

Final Exam Review:I. Magnetohydrodynamics (MHD)A. MHD Approximation

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|-----------------------|--------------------------------------|--|
| 1. Strong collisions: | $\lambda_m \ll L, \tau \gg \nu_{ei}$ | } Quasi-neutrality
$\sum_s n_s q_s = 0$ |
| 2. Non-relativistic: | $v_0^2/c^2 \ll 1$ | |
| 3. Magnetized: | $r_{Li} \ll L$ | |

B. Derivation from Two-fluid Equations

- Drop $\underline{\underline{\Pi}}$, viscous terms from $\underline{\underline{P}}$ tensor (Collisional)
- Zonarpic pressure (Collisional)
- Drop Displacement Current (Non-relativistic)
- Drop Hall term in Ohm's Law (Magnetized $r_{Li} \gg L$)
- Adiabatic Equation of State

C. MHD Equations

- Continuity $\frac{\partial \rho}{\partial t} + \underline{U} \cdot \nabla \rho = -\rho \nabla \cdot \underline{U}$
- Momentum $\rho \frac{\partial \underline{U}}{\partial t} + \rho \underline{U} \cdot \nabla \underline{U} = -\nabla \left(p + \frac{B^2}{2\mu_0} \right) + \frac{(\underline{B} \cdot \nabla) \underline{B}}{\mu_0}$
- Induction $\frac{\partial \underline{B}}{\partial t} = \nabla \times (\underline{U} \times \underline{B}) + \frac{c}{\mu_0} \nabla^2 \underline{B}$
- Adiabatic Equation of State $\frac{\partial p}{\partial t} + \underline{U} \cdot \nabla p = -\gamma p \nabla \cdot \underline{U}$

D. Frozen-In Flux

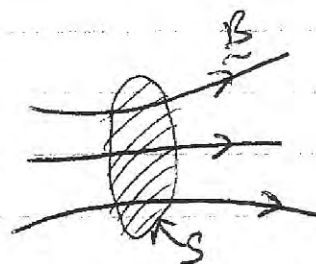
1. Magnetic Reynolds Number: $Re_M = \frac{\mu_0 L V_0}{\eta}$

2. Limits: a) $Re_M \gg 1 \Rightarrow$ Ideal MHD

b) $Re_M \ll 1 \Rightarrow$ Diffusive

3. Frozen-In Flux Theorem: (Ideal Limit)

a. Magnetic Flux $\Phi_B = \int_S \vec{B} \cdot d\vec{A}$



b. Magnetic field lines are frozen to the flow

4. Clebsch Coordinates:

a. $\vec{B} = \nabla\alpha \times \nabla\beta$

b. $\alpha(x) = \text{const}$
 $\beta(x) = \text{const}$ } Defines field lines $\vec{B}(x)$

5. Applications of Frozen-In Flux: $\Phi_B = \text{constant}$ as dA changes!

E. MHD Linear Dispersion Relation

1. Know the steps to compute the linear dispersion relation!

a. Linearize the equations: $\rho = \rho_0 + \epsilon \rho_1(x)$, $\underline{U} = \underline{U}_0 + \epsilon \underline{U}_1(x)$
 ρ_0 uniform, steady $\epsilon \rho_1(x)$ ecc.

b. Fourier transform (plane-wave solution)

$$e^{i(\underline{k}\cdot\underline{x} - \omega t)} \Rightarrow \nabla \Rightarrow i\underline{k}$$

$$\frac{\partial}{\partial t} \Rightarrow -i\omega$$

I. Continued)

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c. Eliminate $\rho_1, \underline{B}_1, p_1$ in terms of \underline{U}_1

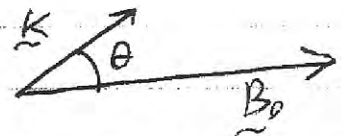
d. Collect to obtain matrix equation

$$\underline{M} \cdot \underline{U} = 0$$

e. Set Determinant $|\underline{M}| = 0 \Rightarrow$ yields $D(\omega, \underline{k}) = 0$

f. Solutions for ω (in terms of \underline{k} and other parameters) yield the linear wave modes!

