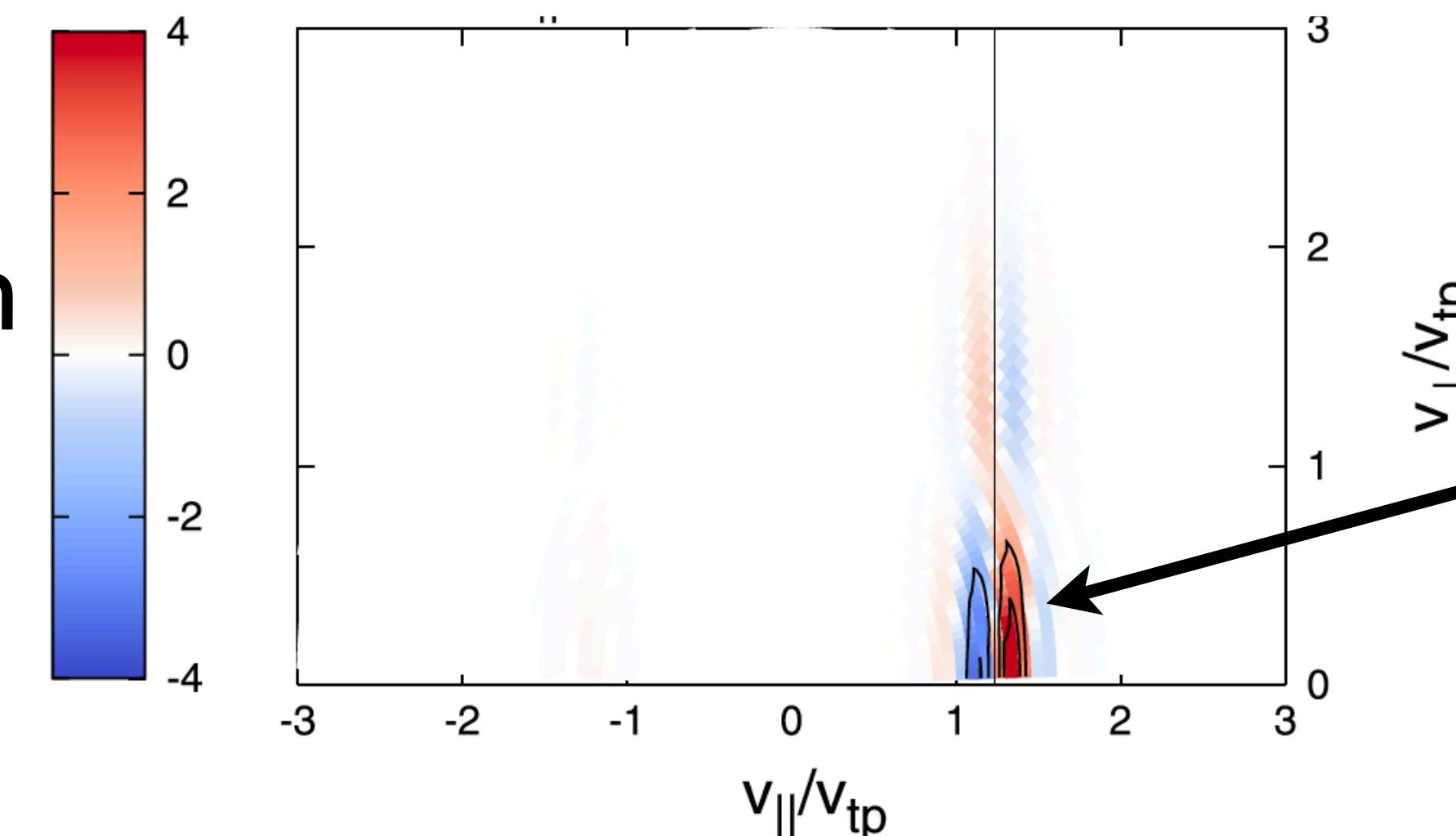
$$C_{E_{\parallel}}(\mathbf{v}, t, \tau) = C \left(-q_s \frac{v_{\parallel}^2}{2} \frac{\partial f_s(\mathbf{r}_0, \mathbf{v}, t)}{\partial v_{\parallel}}, E_{\parallel}(\mathbf{r}_0, t) \right)$$

(Klein & Howes, 2016; Howes, Klein, & Li, 2017;
Klein, Howes, & TenBarge, 2017)

Depending on the problem, one can alternatively, can compute correlations
with E_{\perp} or with (E_x, E_y, E_z)

Scientific Insight from Field-Particle Correlations

I) Distinguish and identify kinetic energization mechanisms through unique **velocity-space signature**



Bipolar signature
of Landau damping

2) Determine rate of change of spatial energy density

$$W_s(\mathbf{r}, t) \equiv \int d^3\mathbf{v} \frac{1}{2} m_s v^2 f_s(\mathbf{r}, \mathbf{v}, t)$$
$$\frac{\partial W_s}{\partial t} = \int d^3\mathbf{v} C_{E_{\parallel}}(\mathbf{v}, t) = - \int d^3\mathbf{v} q_s \frac{v_{\parallel}^2}{2} \frac{\partial f_s}{\partial v_{\parallel}} E_{\parallel} = \int d^3\mathbf{v} q_s v_{\parallel} E_{\parallel} f_s = j_{\parallel, s} E_{\parallel}$$

Integrate correlation
over velocity space

Rate of work done on species s
by E_{\parallel}

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Strong Plasma Turbulence

Driven nonlinear gyrokinetic simulation of solar wind turbulence

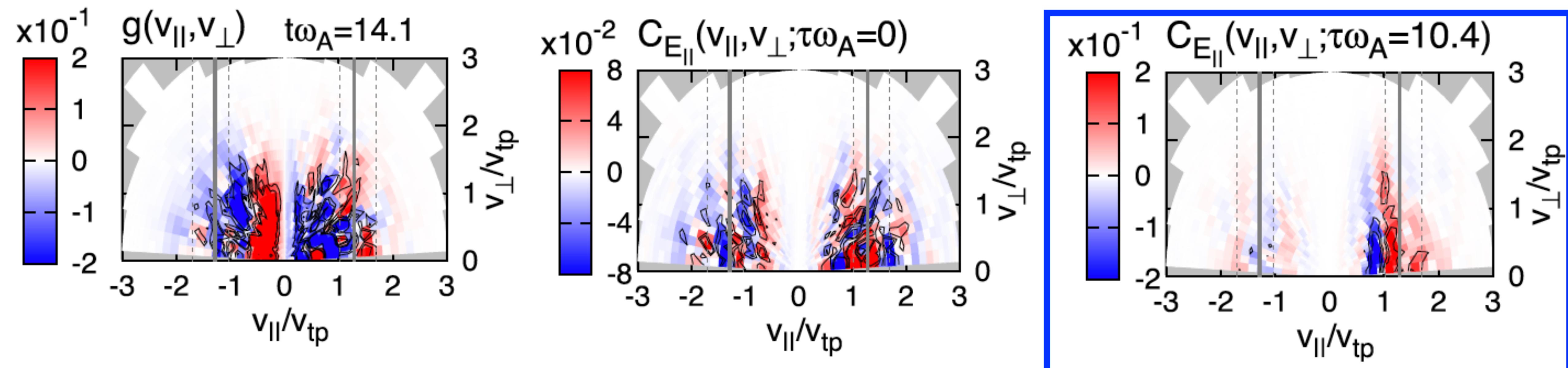
(Klein, Howes, & Tenbarge, 2017)

Plasma parameters: $\beta_i = 1$

$T_i/T_e = 1$

Turbulence parameters: $k_\perp \gg k_\parallel$ (Anisotropic)

$\chi \sim 1$ (Critical balance)



Evidence of Landau Damping
in strong plasma turbulence

Proposed Dissipation Mechanisms in Turbulence



Weakly Collisional Kinetic Turbulence

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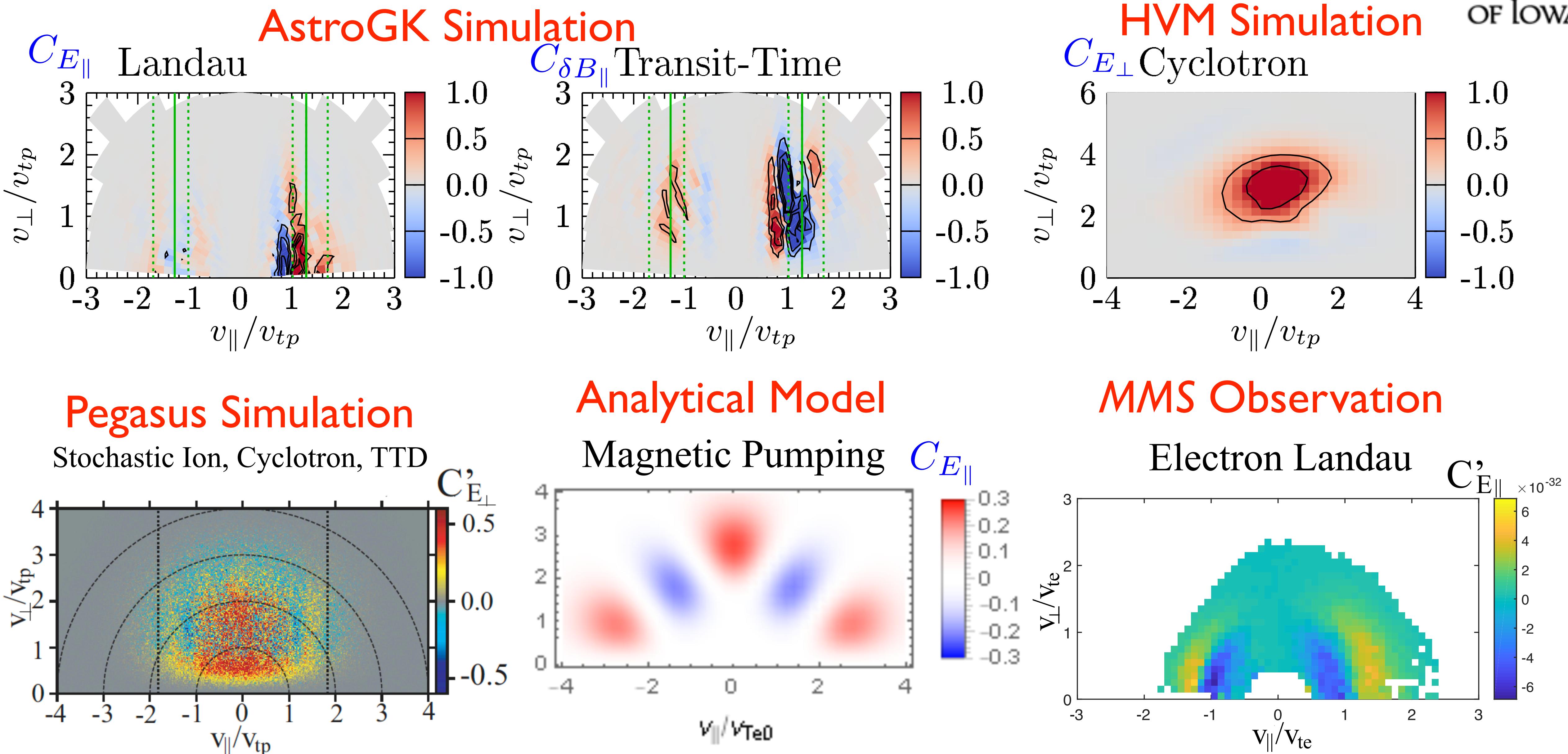
(Berger, 1958; Johnson & Cheng, 2001; Chen *et al.* 2001; White *et al.*, 2002; Voitenko & Goosens, 2004; Bourouaine *et al.*, 2008; Chandran *et al.* 2010; Chandran 2010, Chandran *et al.* 2011; Bourouaine & Chandran 2013,Lichko:2017)

