

Lecture #14: Fourier SeriesI. Fourier SeriesA. General Properties

1. Fourier series are a valuable tool for solving ODEs or PDEs with periodic boundary conditions

2. Consider a piecewise regular function $f(x)$ on $[0, 2\pi]$,
 ↗ may contain finite discontinuities!

a. Define: Fourier Series

$$f(x) = \sum_{n=-\infty}^{\infty} c_n e^{inx}$$

b. To determine complex coefficients c_n , multiply by e^{-inx} and $\int_0^{2\pi}$

$$\int_0^{2\pi} f(x) e^{-inx} dx = \sum_{n'=-\infty}^{\infty} c_{n'} \int_0^{2\pi} e^{in'x} e^{-inx} dx$$

c. Note, $\int_0^{2\pi} e^{i(n'-n)x} dx = 2\pi \delta_{nn'}$

d. Thus $= \sum_{n'=-\infty}^{\infty} c_{n'} 2\pi \delta_{nn'} = 2\pi c_n$

e. So $c_n = \frac{1}{2\pi} \int_0^{2\pi} f(x) e^{-inx} dx$

3. Alternatively, we may define Fourier series by real functions $\cos nx$ & $\sin nx$:

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos nx + \sum_{n=1}^{\infty} b_n \sin nx$$

$$a_n = \frac{1}{\pi} \int_0^{2\pi} f(x) \cos nx dx, \quad n=0, 1, 2, \dots$$

$$b_n = \frac{1}{\pi} \int_0^{2\pi} f(x) \sin nx dx, \quad n=1, 2, 3, \dots$$

I. A. (Continued)

Hawes ②

4. The complex and real representation coefficients are related by

$$c_0 = \frac{a_0}{2} \quad c_n = \frac{a_n - ib_n}{2}, \quad c_{-n} = \frac{a_n + ib_n}{2}, \quad n > 0$$

B. Sturm-Liouville Theory

1. Recall that a Hermitian operator for a problem with periodic boundary conditions is a Sturm-Liouville Problem.

Therefore: a) Eigenfunctions are real,

b) Complete basis of orthogonal eigenfunctions.

2. For this problem

a. ODE: $-y''(x) = \lambda y(x)$

b. $\langle f|g \rangle = \int_0^{2\pi} f^*(x)g(x) dx$

c. Boundary Conditions: $y(0) = y(2\pi), \quad y'(0) = y'(2\pi)$

d. Eigenfunctions: $\phi_n(x) = \frac{1}{\sqrt{2\pi}} e^{inx} \quad n = \dots, -1, 0, 1, \dots$

or $\phi_0 = \frac{1}{\sqrt{2\pi}}, \quad \phi_n = \frac{\cos nx}{\sqrt{\pi}}, \quad \phi_{-n} = \frac{\sin nx}{\sqrt{\pi}} \quad \text{for } n = 1, 2, 3, \dots$

e. Eigenvalues: n^2

3. Normalization of eigenfunctions: Note!

a) For complex representation, $\langle e^{inx} | e^{inx} \rangle = \int_0^{2\pi} e^{-inx} e^{inx} dx = 2\pi$

b) For real representation, $\int_0^{2\pi} \sin^2(nx) dx = \int_0^{2\pi} \cos^2(nx) dx = \pi$

C. Representing Discontinuous Functions

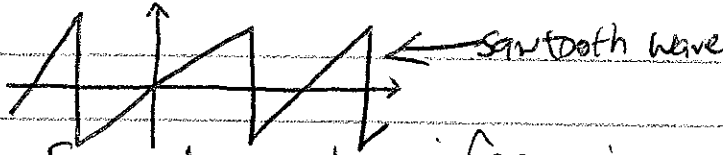
1. Power series cannot represent discontinuous functions

a. Expansion about a point, with radius of convergence up to nearest singularity

I. C. (Continued)

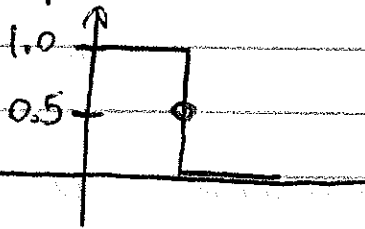
Hanes ③

2. Fourier Series can describe functions with finite discontinuities.



a. Expansion employs integration over entire interval, so it can describe "nonpathological" singularities