F. MHD Wave Modes



1. $\underline{k} = k_{\perp} \hat{x} + k_{\parallel} \hat{z} = k \sin \theta \hat{x} + k \cos \theta \hat{z}$, $\underline{B} = B_0 \hat{z}$

2. General Dispersion Relation

$$(\omega^2 - k_{\parallel}^2 v_A^2) \left[\omega^4 - \omega^2 k^2 (c_s^2 + v_A^2) + k_{\parallel}^2 k^2 c_s^2 v_A^2 \right] = 0$$

Alfvén Modes,
incompressible

Fast and Slow Magnetosonic Waves,
compressible

$$\omega = \pm k_{\parallel} v_A$$

$$\frac{\omega^2}{k^2} = \frac{1}{2} (c_s^2 + v_A^2) \pm \frac{1}{2} \sqrt{(c_s^2 + v_A^2)^2 - 4 c_s^2 v_A^2 \cos^2 \theta}$$

3. Sound speed

$$c_s^2 = \frac{\gamma p_0}{\rho_0}$$

Alfvén Speed

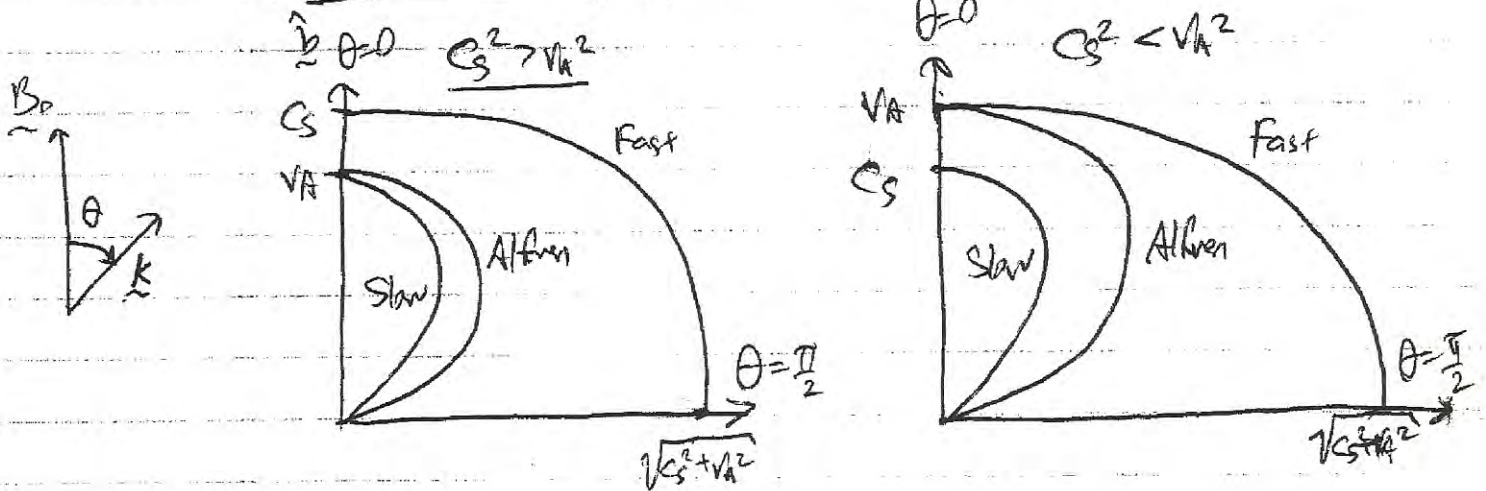
$$v_A^2 = \frac{B_0^2}{\mu_0 \rho_0}$$

4. a. Fast: Thermal & Magnetic Pressure add
Slow: Thermal & Magnetic Pressure oppose

Z. F. Continued)

Hawes (4)

5. Polar Plots: $v_p = \frac{\omega}{k} = \text{phase speed}$



G. Conserved Energy

$$1. E = \int d^3x \left(\frac{1}{2} \rho U^2 + \frac{p}{\gamma - 1} + \frac{B^2}{2\mu_0} \right)$$

H. Eigenfunctions of MHD Waves

$$1. \begin{pmatrix} (-) & 0 & (-) \\ 0 & (-) & 0 \\ (-) & 0 & (-) \end{pmatrix} \begin{pmatrix} U_x \\ U_y \\ U_z \end{pmatrix} = 0$$

2. Take $U_y = 0$, $U_x = U_0$

3. Solve for all other quantities in terms of U_x using the linearized, Fourier-transformed equations

$\Rightarrow p_1, \underline{u}_1, \underline{B}_1, p_1$ in terms of $U_x = U_0!$