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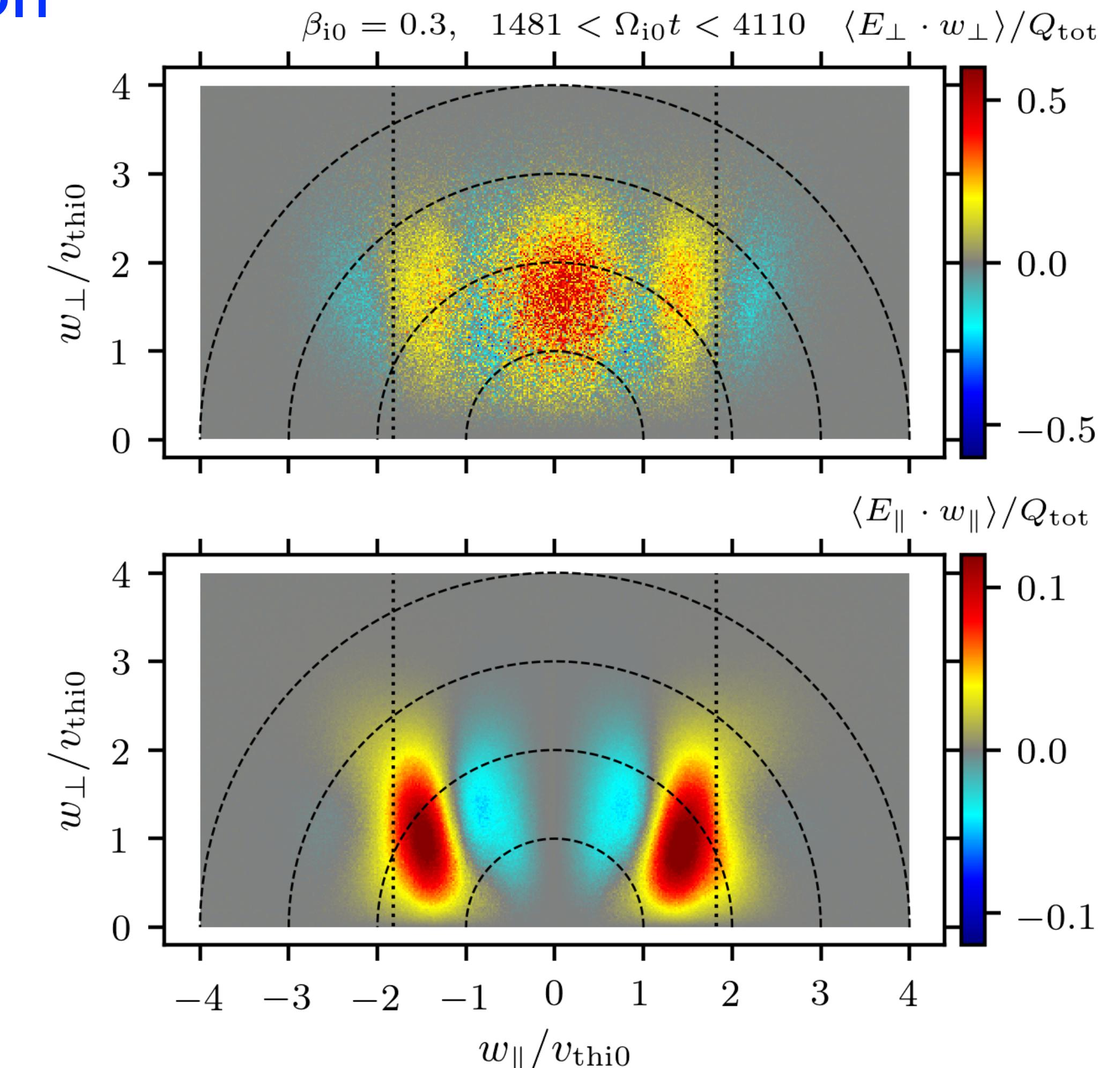
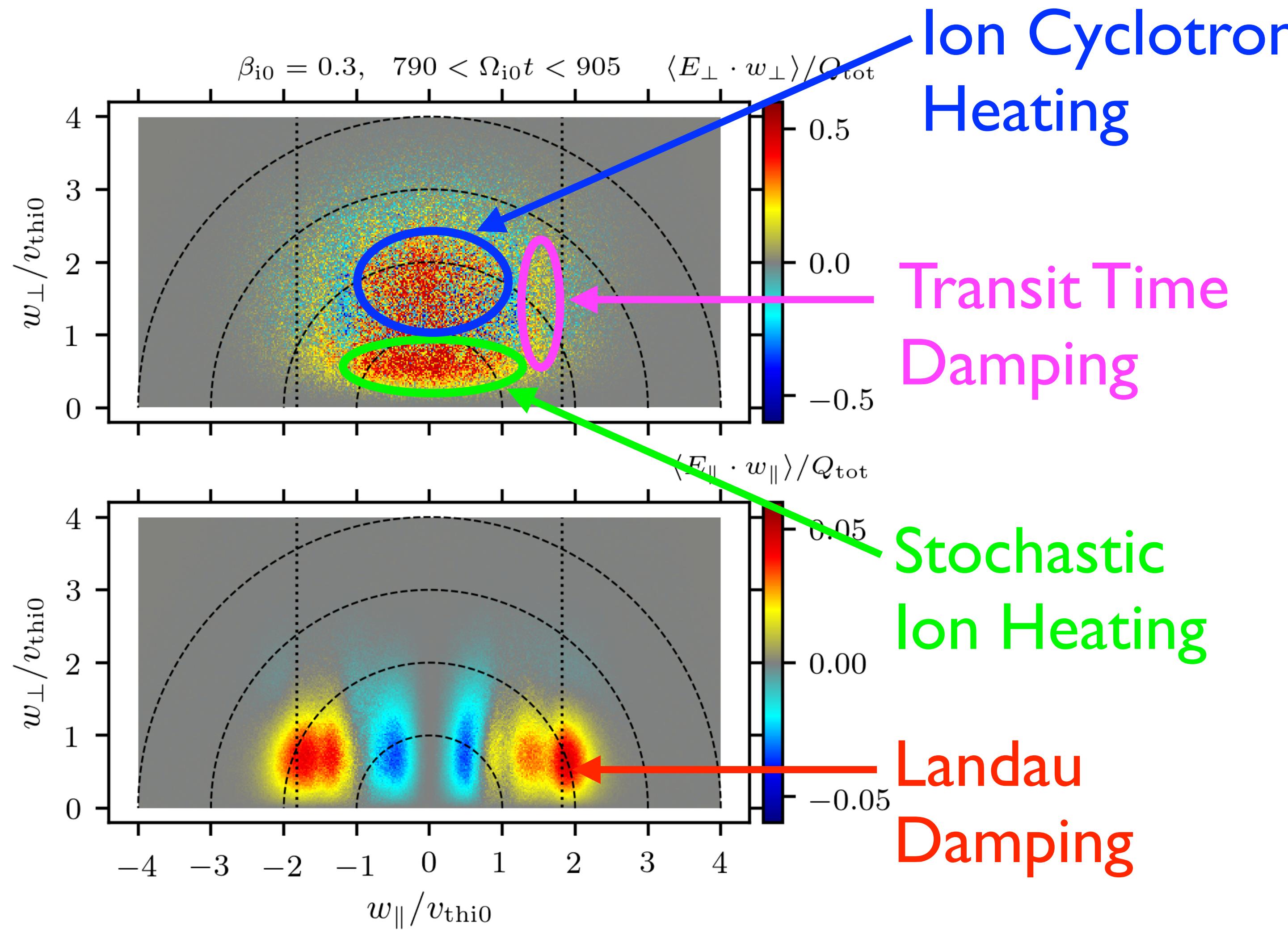
The **velocity-space signature** obtained from **field-particle correlations** has the potential to distinguish between different mechanisms.

Distinguishing Energization Mechanisms



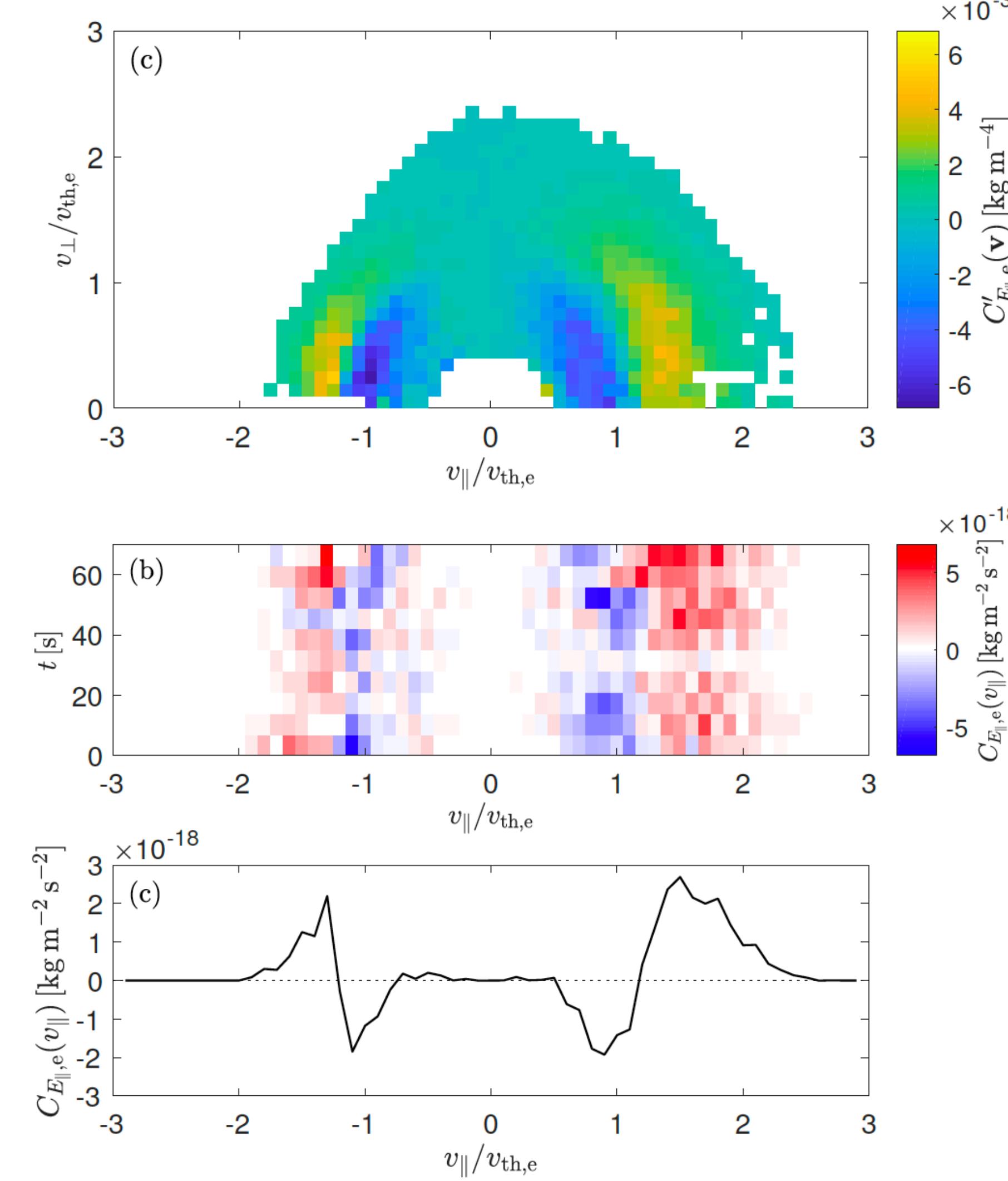
Identifying Mechanisms by Velocity-Space Signatures

