

Lecture #10 Tensor Analysis

I. Tensors

A. Introduction

1. Tensors arise in a number of important areas in physics: general relativity, electrodynamics, stress and strain, moment of inertia.
2. Scalars - tensors of rank 0
Vectors - tensors of rank 1
3. Def: A tensor of rank n in a d -dimensional space:
 - a. Components have n indices, each running 1 to d .
Total of d^n components
 - b. Components of tensor transform in a specified manner under coordinate transformations,

B. Covariant and Contravariant Tensors

1. For a rotational transformation from $(\hat{e}_1, \hat{e}_2, \hat{e}_3)$ to $(\hat{e}'_1, \hat{e}'_2, \hat{e}'_3)$

$$(A_i)' = \sum_j (\hat{e}'_i \cdot \hat{e}_j) A_j = \sum_j \left(\frac{\partial x_j}{\partial x_i'} \right) A_j$$

chain rule to convert A_j into A_j'

2. But the gradient of a scalar ϕ transforms slightly differently.

- a. For $(\nabla\phi)_j = \left(\frac{\partial\phi}{\partial x_j} \right) \hat{e}_j$, $(\nabla\phi)'_i = \frac{\partial\phi}{\partial x_i'} = \sum_j \left(\frac{\partial x_j}{\partial x_i'} \right) \frac{\partial\phi}{\partial x_j}$

3. NOTE: There is a subtle difference in these transformations.

- a. $\left(\frac{\partial x_i'}{\partial x_j} \right)_{x_k}$ $\left(\frac{\partial x_j}{\partial x_i'} \right)_{x_k}$ Different quantities are held fixed!

- b. In Cartesian coordinates, these are the same!

- c. But, in non-Cartesian systems, these transformations generally differ.

I. B. (Continued)

Howes ②

4. Contravariant vs Covariant Vectors:

a. Contravariant: $(A')^i = \sum_j \frac{\partial (x')^i}{\partial x^j} A^j$ Superscript

b. Covariant: $(A')_i = \sum_j \frac{\partial x^j}{\partial (x')^i} A_j$ subscript

5. Einstein Convention for Summation Notation

a. Unsummed index above (i) appears in same position on both sides.

⇒ NOTE: A superscript in the denominator is treated as a subscript.

b. The summed index (j) occurs once as upper, once as lower.

c. Einstein Convention: Omit summation sign. A repeated index appearing as an upper & lower index is assumed to be summed.

$$(A')^i = \frac{\partial (x')^i}{\partial x^j} A^j \leftarrow \begin{array}{l} \text{upper} \\ \text{lower} \end{array} \Rightarrow \text{implicit } \sum_{j=1}^d$$

C. Tensors of Rank 2 (3x3 Matrices)

1. When dealing with tensors, summation notation is often easier to manage than matrix multiplication notation.

2. Different Types:

a. Contravariant: $(A')^{ij} = \frac{\partial (x')^i}{\partial x^k} \frac{\partial (x')^j}{\partial x^l} A^{kl}$ k, l repeated ⇒ $\sum_{k,l}$

b. Mixed: $(B')^i_j = \frac{\partial (x')^i}{\partial x^k} \frac{\partial x^l}{\partial (x')^j} B^k_l$

c. Covariant: $(C)_{ij} = \frac{\partial x^k}{\partial (x')^i} \frac{\partial x^l}{\partial (x')^j} C_{kl}$

I.C. (Continued)

Howes ③

3. Invariance under transformations is used to express universal physical laws \Rightarrow This makes tensor analysis important.