4. Plug in ω for particular wave mode.

II. MHD Equilibrium

1. Force-Balanced Equilibrium: $\nabla p = \underline{j} \times \underline{B}$

2. Pfaff's Thm: Torus is simple surface
such that $\underline{B} \cdot \nabla p = 0$.

3. Force-Free Equilibrium: $\underline{j} \times \underline{B} = 0$

$$a \Rightarrow \nabla \times \underline{B} = \alpha \underline{B} \quad \Rightarrow \underline{B} \cdot \nabla \alpha = 0$$

surfaces of constant α

4. Cylindrical Equilibria: $\nabla \left(p + \frac{B^2}{2\mu_0} \right) = \frac{(\underline{B} \cdot \nabla) \underline{B}}{\mu_0}$

ρ_r ρ_B μ_0 Mag. Energy

a. $\underline{B} = B_\phi(r) \hat{\phi} + B_z(r) \hat{z}$ (Ex. Flux rope, RFP)

b. $\frac{d}{dr} \left(p + \frac{B_\phi^2}{2\mu_0} + \frac{B_z^2}{2\mu_0} \right) = -\frac{B_\phi^2}{\mu_0 r}$

c. $B_z = 0 \Rightarrow z$ -pinch

$B_\phi = 0 \Rightarrow$ Theta-pinch

5. Toroidal Equilibria: (Nac on Eqm)

a. Magnetic Flux Coordinates $(\psi(r, z), F(r))$

b. Grad-Shafranov Equation

\Rightarrow Magnetostatic Equilibria

III. Multiple Timescale Methods

A. General Approach

1. Two timescales: $t, \tau = \epsilon^2 t$
2. $\frac{d}{dt} = \frac{\partial}{\partial t} + \epsilon^2 \frac{\partial}{\partial \tau}$
3. Expand variables: $X = \cancel{X_0} + \epsilon X_1 + \epsilon^2 X_2 + \dots$
4. Substitute into equation and multiply out to get all terms
5. Collect terms order by order to obtain
 - $O(1)$
 - $O(\epsilon)$ equations
 - $O(\epsilon^2)$
6. Solve Equations order by order for x_1, x_2, \dots
 \Rightarrow set $\epsilon = 1!$
7. NOTE! Every problem is different. I will try to walk you through any problems I give you (Krook, Homes)

III. (Continued)

Hanes 7

B. Polarization Drift

1. Slowly varying electric field, $\underline{E}(\tau)$ with constant $\underline{B} = B_0 \hat{z}$.

$$2. \underline{v} = v_L (\cos(\omega t + \phi) \hat{x} - \sin(\omega t + \phi) \hat{y}) + \underbrace{\frac{\underline{E}(\tau) \times \underline{B}}{B_0^2}}_{\substack{\mathcal{O}(E) \\ \underline{E} \times \underline{B} \text{ drift}}} + \underbrace{\frac{1}{\omega_c B_0} \frac{d\underline{E}}{d\tau}}_{\substack{\mathcal{O}(E^2) \\ \text{Polarization} \\ \text{Drift, } \underline{v}_p}}$$

$$3. \underline{v}_p = \frac{1}{\omega_c B} \frac{d\underline{E}}{d\tau} = \frac{m}{2B^2} \frac{d\underline{E}}{d\tau}$$

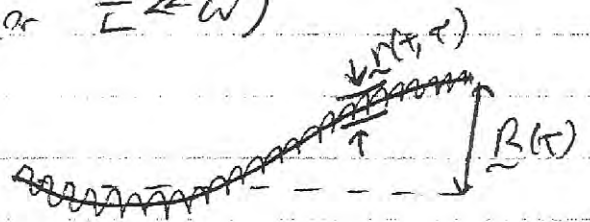
C. Ponderomotive Force

1. Particle motion in rapidly varying electromagnetic field

$$\underline{E}(\underline{x}, t, \tau) = \underline{E}_0(\underline{x}, \tau) \cos(\omega t - \underline{k} \cdot \underline{x})$$

2. Ordering: $\epsilon = \frac{v}{L\omega} \ll 1$ (or $\frac{v}{L} \ll \omega$)

3. $\underline{x} = \underbrace{\underline{R}(\tau)}_{\text{oscillation center}} + \epsilon \underline{\rho}(t, \tau)$



4. $\underline{F}_{\text{pond}} = m \frac{d\underline{u}}{d\tau} = \frac{-q^2}{4\pi\epsilon_0 \omega^2} \nabla |\underline{E}_0|^2$

5. $\underline{F}_{\text{pond}} = -\nabla \Phi_{\text{pond}} \quad \Phi_{\text{pond}} = \frac{q^2}{4\pi\epsilon_0 \omega^2} |\underline{E}_0|^2$

NOTE:
Drop $\cos(\omega t - \underline{k} \cdot \underline{x})$ in this part!