Pegasus hybrid simulations of anisotropic turbulence

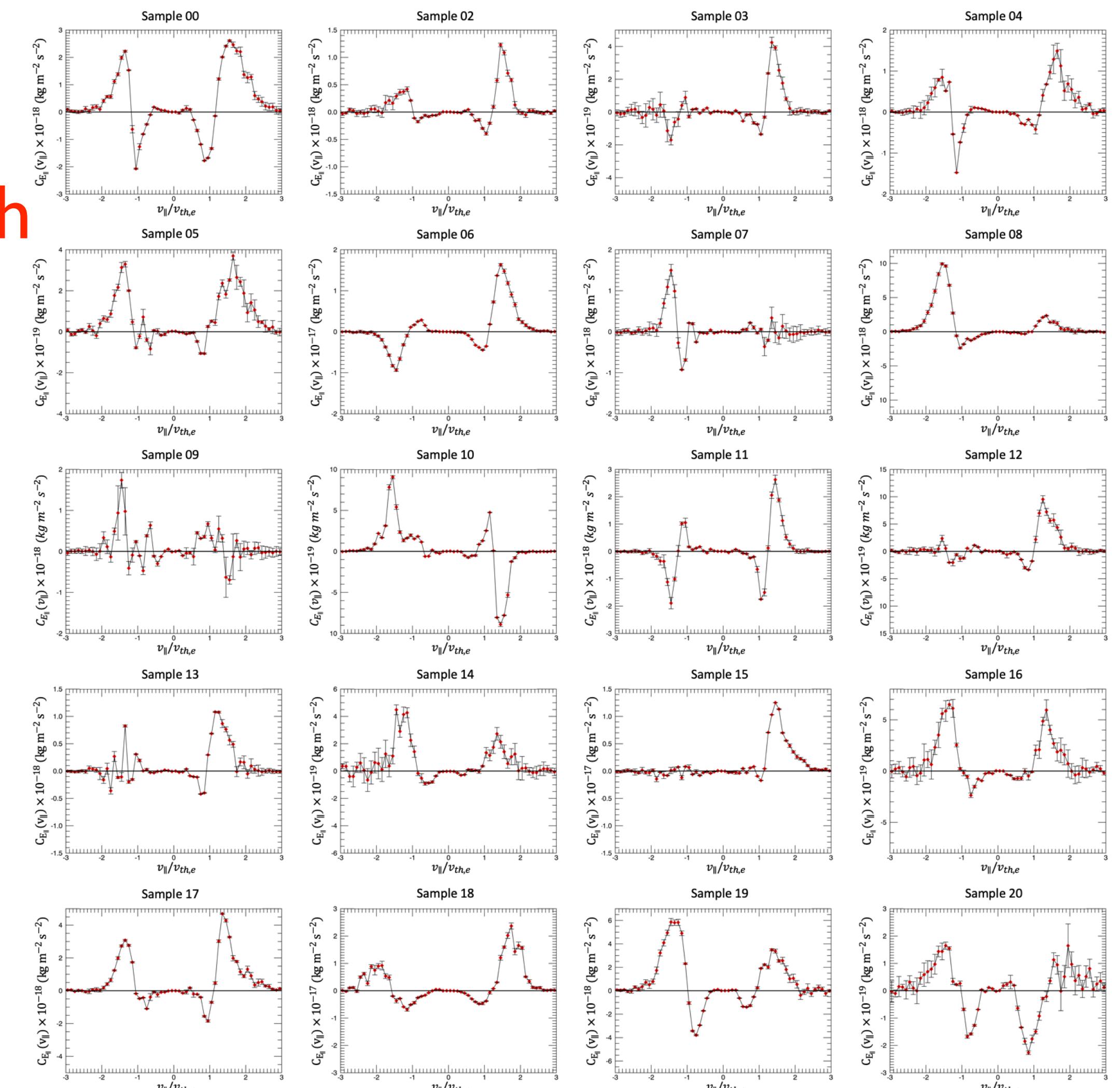


Determining Rate of Particle Energization

Observations of electron energization in Earth's magnetosheath using MMS

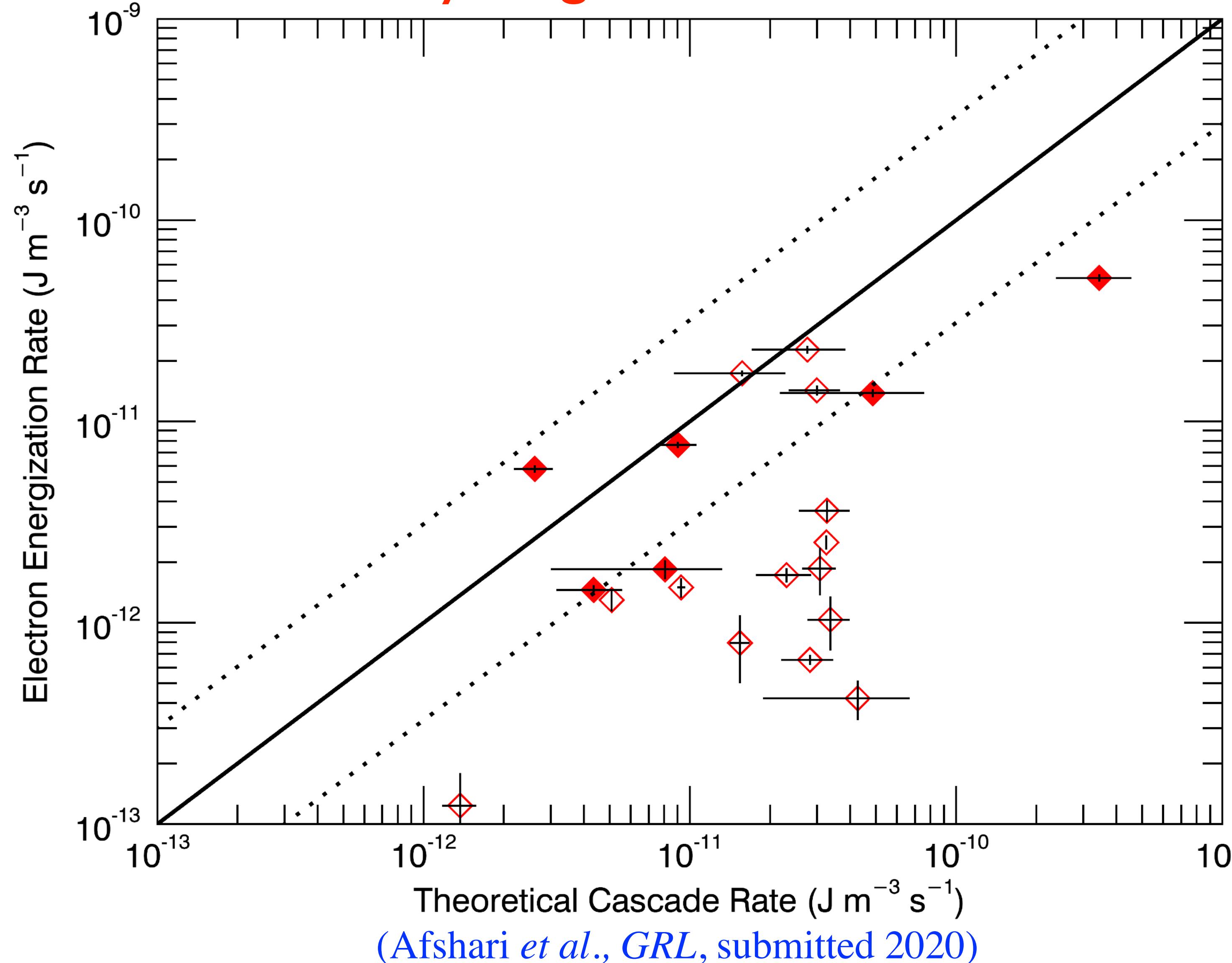


Twenty
magnetosheath
intervals



Determining Rate of Particle Energization

Twenty magnetosheath intervals



Two Major Findings:

- 19 of 20 intervals show velocity-space signatures of **electron Landau damping**
- Nearly half of intervals have **electron Landau damping rates comparable to the estimated turbulent cascade rate**

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Simulation of Perpendicular Collisionless Shock

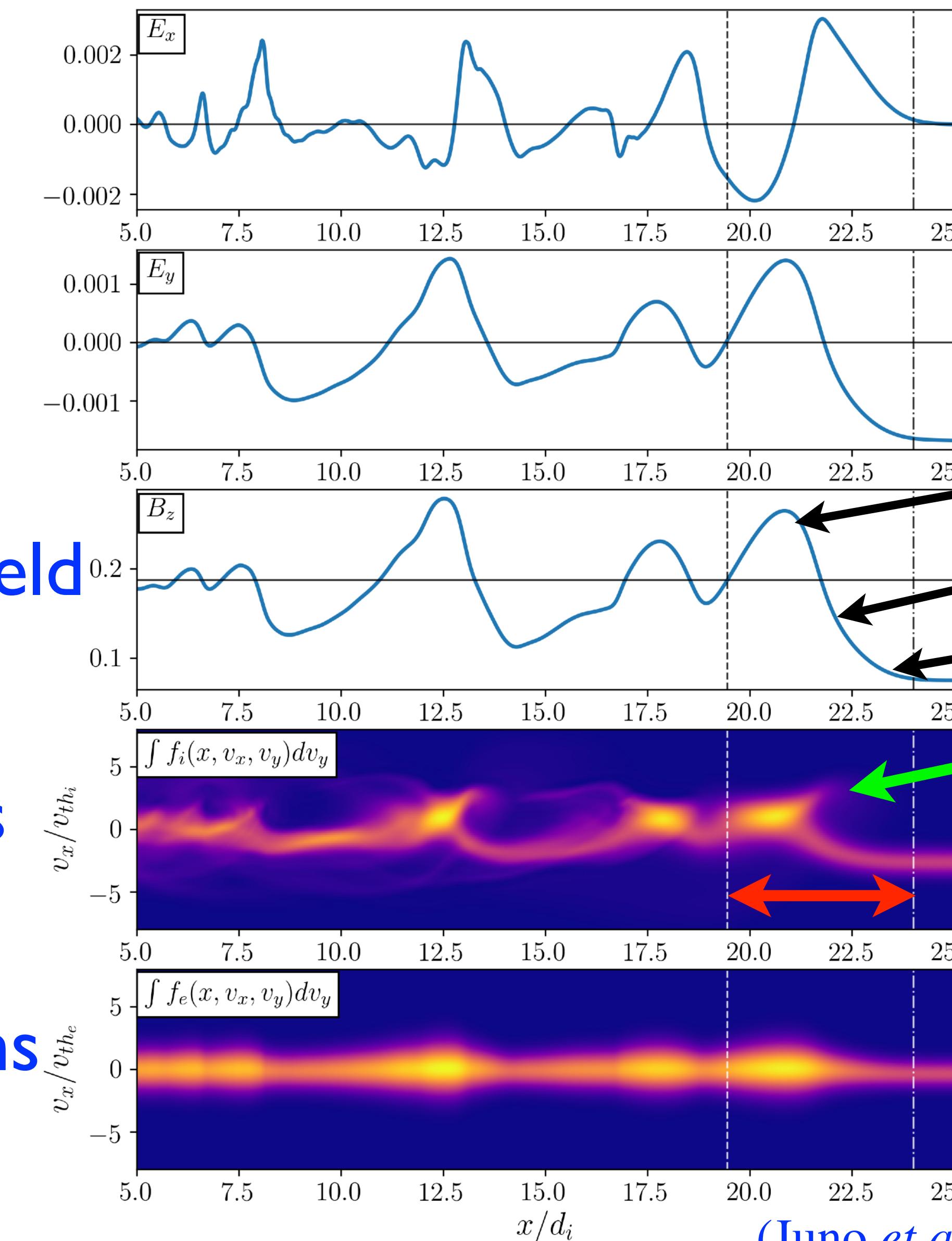
Gkeyll Simulation

- Exactly Perpendicular
 $\theta_{nB} = 90^\circ$
- 1D-2V simulation
- Supercritical
 $M_f \approx 3$
 $M_A \approx 5$
- $\beta_i = 1.3$
- $\beta_e = 0.7$
- $m_i/m_e = 100$

Magnetic Field

Ions

Electrons



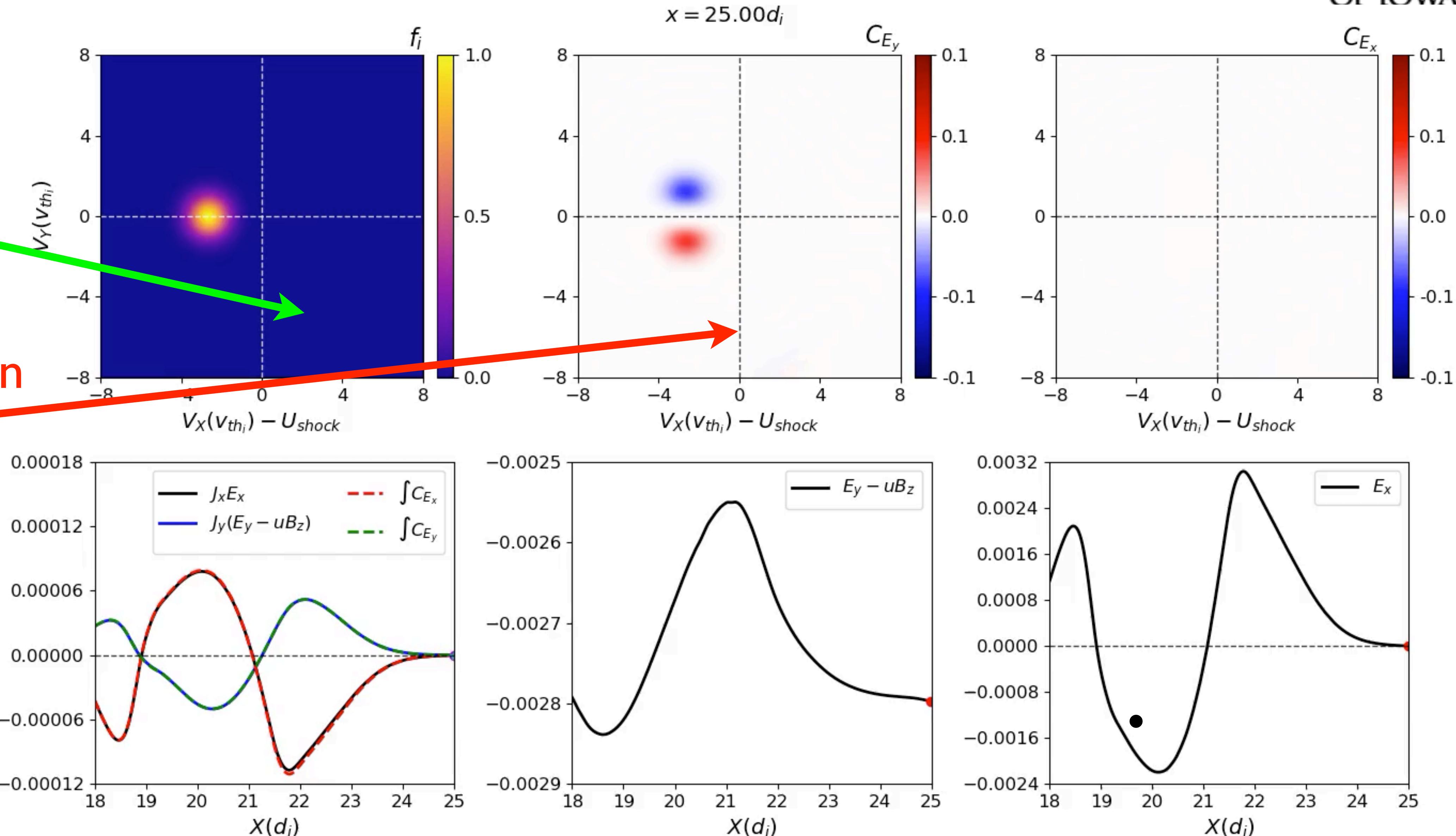
(Juno *et al.*, in prep, 2020)