4. Matrix Form

$$a. \tilde{A} = \begin{pmatrix} A^{11} & A^{12} & A^{13} \\ A^{21} & A^{22} & A^{23} \\ A^{31} & A^{32} & A^{33} \end{pmatrix}$$

b. Similarity transformation

$$(A')^{ij} = \sum_{kl} S_{ik} A^{kl} (S^T)_{lj} \Leftrightarrow \tilde{A}' = \tilde{S} \tilde{A} \tilde{S}^T$$

D. Tensor Properties and Operations

1. Addition as expected for matrices, $A^{ij} + B^{ij} = C^{ij}$

2. Symmetry: a. Symmetric if $A^{mn} = A^{nm}$

b. Anti-symmetric if $A^{mn} = -A^{nm}$

c. Decomposition:

$$A^{mn} = \underbrace{\frac{1}{2}(A^{mn} + A^{nm})}_{\text{Symmetric}} + \underbrace{\frac{1}{2}(A^{mn} - A^{nm})}_{\text{Anti-symmetric}}$$

3. Kronecker-Delta, Mixed rank 2 tensor, δ_l^k

b. Does it transform properly? $(\delta')^i_j = \frac{\partial(x')^i}{\partial x^k} \frac{\partial x^l}{\partial(x')^j} \delta_l^k = \frac{\partial(x')^i}{\partial x^k} \frac{\partial x^k}{\partial(x')^j} = \frac{\partial(x')^i}{\partial(x')^j} = (\delta^j_i)^i_j$

c. Isotropic: δ_l^k has same components in all rotated coordinates.

4. Contraction: a. For Vectors, $\underline{A} \cdot \underline{B} = \sum_i A_i B_i$ $i \leftarrow$ upper

b. For tensor, set two indices equal to each other, B_i^i , and sum according to Einsteinian convention $i \leftarrow$ lower

$$c. \text{Ex: } (B')^i_i = \frac{\partial(x')^i}{\partial x^k} \frac{\partial x^l}{\partial(x')^i} B_l^k = \frac{\partial x^l}{\partial x^k} B_l^k = \delta_l^k B_l^k = B_k^k$$

$$= \sum_k B_k^k$$

d. Contraction Reduces a tensor rank by 2 \Rightarrow Rank 2 Tensor \rightarrow Scalar.

e. NOTE: In Matrix Analysis, $B_k^k \Rightarrow \sum_k B_k^k = B_1^1 + B_2^2 + B_3^3 = \text{Trace}(\underline{B})$

I. D. (Continued)

Howes (4)

5. Direct Product: (Sum ranks of tensors)

a. $C_{kem}^{ij} = A_k^i B_{em}^j$ - Simply multiply, compare by comparison.
 - Easy in summation notation.

Rank: $5 = 2 + 3$

b. Contravariance or covariance of each index must be maintained.

6. Inverse Transformation:

a. $(A')^j = \frac{\partial(x')^j}{\partial x^i} A^i \xrightarrow{\text{inverse}} A^i = \frac{\partial x^i}{\partial(x')^j} (A')^j$

b. BUT \Rightarrow \uparrow
 x^k fixed \uparrow
 $(x')^k$ fixed

c. Check Inverse: $\sum_k \frac{\partial(x')^k}{\partial x^i} [A^i] = \frac{\partial(x')^k}{\partial x^i} \left[\frac{\partial x^i}{\partial(x')^j} (A')^j \right] = \frac{\partial(x')^k}{\partial(x')^j} (A')^j = \delta_j^k (A')^j = (A')^k$

7. Quotient Rule: Used to establish tensor nature.

a. If A_{ij} & B^{ke} are tensors, then direct product $C_{ij}^{ke} = A_{ij} B^{ke}$ is a tensor

b. What about inverse problem?

$$K_{ke} C_{ij}^{ke} = A_{ij} \quad (K_{ke} \text{ is inverse of } B^{ke})$$

c. Quotient Rule: If the equation of inverse holds in all coordinate systems, then \underline{K} is a tensor of indicated rank & contra/covariant character.

E. Pseudotensors:

1. For vectors, $\underline{A}' = \underline{S} \underline{A}$ vector
 $\underline{A}' = \det(\underline{S}) \underline{S} \underline{A}$ pseudovector

2. Pseudotensors are the extension of this concept.

a. Require additional sign factor in transformation rule.