6. Energy Picture: $E_{\text{osc}} = \frac{1}{2} m u^2 + \Phi_{\text{pond}} = \text{constant!}$

IV MHD Stability

A. Will not be tested on the Final Exam

V. Kinetic Theory: Linear Waves and Landau Damping

A. Equations:

1. Boltzmann Eq: $\frac{\partial f_s}{\partial t} + \underline{v} \cdot \nabla f_s + \frac{q_s}{m_s} (\underline{E} + \underline{v} \times \underline{B}) \cdot \frac{\partial f_s}{\partial \underline{v}} = \left(\frac{\partial f_s}{\partial t} \right)_{\text{coll}}$

a. Vlasov Equation when $\left(\frac{\partial f_s}{\partial t} \right)_{\text{coll}} = 0$.

b. $f_s(x, y, t)$: 6D (3D-3V) Distribution Function

2. Maxwell's Equations:

$$\nabla \cdot \underline{E} = \frac{\rho_2}{\epsilon_0}$$

$$\rho_2 = \sum_s \int d^3v q_s f_s$$

$$\nabla \times \underline{E} = -\frac{\partial \underline{B}}{\partial t}$$

$$\nabla \times \underline{B} = \mu_0 \underline{j} + \mu_0 \epsilon_0 \frac{\partial \underline{E}}{\partial t}$$

$$\underline{j} = \sum_s \int d^3v q_s \underline{v} f_s$$

$$\nabla \cdot \underline{B} = 0$$

B. Equilibrium and Moments:

1. Maxwellian Equilibrium: $f_{sm}(x, y, t) = \frac{n_s(x, y, t)}{\pi^{3/2} v_{Ts}^{3/2}} e^{-\frac{m_s |v - U_s(x, y)|^2}{2 T_s(x, y)}}$

a. $v_{Ts}^2 = \frac{2 T_s}{m_s}$

b. Uniform: $f_{sm}(v) = \frac{n_s}{\pi^{3/2} v_{Ts}^{3/2}} e^{-\frac{v^2}{v_{Ts}^2}}$

VB (Continued)

Hand 9

2. Moments: a. $n_{s0} = \int d^3v f_s(v)$

b. $\frac{3}{2} n_{s0} T_{s0} = \int d^3v \frac{1}{2} m_s v^2 f_{sm}(v)$

3. Reduced Distribution Function:

a. $F_s(v_z) \equiv \frac{1}{n_{s0}} \int_{-\infty}^{\infty} dv_x \int_{-\infty}^{\infty} dv_y f_s(v)$

b. $\int_{-\infty}^{\infty} dv_z F_s(v_z) = 1$

C. Electrostatic Waves in an Unmagnetized Plasma

1. $\underline{E} = -\nabla\phi$

2. Linearized: a. $\frac{\partial f_{s1}}{\partial t} + \underline{v} \cdot \nabla f_{s1} - \frac{q_s}{m_s} \nabla\phi \cdot \frac{\partial f_{s0}}{\partial \underline{v}} = 0$

b. $-\nabla^2 \phi_1 = \frac{1}{\epsilon_0} \sum_s q_s \int d^3v f_{s1}$

3. Fourier Transform Solution: $\frac{\partial}{\partial t} \rightarrow -i\omega$, $\nabla \rightarrow i\vec{k}$

a. $D(\vec{k}, \omega) = 1 - \sum_s \frac{q_s n_{s0}}{\epsilon_0 k^2} \int_{-\infty}^{\infty} dv_z \frac{\partial f_{s0} / \partial v_z}{(v_z - \frac{\omega}{k})} = 0$

b. Only valid when $f_{s0}(v_z = \frac{\omega}{k}) = 0$ and $\left. \frac{\partial f_{s0}}{\partial v_z} \right|_{v_z = \frac{\omega}{k}} = 0$

c. Integrate by parts: $D(\vec{k}, \omega) = 1 - \sum_s \frac{q_s n_{s0}}{\epsilon_0 k^2} \int_{-\infty}^{\infty} dv_z \frac{f_{s0}}{(v_z - \frac{\omega}{k})^2} = 0$