Distribution Function and Velocity-Space Signatures

Reflected
Ions

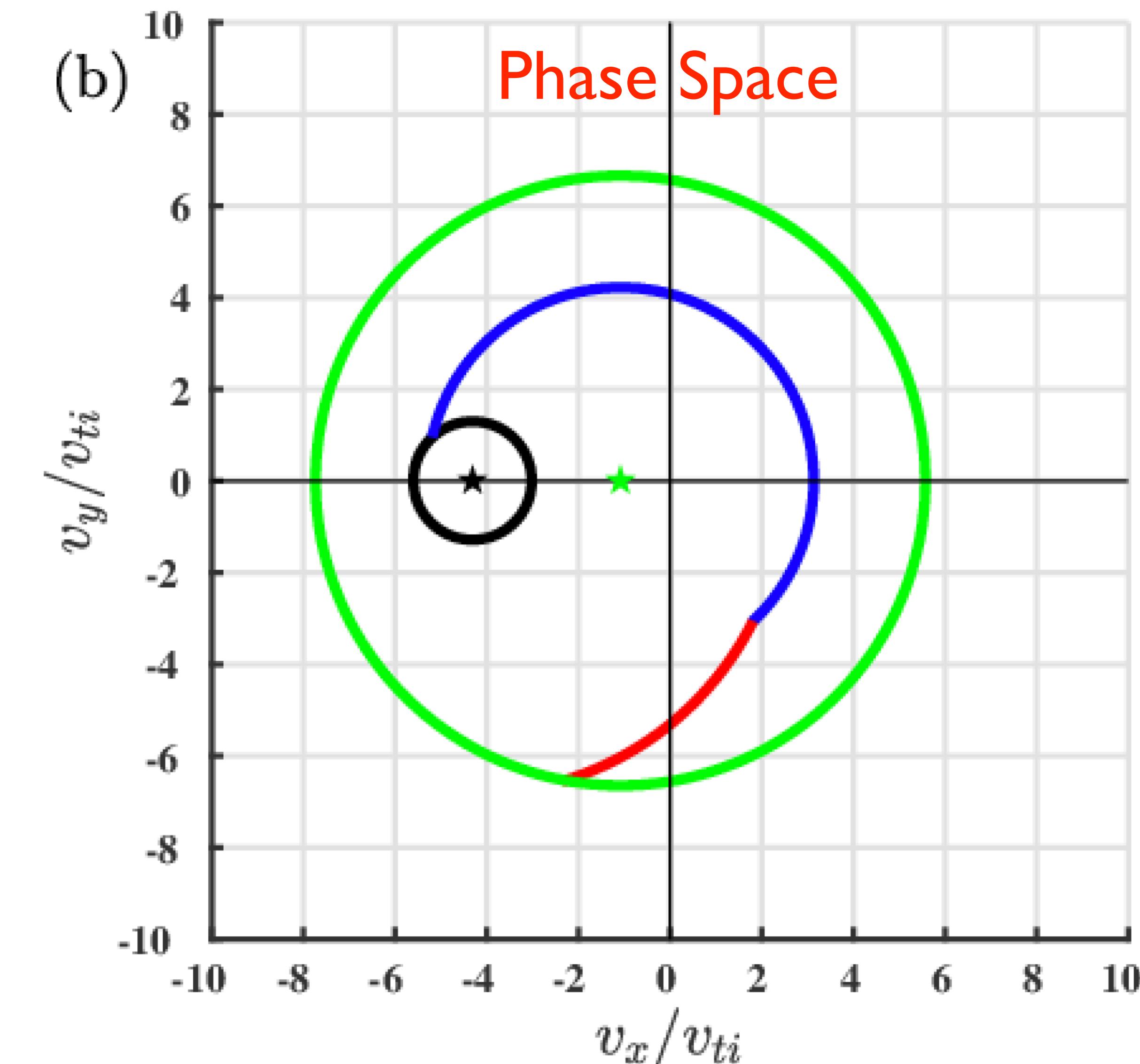
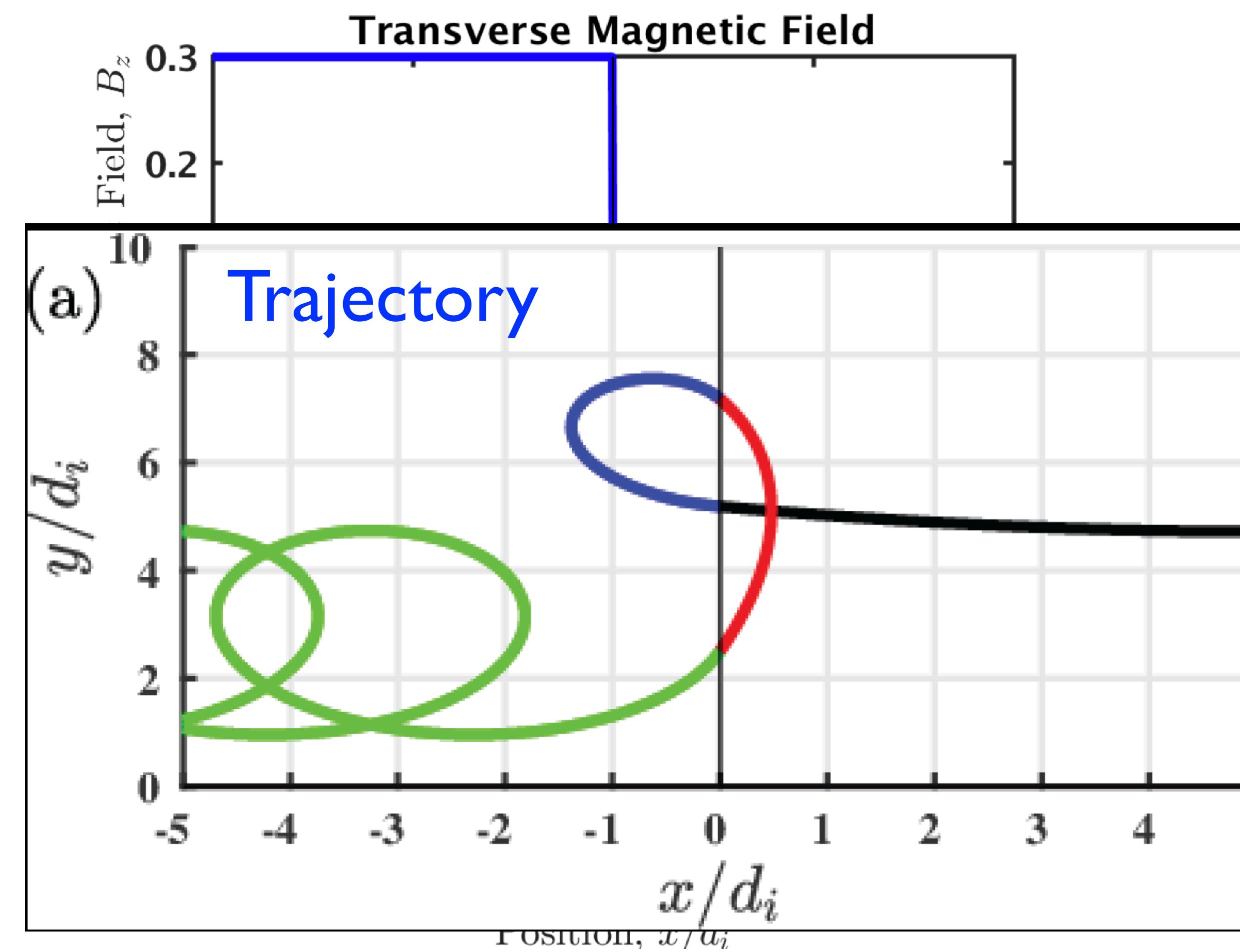
Ion Energization
by motional
electric field