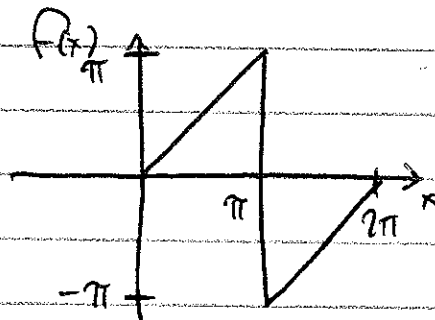
3. For a piecewise regular function $f(x)$ (also known as satisfying Dirichlet conditions), Fourier series at a discontinuous point is limit of left and right approaches



$$f_{\text{Fourier Series}}(x_0) = \lim_{\epsilon \rightarrow 0} \left[\frac{f(x_0 + \epsilon) + f(x_0 - \epsilon)}{2} \right]$$

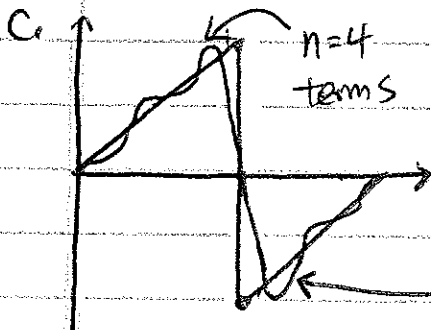
4. Ex: Sawtooth Wave

$$a. f(x) = \begin{cases} x & 0 \leq x < \pi \\ x - 2\pi & \pi < x \leq 2\pi \end{cases}$$



b. Using real representation, $a_n = 0$, $b_n = \frac{1}{\pi} \int_0^{2\pi} f(x) \sin nx \, dx$ leads to

$$f(x) = 2 \left[\sin x - \frac{\sin 2x}{2} + \frac{\sin 3x}{3} - \dots + (-1)^{n+1} \frac{\sin nx}{n} + \dots \right]$$



d. Comments:

- i. Accuracy of $f(x)$ increases with number of terms
- ii. All curves pass through $f(\pi) = 0$
- iii. Overshoot near discontinuity \Rightarrow Gibbs phenomenon!

I. (Continued)

Homes (4)

D. General Periodic Interval $[-L, L]$

1. By taking interval ~~over~~ $[-\pi, \pi]$ instead of $[0, 2\pi]$ and substituting $y = x \frac{L}{\pi}$, we may obtain

$$f(x) = \sum_{n=-\infty}^{\infty} c_n e^{i \frac{n\pi x}{L}}$$

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos \frac{n\pi x}{L} + \sum_{n=1}^{\infty} b_n \sin \frac{n\pi x}{L}$$

$$c_n = \frac{1}{2L} \int_{-L}^L f(x) e^{-i \frac{n\pi x}{L}} dx$$

$$a_n = \frac{1}{L} \int_{-L}^L f(x) \cos \frac{n\pi x}{L} dx$$

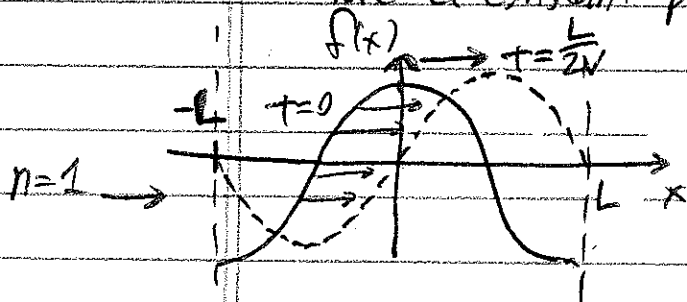
$$b_n = \frac{1}{L} \int_{-L}^L f(x) \sin \frac{n\pi x}{L} dx$$