V. (Continued)

Haves (11)

F. Laplace-Fourier Solution for Electromagnetic Plasma Waves

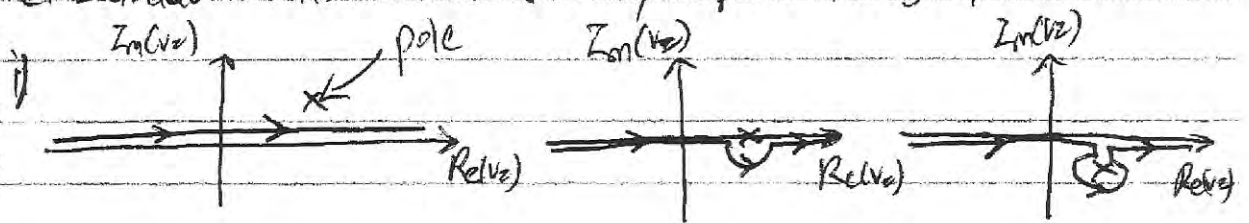
1. Solution: $\tilde{\Phi}_i(k, p) = \frac{N(k, p)}{D(k, p)}$

a. $N(k, p) = -i \sum_s \frac{q_s n_0}{\epsilon_0 k^3} \int_{-\infty}^{\infty} dv_z \frac{F_s(v_z)}{v_z - i\frac{p}{k}}$

b. Dispersion Relation: $D(k, p) = 0$

$$D(k, p) = 1 - \sum_s \frac{q_s^2 n_0^2}{k^2} \int_{-\infty}^{\infty} dv_z \frac{\partial F_s / \partial v_z}{v_z - i\frac{p}{k}} = 0$$

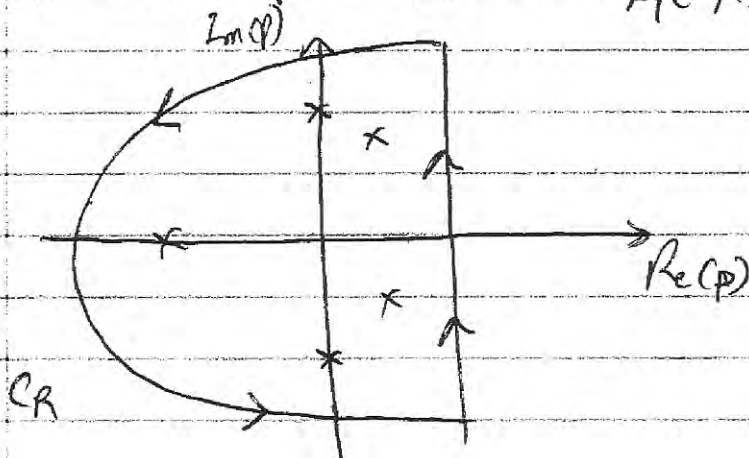
c. Landau Contour for velocity space integral:



ii) Contour always goes under the pole,

iii) Achieves analytic continuation of $D(k, p)$ to $\text{Re}(p) < 0$!

d. Inverse Laplace Transform: $\hat{\Phi}_i(k, t) = \frac{1}{2\pi i} \int_{\sigma-i\infty}^{\sigma+i\infty} \tilde{\Phi}_i(k, p) e^{pt} dp$



e. Close Contour (C_R)

f. Evaluate using Residue Theorem

V.E. (Continued)