E_y



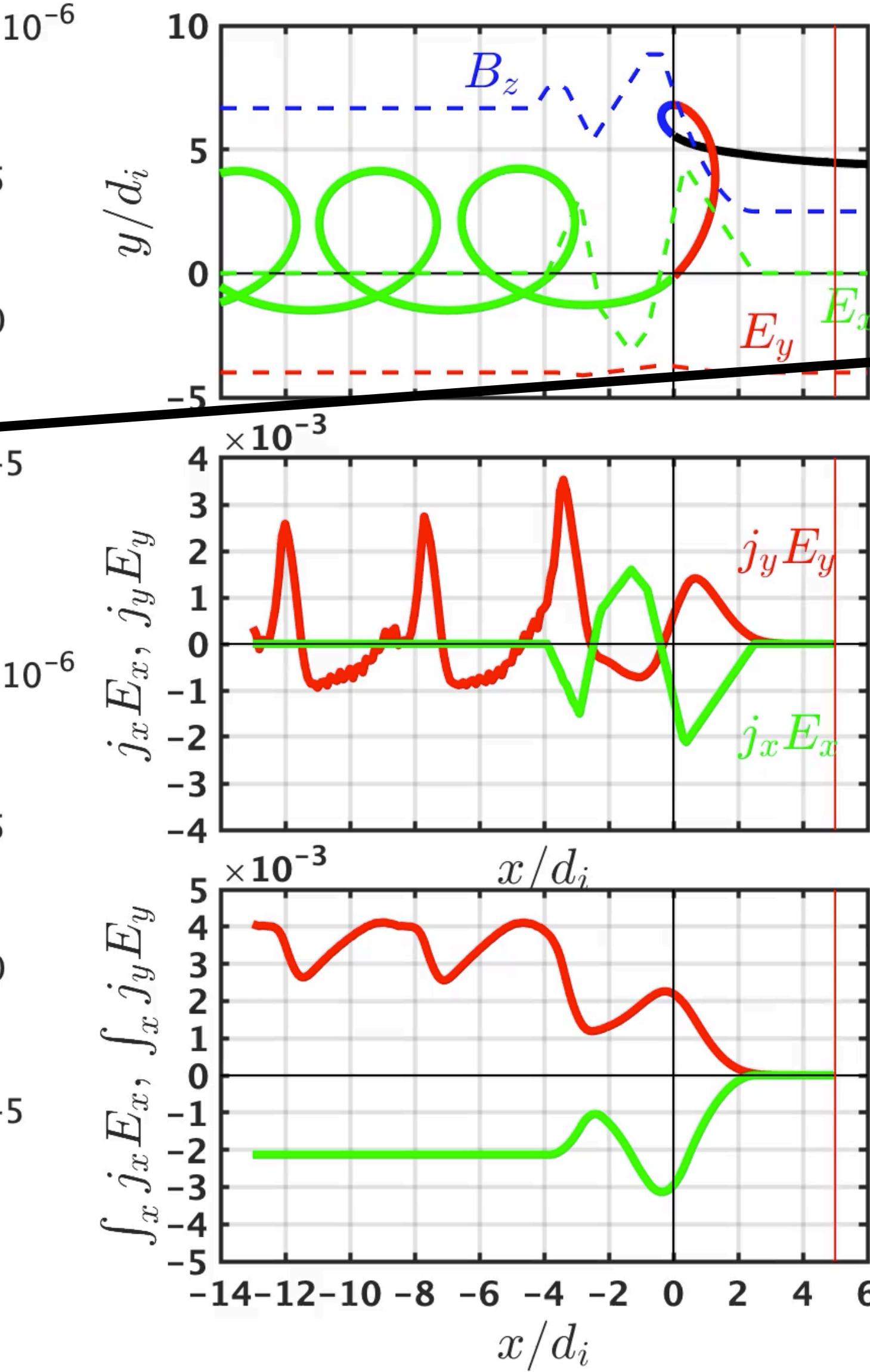
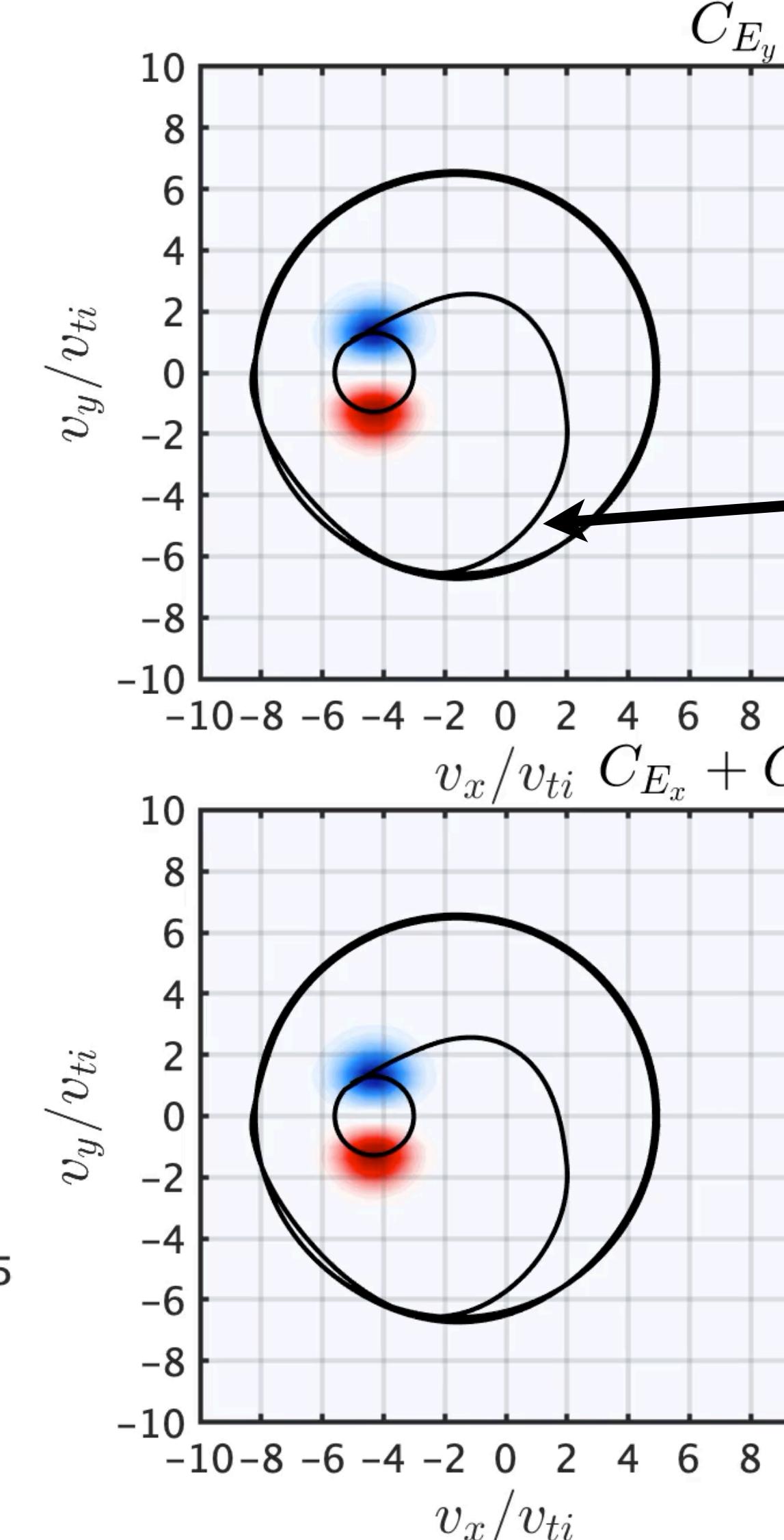
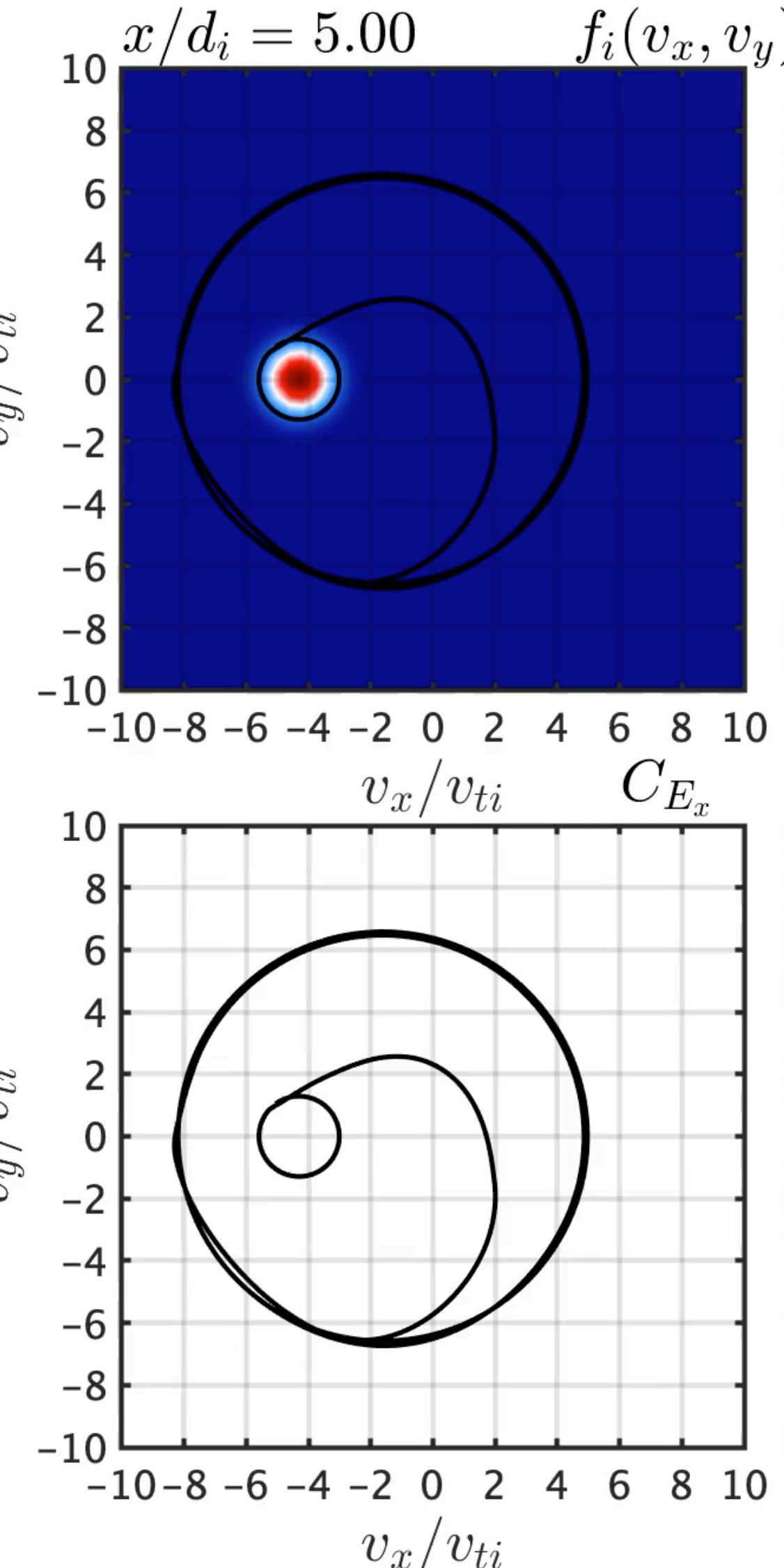
Model of Perpendicular Collisionless Shock