I. E (continued)

Notes ⑤

3. Quotient Rule and psuedotensors: $\underline{T} \rightarrow$ tensor
 $\underline{P} \rightarrow$ psuedo tensor:

a. $\underline{T} \otimes \underline{T} = \underline{T}$ b. $\underline{P} \otimes \underline{T} = \underline{P}$
 $\underline{P} \otimes \underline{P} = \underline{T}$ $\underline{T} \otimes \underline{P} = \underline{P}$

4. Ex: Levi-Civita Symbol ϵ_{ijk}

a. ϵ_{ijk} is a rank-3 psuedotensor and is isotropic

II. Tensors in General Coordinates

A. The Metric Tensor

1. For non-Cartesian coordinate systems, we want to develop a systematic way of handling these more general metric spaces.

2. Consider a general (3D) coordinate system (q_1, q_2, q_3)

a. Covariant basis vectors $\underline{\epsilon}_i$

$$\underline{\epsilon}_i = \frac{\partial \underline{x}}{\partial q_i} \hat{\underline{e}}_x + \frac{\partial \underline{y}}{\partial q_i} \hat{\underline{e}}_y + \frac{\partial \underline{z}}{\partial q_i} \hat{\underline{e}}_z$$

b. Arbitrary Vector: $\underline{A} = A^1 \underline{\epsilon}_1 + A^2 \underline{\epsilon}_2 + A^3 \underline{\epsilon}_3$

c. NOTE: i. \underline{A} is fixed object (unchanged by choice of coordinates)

ii. A^i coefficients form a contravariant vector

iii. $\underline{\epsilon}_i$ is a covariant basis vector.

iv. When coordinate system is changed, A^i and $\underline{\epsilon}_i$ change in complementary ways to yield a fixed \underline{A} .

II A. (Continued)

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3. $(ds)^2 = \sum_{ij} (\underline{\xi}_i da^i) (\underline{\xi}_j da^j) = \underbrace{g_{ij} da^i da^j}_{\text{summation convention.}}$

4. Covariant Metric Tensor: a. $\boxed{g_{ij} = \underline{\xi}_i \cdot \underline{\xi}_j}$

b. Since $(ds)^2$ is a scalar, quadratic rule tells us that g_{ij} must be a rank 2 covariant tensor.

5. Raising and Lowering Operations

a. Define: Contravariant Metric Tensor g^{ij} by $g^{ik} g_{kj} = \delta^i_j$

b. g^{ij} is the inverse of g_{ij} .

c. Convert contravariant to covariant tensor $g_{ij} F^j = F_i$ "Lowering"

d. Convert covariant to contravariant tensor $g^{ij} F_j = F^i$ "Raising"

6. Ex: a. $\underline{A} = A^i \underline{\xi}_i = A^i [\delta^i_k] \underline{\xi}_k = A^i [g_{ij} g^{jk}] \underline{\xi}_k = A_j \underline{\xi}^j$

b. Thus, the same vector can be represented in either contravariant or covariant bases.

7. To wrap up, we provide the following definitions.

a. Contravariant basis vectors $\underline{\xi}^i = \frac{\partial x^i}{\partial x} \hat{e}_x + \frac{\partial x^i}{\partial y} \hat{e}_y + \frac{\partial x^i}{\partial z} \hat{e}_z$

b. Contravariant metric tensor $\boxed{g^{ij} = \underline{\xi}^i \cdot \underline{\xi}^j}$

II. A. (Continued)

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8. Ex: Spherical Polar Metric Tensor: $(q^1, q^2, q^3) = (r, \theta, \phi)$

a. $x = r \sin \theta \cos \phi$, $y = r \sin \theta \sin \phi$, $z = r \cos \theta$

b. Covariant Basis Vectors:

$$\underline{\xi}_r = \sin \theta \cos \phi \hat{e}_x + \sin \theta \sin \phi \hat{e}_y + \cos \theta \hat{e}_z$$

$$\underline{\xi}_\theta = r \cos \theta \cos \phi \hat{e}_x + r \cos \theta \sin \phi \hat{e}_y - r \sin \theta \hat{e}_z$$