E. Traveling Waves

1. A waveform may be traveling with phase velocity v , so we obtain terms

$$\sim \cos \left[\frac{n\pi}{L} (x - vt) \right]$$

where a constant phase of wave moves at $x - vt = \text{constant}$, or $(v = \frac{x}{t})$



2. For a particular mode n ,

a. Wavelength $\lambda_n = \frac{2L}{n}$

b. Period $T_n = \frac{\lambda_n}{v} = \frac{2L}{nv}$

c. Linear Frequency: $f = \frac{1}{T} = \frac{nv}{2L}$, Angular Frequency $\omega = \frac{2\pi}{T} = \frac{\pi n v}{L}$

d. Fundamental Frequency: $n=1$

Harmonic Frequencies: $n > 1$

I. (Continued)

Howes (5)

F. Linearity

1. Fourier Analysis is most useful for Linear Differential Equations

2a. For a linear problem, the general solution is a linear superposition of each Fourier component (each mode n)

b. Each Fourier component evolves independently of the others!

c. The solution for general mode n is applicable to all modes.

3a. For nonlinear problems, the nonlinear terms lead to coupling among different Fourier modes.

b. Numerically, efficient algorithms, such as the Fast Fourier Transform [Cooley & Tukey, (1965)], enable nonlinear evolution of Fourier components in time (pseudospectral).

c. Pseudospectral algorithms evaluate linear terms in Fourier space and nonlinear terms in real space.

G. Symmetry

1. For functions $f(x)$ that have definite parity, the real representation of Fourier series ($\sin nx$ & $\cos nx$) leads to simplifications

a. For $f(x)$ even,
$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos nx \quad (b_n = 0)$$

b. For $f(x)$ odd,
$$f(x) = \sum_{n=1}^{\infty} b_n \sin nx \quad (a_n = 0)$$

2. Sometimes these are denoted Fourier cosine & Fourier sine series.

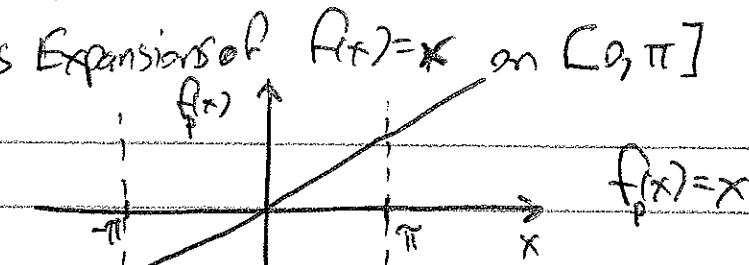
Z.G. (Continued)

Hines (6)

3. Ex: Various Expansions of $f(x) = x$ on $[0, \pi]$

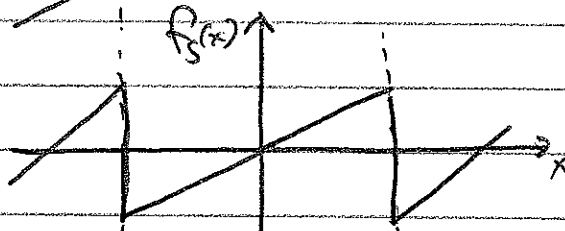
a. Power Series:

Non-periodic



b. Fourier Sine Series:

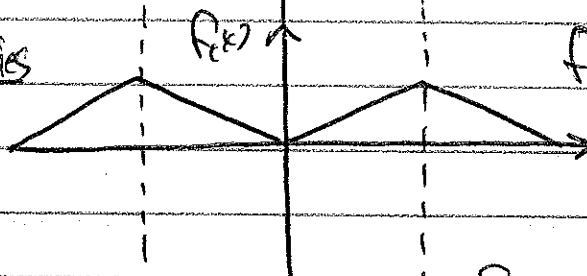
Odd parity



$$f_S(x) = 2 \sum_{n=1}^{\infty} (-1)^{n+1} \frac{\sin nx}{n}$$

c. Fourier Cosine Series:

Even parity



$$f_C(x) = \frac{\pi}{2} - \sum_{n=0}^{\infty} \frac{4}{\pi} \frac{\cos(2n+1)x}{(2n+1)^2}$$

Different outside of $[0, \pi]$ interval

Same form on $[0, \pi]$ for all three functions

H. Integration and Differentiation of Fourier Series

1. Consider term-by-term integration of $f(x) = \sum_{n=-\infty}^{\infty} c_n e^{inx}$

a. $\int f(x) dx = \sum_{n=-\infty}^{\infty} \frac{c_n}{in} e^{inx} + \text{constant}$

b. Additional factor of n^{-1} in each term \Rightarrow more rapid convergence!

2. Thus, A convergent Fourier Series may always be integrated term by term

3. Consider term-by-term differentiation of Sawtooth wave (a convergent Fourier Series)

a. $f(x) = x = 2 \sum_{n=1}^{\infty} (-1)^{n+1} \frac{\sin nx}{n}$ on $-\pi \leq x \leq \pi$

b. Differentiating: $f'(x) = 1 = 2 \sum_{n=1}^{\infty} (-1)^{n+1} \cos(nx) \leftarrow \text{Not convergent!}$

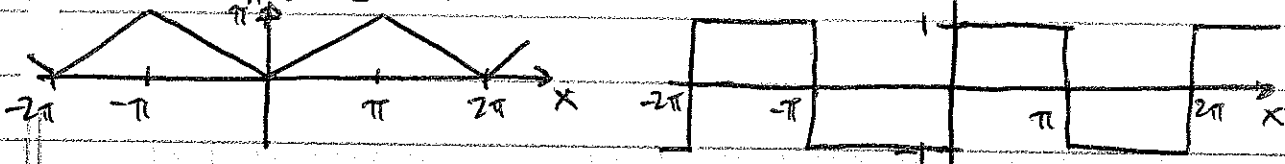
4a. Term-by-term differentiation of a convergent Fourier Series may not lead to a convergent series for the derivative

b. Always check the resulting series for convergence!

5. Ex: Derivative of triangular wave \Rightarrow Square wave

a. $f(x) = \frac{\pi}{2} - \frac{4}{\pi} \sum_{n=0}^{\infty} \frac{\cos(2n+1)x}{(2n+1)^2}$

b. $f'(x) = \frac{4}{\pi} \sum_{n=0}^{\infty} \frac{\sin(2n+1)x}{(2n+1)}$



6. Generally:

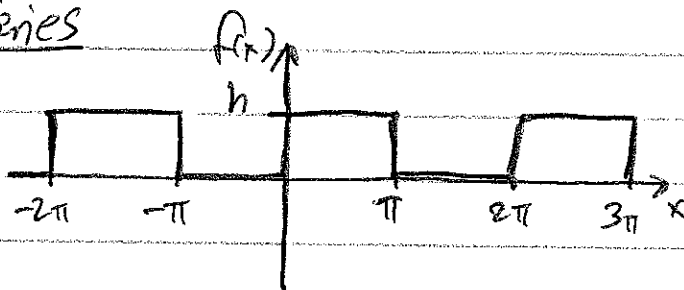
a. Differentiation reduces rate of convergence \Rightarrow may yield divergent series

b. Term-by-term differentiation usually permitted for a uniformly convergent series

II. Applications of Fourier Series

A. Example: Square Wave

1. $f(x) = \begin{cases} 0 & -\pi < x < 0 \\ h & 0 < x < \pi \end{cases}$



2. NOTE: $f(x) - \frac{h}{2}$ is odd, so definite parity suggests real Fourier representation is appropriate choice (even terms will be zero, except a_0)