Idealized Perpendicular Shock Model



Model of Perpendicular Collisionless Shock

Single-Particle-Motion Model

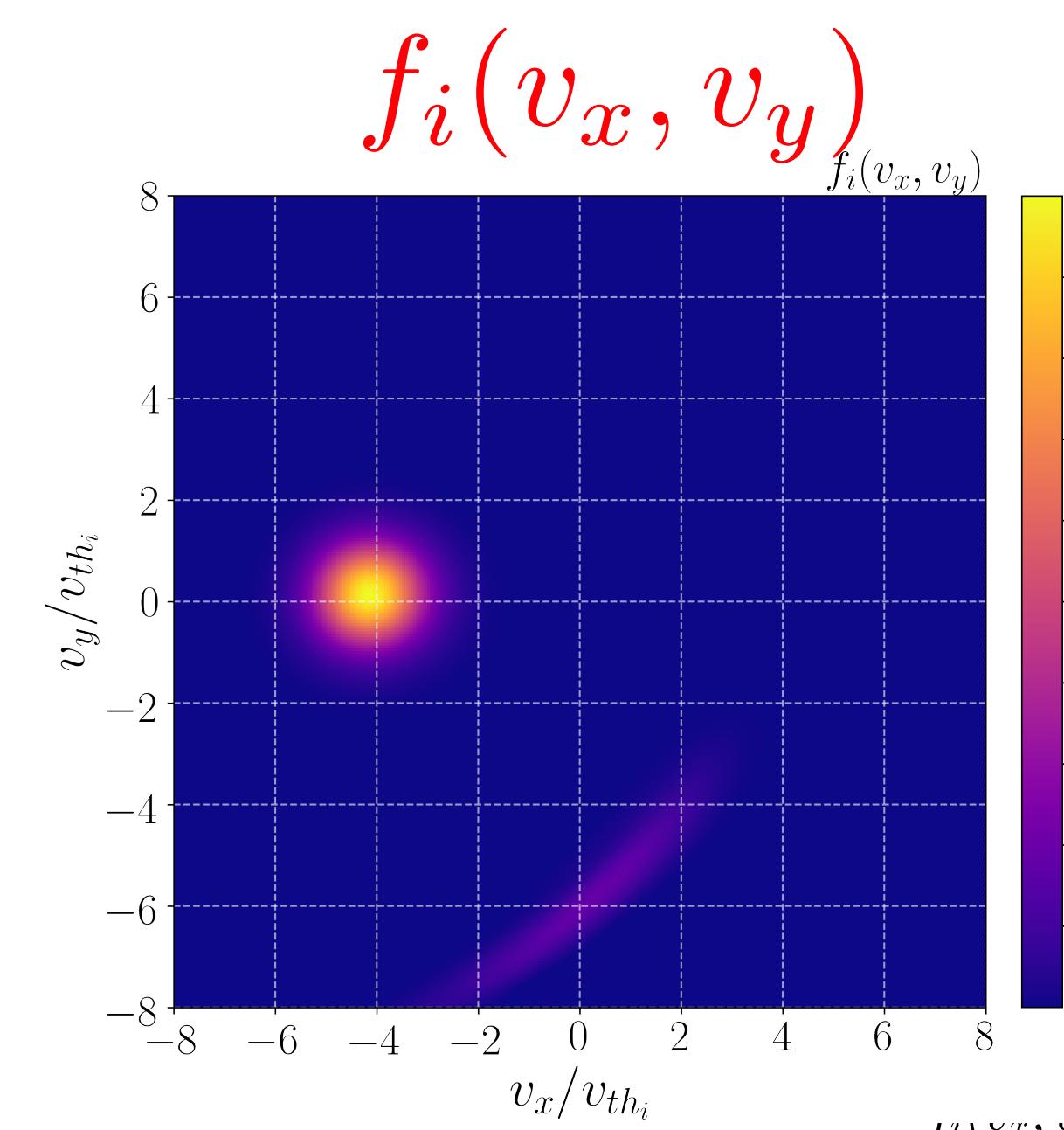


Ion
Energization
by E_y

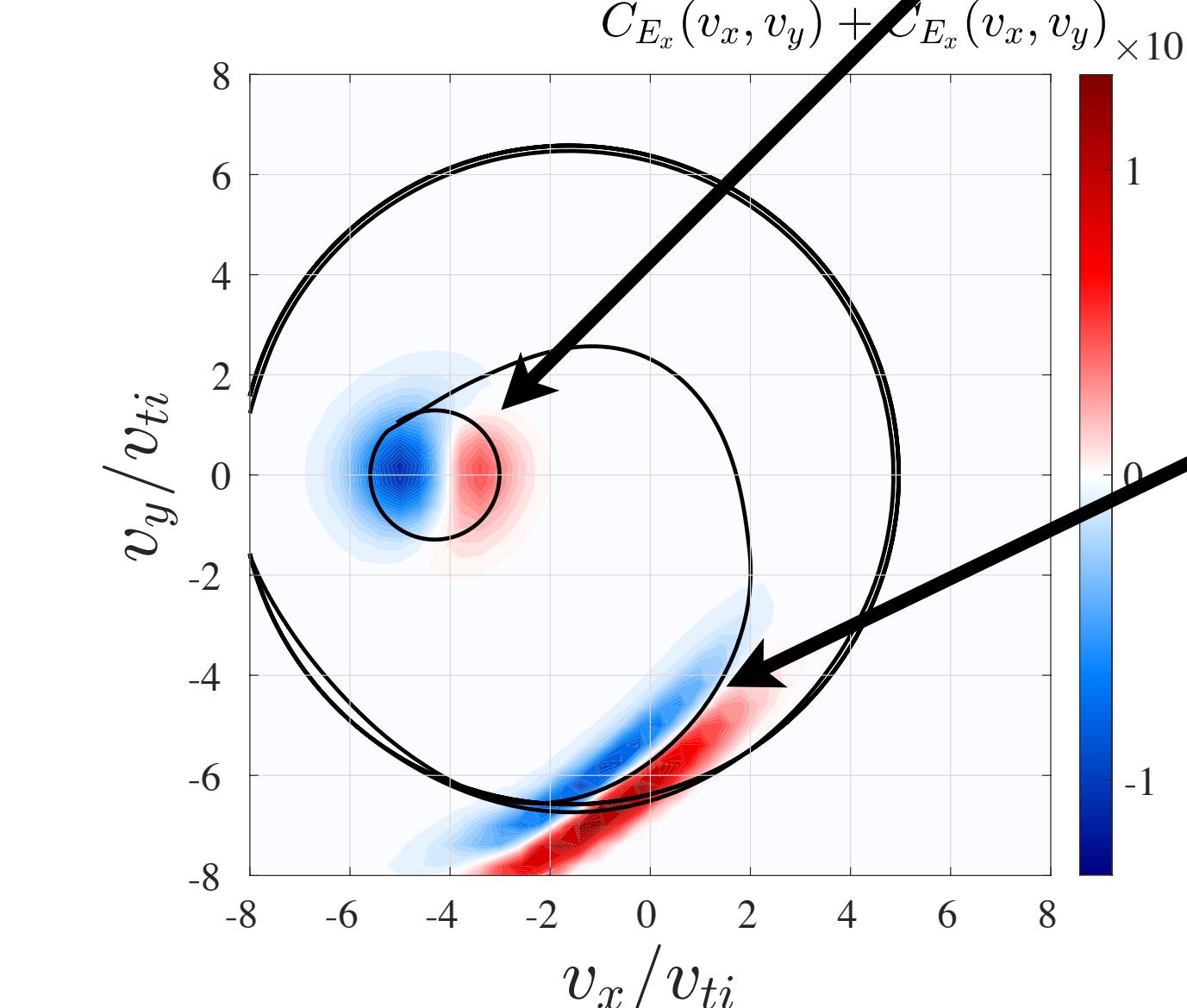
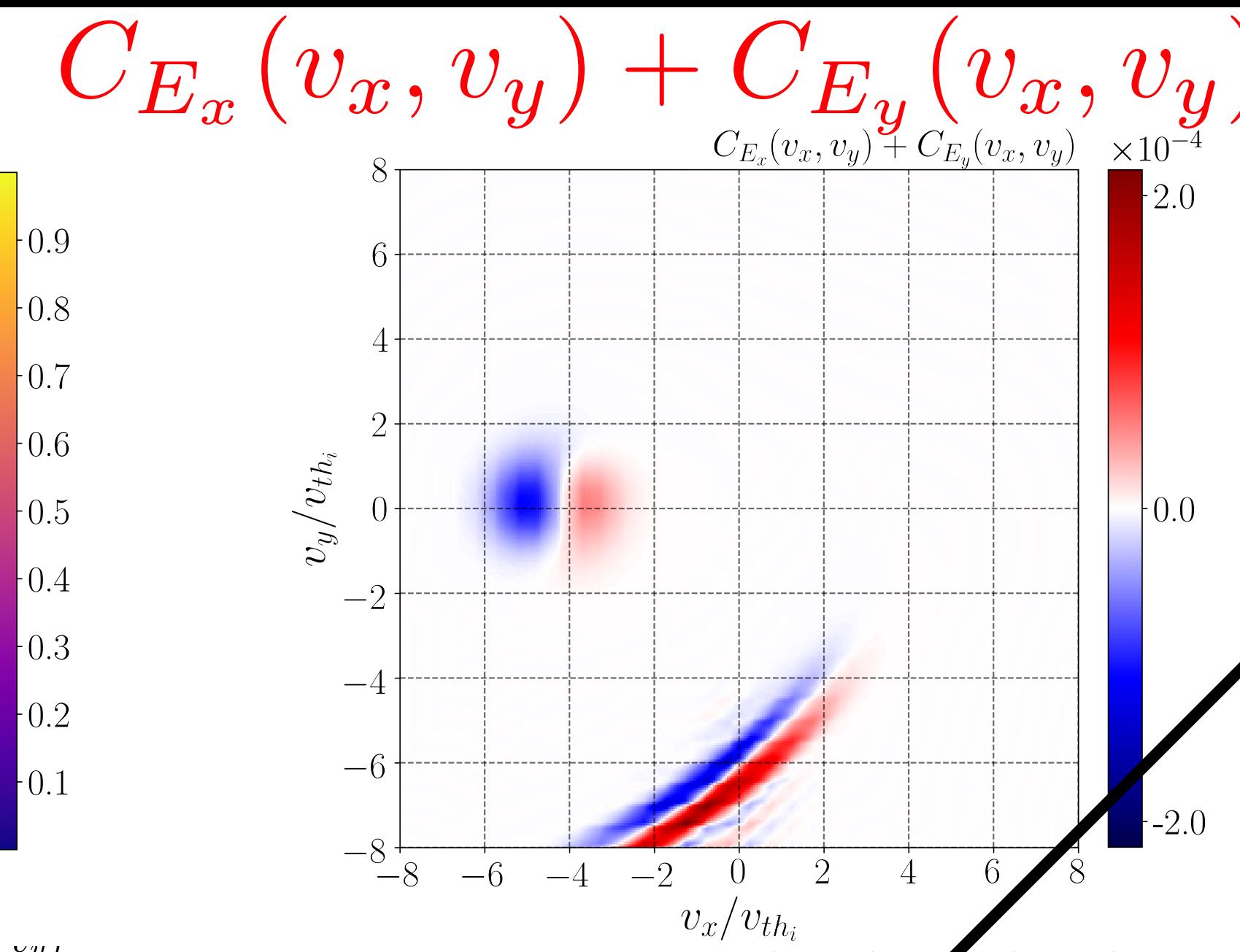
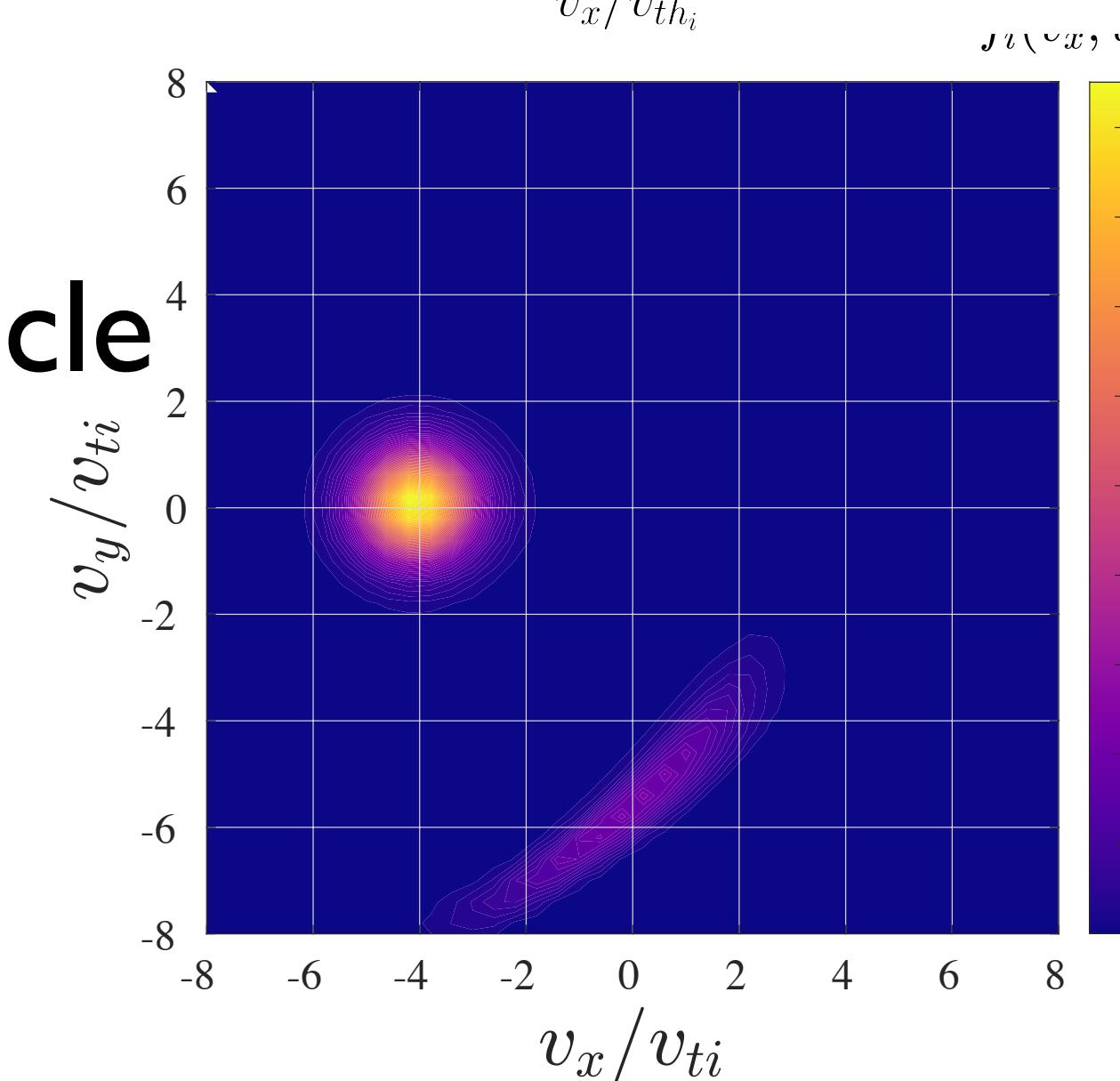
Velocity-space
signature of
Shock Drift
Acceleration

Comparison of Simulation and Model

Gkeyll
Simulation



Single-Particle
Motion
Model



Deceleration of
Incoming beam
by E_x

Shock Drift
Acceleration
by E_y

Outline

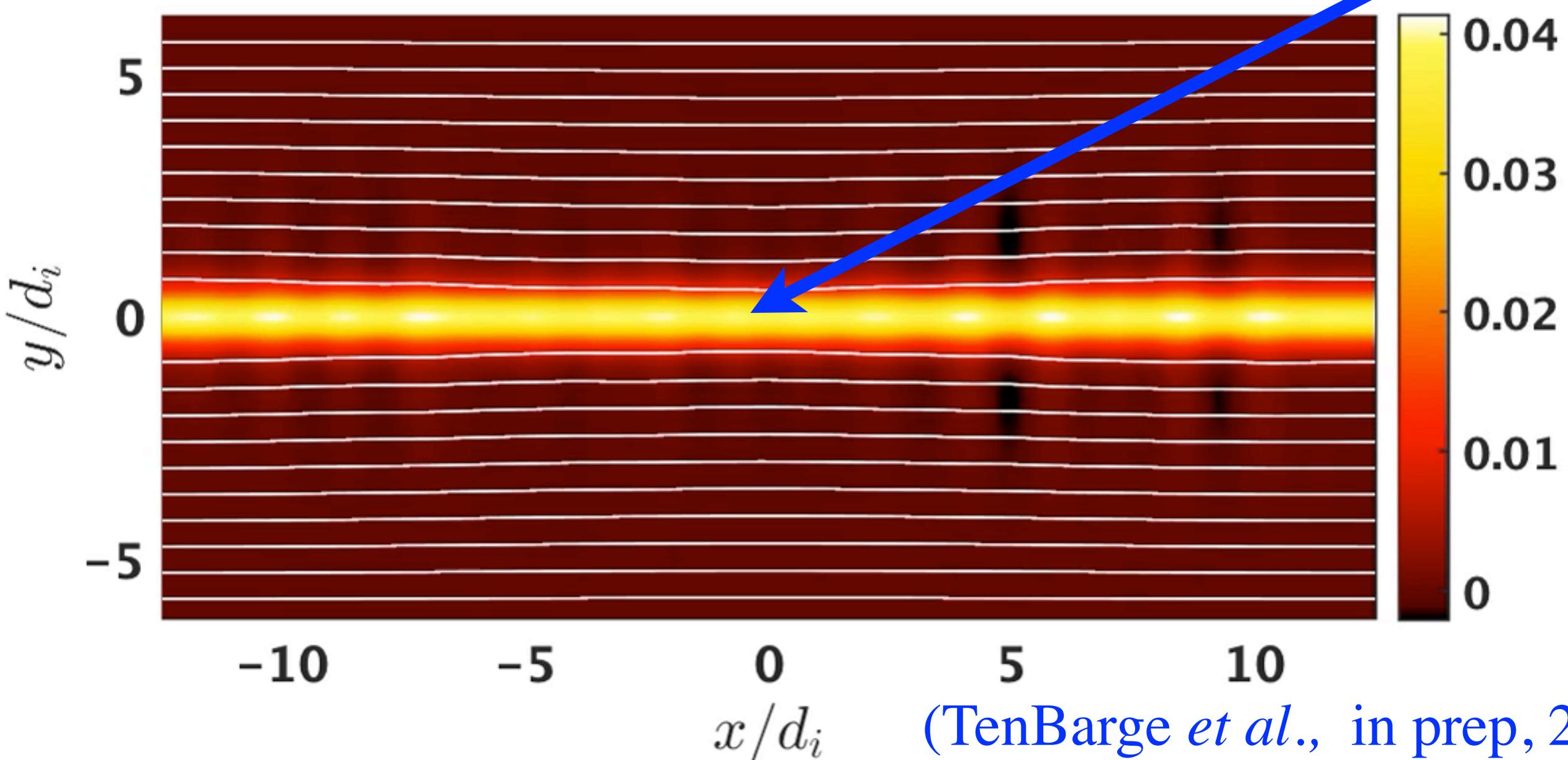
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Collisionless Magnetic Reconnection

Gkeyll Simulation

- Zero Guide-Field Magnetic Reconnection
- 2D, GEM initial conditions

$t = 0000\Omega_{ci}^{-1}$



- $\beta_i = 5/6$
- $T_i/T_e = 5$
- $m_i/m_e = 25$

How are electrons energized in the exhaust?