$$\underline{\xi}_\phi = -r \sin \theta \sin \phi \hat{e}_x + r \sin \theta \cos \phi \hat{e}_y$$

c. Can also compute contravariant basis vectors from

$$r \cong \sqrt{x^2 + y^2 + z^2}, \quad \cos \theta = \frac{z}{r}, \quad \tan \phi = \frac{y}{x}$$

d. Covariant Metric Tensor

$$g_{11} = \underline{\xi}_r \cdot \underline{\xi}_r = 1$$

$$g_{22} = \underline{\xi}_\theta \cdot \underline{\xi}_\theta = r^2$$

$$g_{33} = \underline{\xi}_\phi \cdot \underline{\xi}_\phi = r^2 \sin^2 \theta$$

$$\Rightarrow (g_{ij}) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & r^2 & 0 \\ 0 & 0 & r^2 \sin^2 \theta \end{pmatrix}$$

e. Similarly contravariant metric tensor

$$(g^{ij}) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & r^{-2} & 0 \\ 0 & 0 & (r \sin \theta)^{-2} \end{pmatrix}$$

9. Ex: Minkowski Special Relativity Metric (4-vectors)

$$(g_{ij}) = (g^{ij}) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \left. \begin{array}{l} \leftarrow \text{time} \\ \left. \vphantom{\begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}} \right\} \text{3D space} \end{array} \right\}$$

B. Covariant Derivatives

1. Much more complicated because $\underline{\xi}_i$ are, in general, not constant!

2a. Begin with transformation law: $(V^i)' = \frac{\partial x^i}{\partial q^k} V^k$

b. Differentiate:

$$\frac{\partial (V^i)'}{\partial q^j} = \frac{\partial x^i}{\partial q^k} \frac{\partial V^k}{\partial q^j} + \frac{\partial^2 x^i}{\partial q^j \partial q^k} V^k$$

III. B2 (Continued)

Hawes (8)

c. Write as a vector equation $\underline{\xi}_k = \frac{\partial x^i}{\partial q^k}$, so
in x^i coordinates:

$$\underline{V} = V_x \hat{e}_x + V_y \hat{e}_y + V_z \hat{e}_z \quad \frac{\partial \underline{V}'}{\partial q^j} = \frac{\partial V^k}{\partial q^j} \underline{\xi}_k + V^k \frac{\partial \underline{\xi}_k}{\partial q^j}$$

d. NOTE: i. $\frac{\partial \underline{\xi}_k}{\partial q^j}$ is a vector in the space spanned by $\underline{\xi}_i$.

ii. Thus, $\frac{\partial \underline{\xi}_k}{\partial q^j} = \Gamma_{jk}^{\mu} \underline{\xi}_{\mu}$ Christoffel symbol of second kind

3. Christoffel Symbol: a. $\Gamma_{jk}^m = \underline{\xi}^m \cdot \frac{\partial \underline{\xi}_k}{\partial q^j}$

b. NOTE: $\Gamma_{kj}^m = \Gamma_{jk}^m$

4. Thus, $\frac{\partial \underline{V}'}{\partial q^j} = \frac{\partial V^k}{\partial q^j} \underline{\xi}_k + V^k \Gamma_{jk}^{\mu} \underline{\xi}_{\mu}$

5. Covariant Derivative: a. $\frac{\partial \underline{V}'}{\partial q^j} = \left(\frac{\partial V^k}{\partial q^j} + V^{\mu} \Gamma_{j\mu}^k \right) \underline{\xi}_k$

b. DEF: $V_{ij}^k \equiv \frac{\partial V^k}{\partial q^j} + V^{\mu} \Gamma_{j\mu}^k$

c. Mixed, Second Rank Tensor: i. Only combination V_{ij}^k transforms as a tensor.

ii. $\frac{\partial V^k}{\partial q^j}$ and $V^{\mu} \Gamma_{j\mu}^k$ do not individually transform as tensors.

d. Includes changes in basis vectors $\underline{\xi}_k$ as q^i changes (dq^j) to determine derivative.