3. $a_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \cos nx \, dx$

a. For $n=0$, $a_0 = \frac{1}{\pi} \int_{-\pi}^0 (0) \, dx + \frac{1}{\pi} \int_0^{\pi} h \, dx = \frac{h(\pi-0)}{\pi} = h$

b. For $n > 0$, $a_n = \frac{1}{\pi} \int_{-\pi}^0 (0) \cos nx \, dx + \frac{1}{\pi} \int_0^{\pi} h \cos nx \, dx = \frac{h}{\pi} \left[\frac{\sin nx}{n} \right]_0^{\pi} = \frac{h}{\pi} \left(\frac{\sin n\pi - \sin 0}{n} \right) = 0!$

IIA (Continued)

Howes (8)

$$4. b_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin(nx) dx$$

$$a. b_n = \frac{1}{\pi} \int_{-\pi}^0 (0) \sin(nx) dx + \frac{1}{\pi} \int_0^{\pi} h \sin(nx) dx = \frac{h}{\pi} \left[\frac{-\cos nx}{n} \right]_0^{\pi} = \frac{h}{\pi} \left[\frac{1 - \cos n\pi}{n} \right]$$

$$b. \text{NOTE: } \cos n\pi = (-1)^n, \text{ so } b_n = \frac{h}{\pi} \left[\frac{1 - (-1)^n}{n} \right] = \begin{cases} 0, & n \text{ even} \\ \frac{2h}{n\pi}, & n \text{ odd} \end{cases}$$

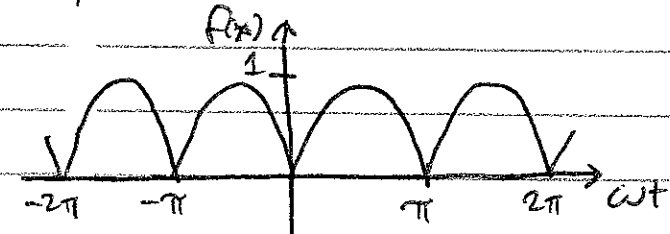
$$5. f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos nx + \sum_{n=1}^{\infty} b_n \sin nx$$

$$f(x) = \frac{h}{2} + \frac{2h}{\pi} \sum_{n=0}^{\infty} \frac{\sin(2n+1)x}{(2n+1)}$$

6. IMPORTANT NOTE:

- Coefficients decrease only as $n^{-1} \Rightarrow$ Convergence is slow!
- Thus, high frequency harmonics do not drop off very fast
 \Rightarrow Signal contains significant contribution from high frequencies!
- For numerical simulations using Fourier representation, resolution of discontinuities requires many modes \Rightarrow not well suited!

B. Ex: Full Wave Rectifier

$$1. f(t) = \begin{cases} \sin \omega t & 0 < \omega t < \pi \\ -\sin \omega t & -\pi < \omega t < 0 \end{cases}$$


2. Even parity, so $b_n = 0$ for all n !

$$3. a. a_0 = \frac{1}{\pi} \int_{-\pi}^0 (-\sin \omega t) d(\omega t) + \frac{1}{\pi} \int_0^{\pi} \sin \omega t d(\omega t) = \frac{2}{\pi} \int_0^{\pi} \sin \omega t d(\omega t) = \frac{4}{\pi}$$

$$b. a_n = \frac{2}{\pi} \int_0^{\pi} \sin \omega t \cos(n\omega t) d(\omega t) = \begin{cases} -\frac{2}{\pi} \frac{2}{n^2-1}, & n \text{ even} \\ 0, & n \text{ odd} \end{cases}$$

$$4. f(t) = \frac{2}{\pi} - \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{\cos(2n\omega t)}{(2n)^2-1} \leftarrow \text{coefficients decrease as } n^{-2}!$$

- C. Generally,
- If $f(x)$ has discontinuities, $c_n \propto O\left(\frac{1}{n}\right) \Rightarrow$ Slow convergence!
 - If $f(x)$ continuous [but $f'(x)$ may be discontinuous], $c_n \propto O\left(\frac{1}{n^2}\right) \Rightarrow$ Faster convergence!