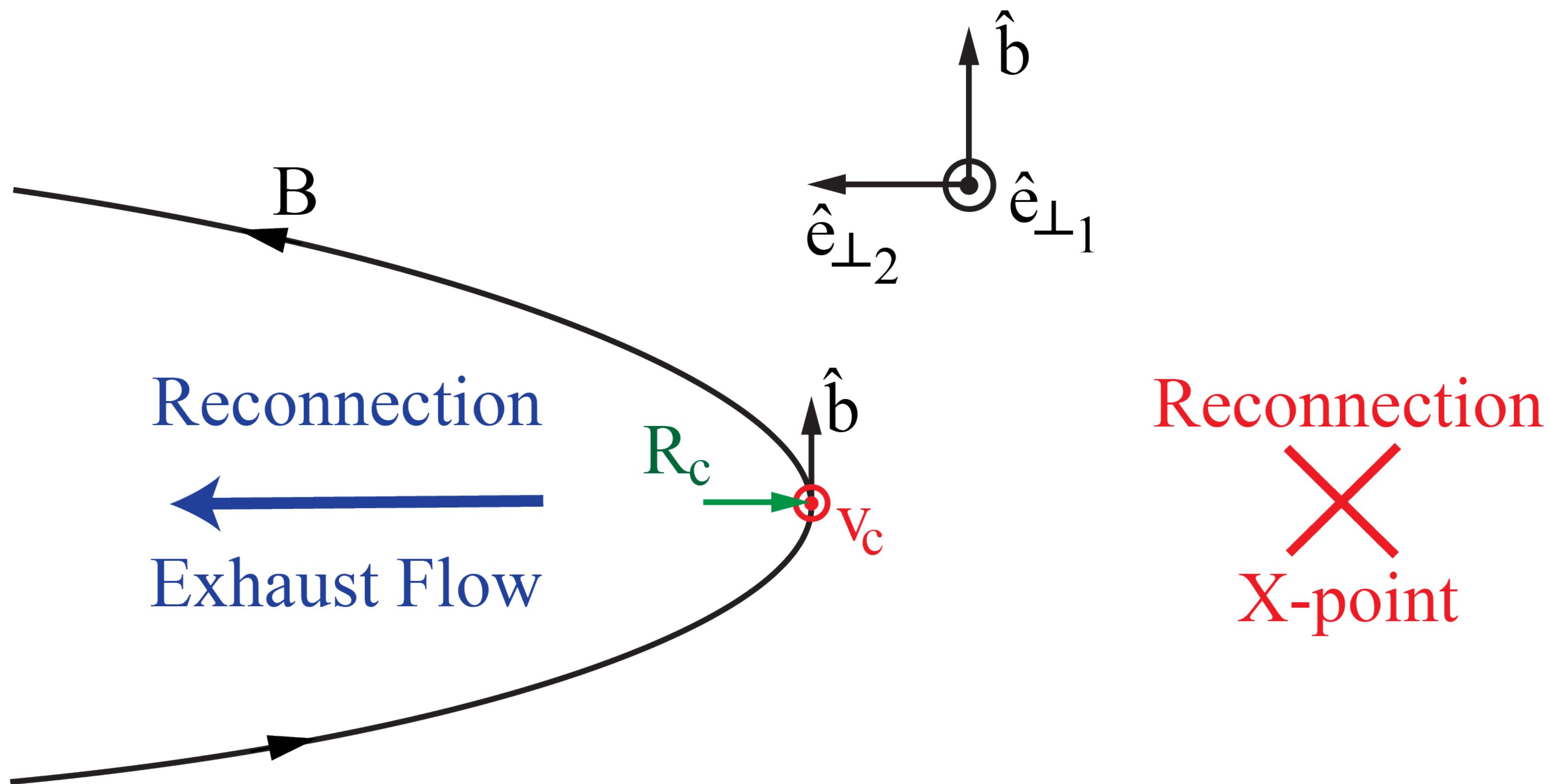
Dahlin, Drake, & Swisdak have demonstrated Fermi acceleration of electrons by the curvature drift

(Dahlin, Drake, & Swisdak, 2014, 2015, 2016, 2017)

What is the velocity-space signature of this curvature drift acceleration?

Collisionless Magnetic Reconnection

Field-Aligned Coordinate System



$$\mathbf{v}_c = \frac{mv_{\parallel}^2}{qB^2} \frac{\mathbf{R}_c \times \mathbf{B}}{R_c^2}$$

- Curvature drift is in the $\hat{\mathbf{e}}_{\perp 1}$ direction
- Compute Correlations

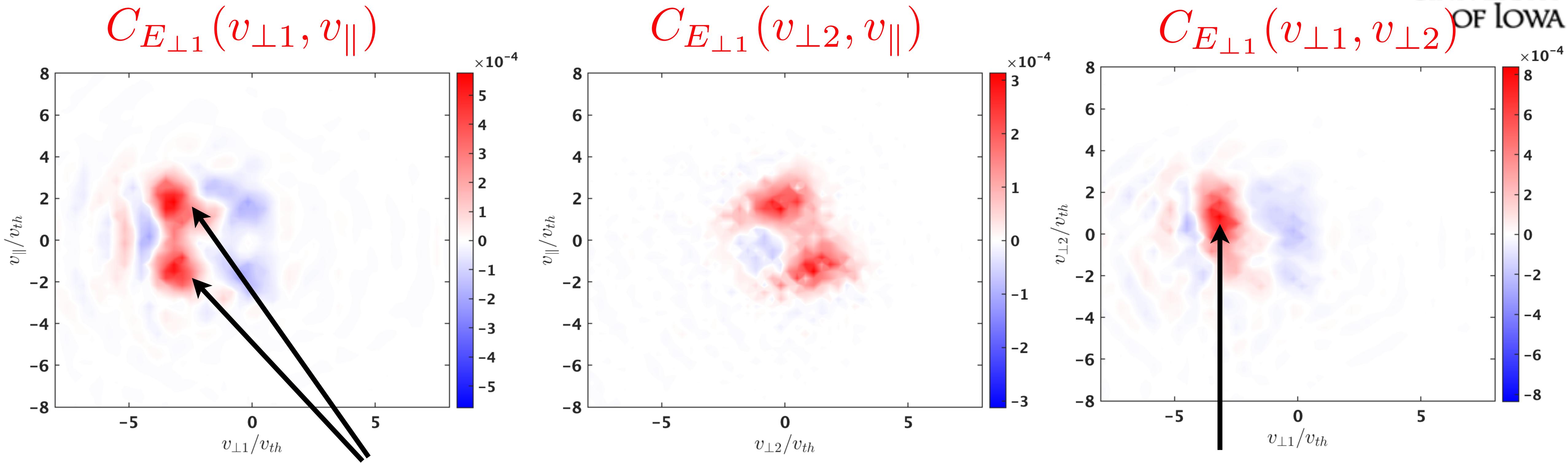
$$C_{E_{\perp 1}}(v_{\perp 1}, v_{\parallel})$$

$$C_{E_{\perp 1}}(v_{\perp 2}, v_{\parallel})$$

$$C_{E_{\perp 1}}(v_{\perp 1}, v_{\perp 2})$$

Velocity-Space Signature of Curvature Drift Acceleration

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Non-zero v_{\parallel} needed for significant curvature drift

Accelerated electrons have $2 \lesssim v_{\perp 1}/v_{te} \lesssim 4$

Tail of the distribution
Electron acceleration!

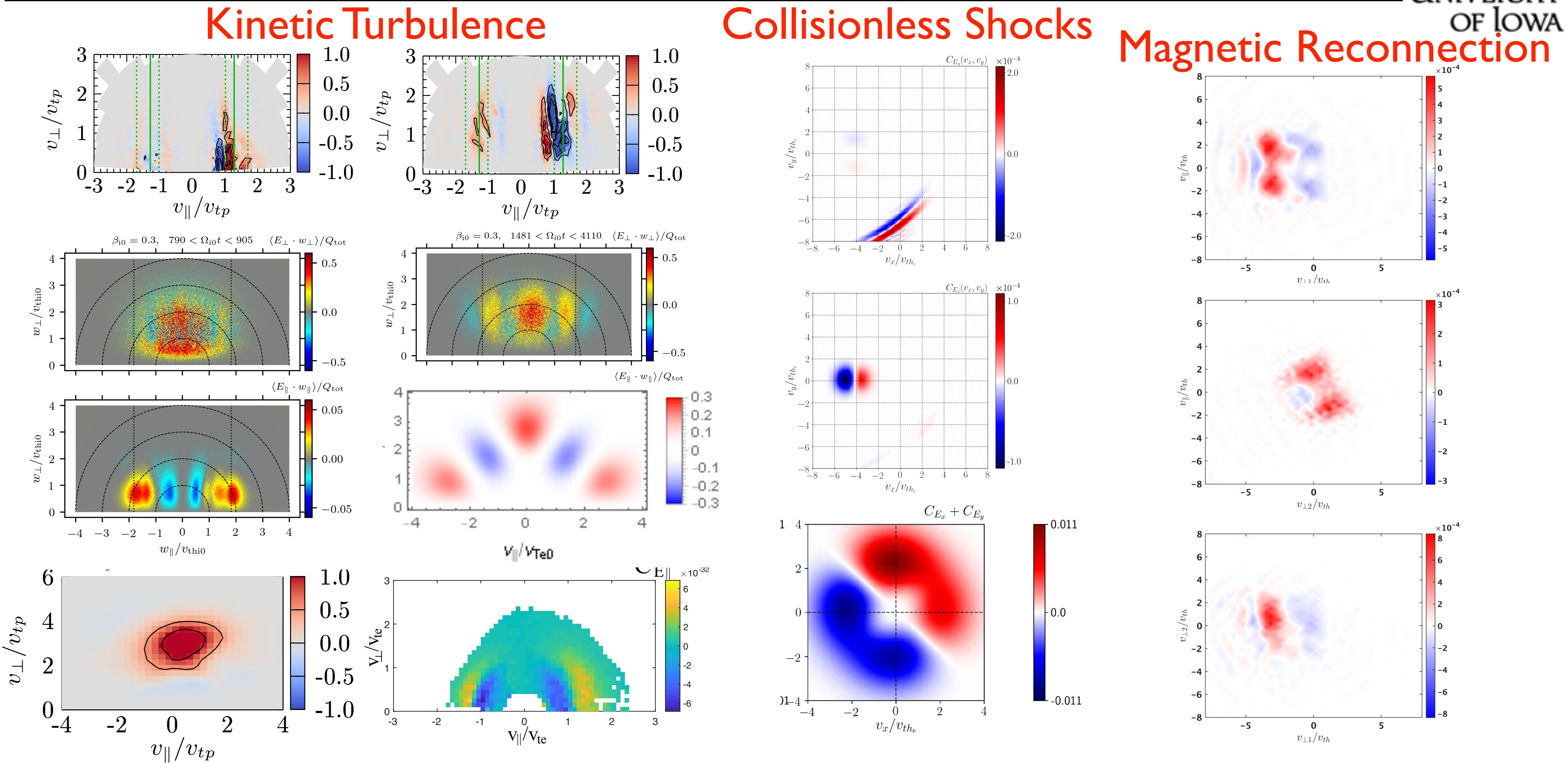
Velocity-space signature of Fermi Acceleration by the Curvature Drift

Outline

- Plasma Heating and Particle Acceleration in the Heliosphere
- Kinetic Theory of Particle Energization
 - Field-Particle Correlation Technique
- Three Applications of the Field-Particle Correlation Technique
 - Plasma Heating by Dissipation of Plasma Turbulence
 - Ion Energization in Collisionless Shocks
 - Electron Fermi Acceleration in Collisionless Magnetic Reconnection
- Constructing a “Rosetta Stone” for particle energization
- Conclusions

A Rosetta Stone for Particle Energization

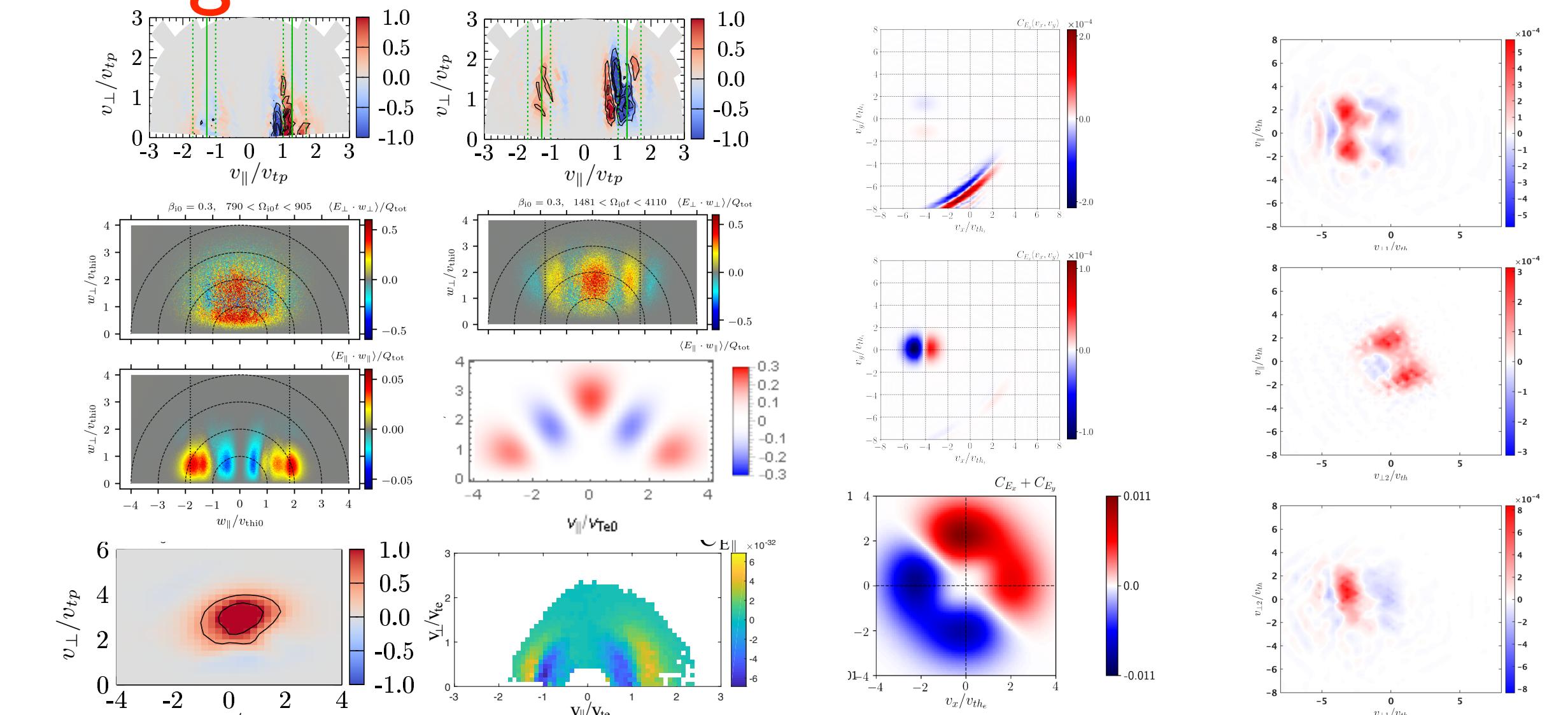
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Conclusions

- Turbulence, reconnection, and shocks govern energy flow in the heliosphere
 - Understanding the mechanisms of particle energization in these processes is essential to develop a predictive understanding of the heliospheric evolution
- Full velocity-space information of 3D-3V phase space is often under-utilized
- The field-particle correlation method fully exploits velocity-space (single-point)
 - Velocity-space signature can be used to identify mechanisms of energization
 - Quantify the rate of particle energization

A Rosetta Stone
for Particle Energization
in Heliospheric Plasmas



The End

Alternative Lagrangian Formulation

$$\frac{dw_s(\mathbf{r}, \mathbf{v}, t)}{dt} = \frac{\partial w_s}{\partial t} + \mathbf{v} \cdot \nabla w_s + q_s \left(\mathbf{E} + \frac{1}{c} \mathbf{v} \times \mathbf{B} \right) \cdot \frac{\partial w_s}{\partial \mathbf{v}} = q_s \mathbf{v} \cdot \mathbf{E} f_s$$

Along particle trajectory in phase space ...

rate of change of phase-space energy density $w_s(\mathbf{r}, \mathbf{v}, t)$ is force times velocity

$$\frac{dW}{dt} = \mathbf{F} \cdot \mathbf{v} = (q\mathbf{E}) \cdot \mathbf{v}$$

Compare to Eulerian Form

$$\frac{\partial w_s(\mathbf{r}, \mathbf{v}, t)}{\partial t} = -\mathbf{v} \cdot \nabla w_s - q_s \frac{v^2}{2} \mathbf{E} \cdot \frac{\partial f_s}{\partial \mathbf{v}} - \frac{q_s}{c} \frac{v^2}{2} (\mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f_s}{\partial \mathbf{v}}$$

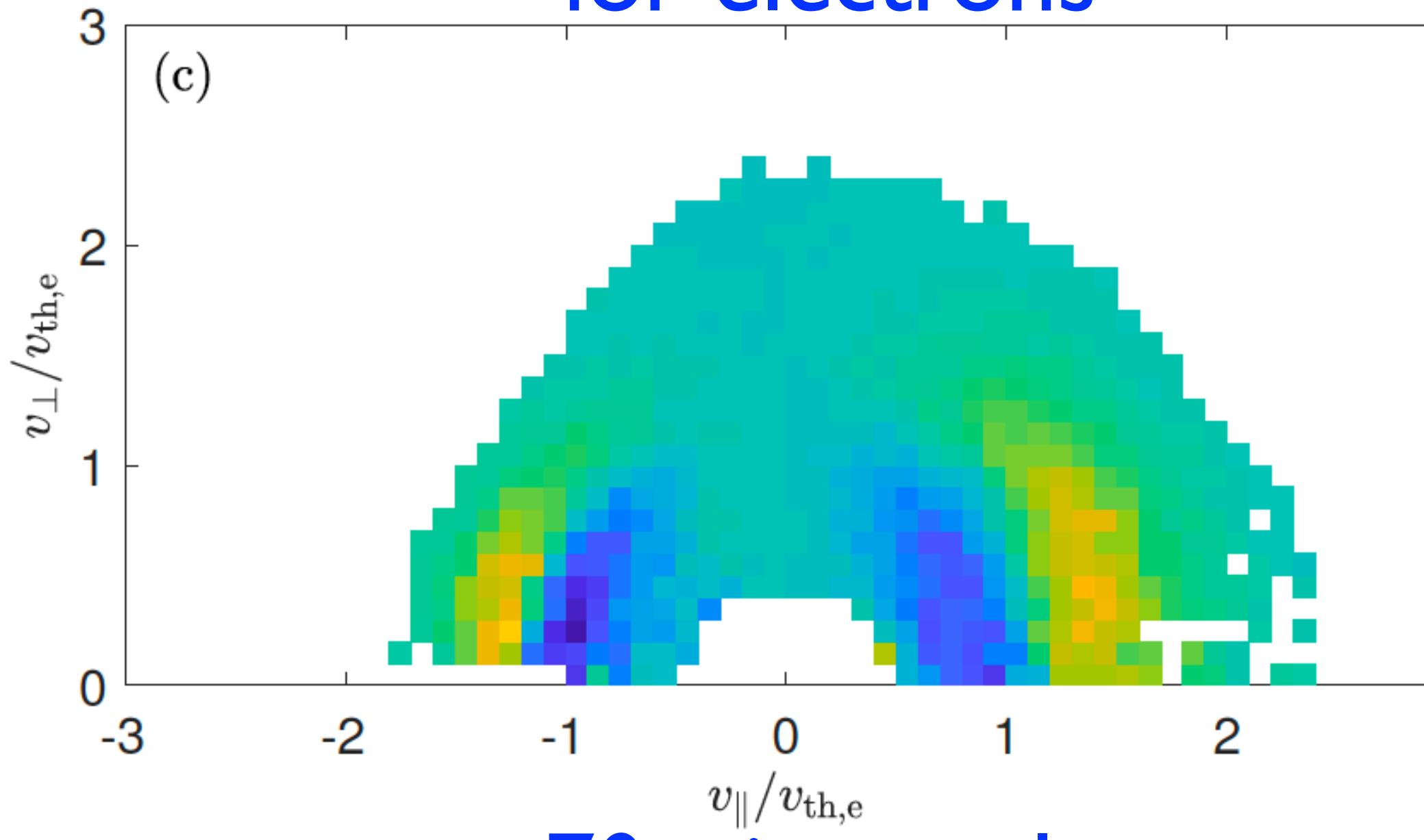
Analysis of Magnetosheath Turbulence with MMS

MMS Observations of Magnetosheath Turbulence

Plasma parameters: $\beta_i = 0.80$

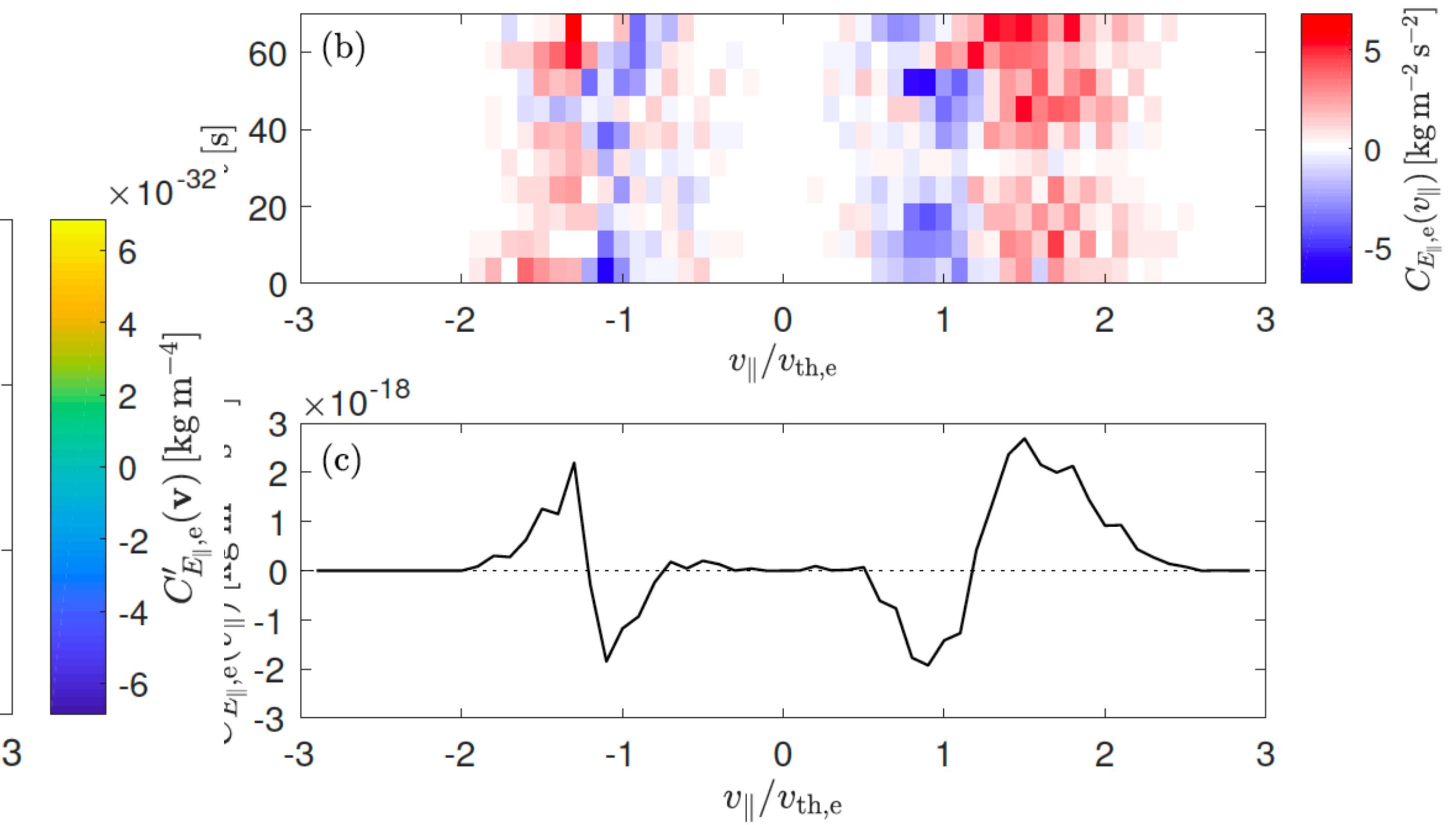
$$T_i/T_e = 10$$

Field-particle correlation
for electrons



70 s interval

(Chen, Klein, & Howes,
Nature Comm., 10:740, 2019)



First definitive evidence of
electron Landau damping in space!