Hines 13

2. Know how to solve for $D(\underline{k}, \rho) = 0$
for a given choice of $F_0(v_z)$!

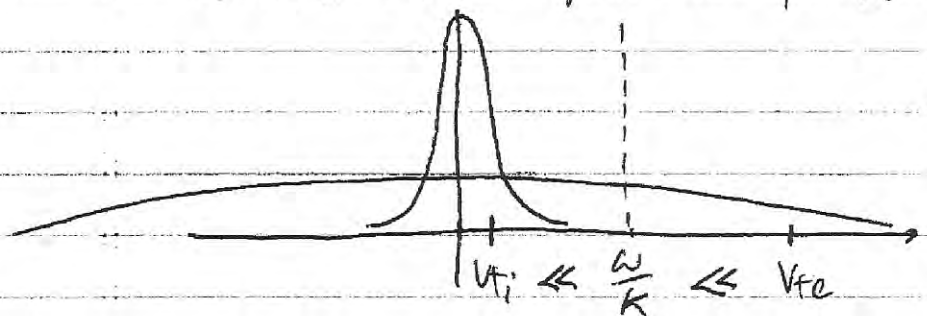
a. Examples: Cauchy Distribution: $F_{0c}(v_z) = \frac{C}{\pi} \frac{1}{C^2 + v_z^2}$

b. For Maxwellian Distribution \Rightarrow Plasma Dispersion Function!
(see below)

F. Weak Growth Rate Approximation

1. Can be used anytime $|\delta| \ll |\omega|$

a. When resonance phase velocity $\frac{\omega}{k}$ is not near v_{te} !



i) $\frac{\omega}{k} \gg v_{ti} \Rightarrow$
weak ion damping!

ii) $\frac{\omega}{k} \ll v_{te} \Rightarrow$
weak electron damping!

2. a. Solve for complex function $D(\underline{k}, \rho) = 0$

b. Separate real and
imaginary parts: $D(\underline{k}, -i\omega) = D_r(\underline{k}, -i\omega) + i D_i(\underline{k}, -i\omega)$

$$D(\underline{k}, -i\omega) = 0$$

$\omega = \omega_r + i\delta$

c. Set $\delta = 0 \rightarrow \omega = \omega_r + i\delta \rightarrow \omega_r!$

V. F. (Continued)

Haves (13)

3. Weak Growth Rate Approximation:

$$\begin{array}{l}
 \text{a. } D_r(k, -i\omega_r) = 0 \Rightarrow \text{solve for } \omega_r! \\
 \text{b. } \gamma = \frac{-D_i(k, -i\omega_r)}{\partial D_r(k, -i\omega_r) / \partial \omega_r} \Rightarrow \text{solve for } \gamma!
 \end{array}$$

4. Application to Electrostatic Waves:

a. Plemelj Relations: i) $\lim_{\epsilon \rightarrow 0} \int_{-\infty}^{\infty} dx \frac{f(x)}{x - (x_0 \pm i\epsilon)} = P \int_{-\infty}^{\infty} dx \frac{f(x)}{x - x_0} \pm i\pi f(x_0)$

ii) Principle Value: $P \int_{-\infty}^{\infty} dx \frac{g(x)}{x - x_0} = \lim_{\delta \rightarrow 0} \left[\int_{-\infty}^{x_0 - \delta} dx \frac{g(x)}{x - x_0} + \int_{x_0 + \delta}^{\infty} dx \frac{g(x)}{x - x_0} \right]$

b. For equilibrium $f_{s0}(v_z)$,

$$\gamma = \pi \frac{k}{|k|} \frac{1}{\partial D_r(k, \omega_r) / \partial \omega_r} \sum_s \frac{\omega_{ps}^2}{k^2} \left. \frac{\partial f_{s0}}{\partial v_z} \right|_{v_z = \frac{\omega}{k}}$$

G. Plasma Dispersion Function: For Maxwellian equilibrium,

1. $Z(\zeta_s) = \int_{-\infty}^{\infty} \frac{dz}{\pi^{1/2}} \frac{e^{-z^2}}{z - \zeta_s}$ $\zeta_s = \frac{i\omega}{kv_{Ts}} = \frac{\omega}{kv_{Ts}} - i \frac{\gamma}{kv_{Ts}}$

2. Limits:

a. $|\zeta| \gg 1$: $Z(\zeta) = i\pi^{1/2} e^{-\zeta^2} - \frac{1}{\zeta} - \frac{1}{2\zeta^3} - \frac{3}{4\zeta^5} - \dots$

b. $|\zeta| \ll 1$: $Z(\zeta) = i\pi^{1/2} e^{-\zeta^2} - 2\zeta + 4\frac{\zeta^3}{3} - 8\frac{\zeta^5}{15} + \dots$

VI. Kinetic Stability

Hand 14

A. Cold Beam Instabilities

1. Cold Beam Distribution: $f_{0s}(v) = \sum_j N_{js} \delta(v - v_j) \delta(v_z - v_{js})$

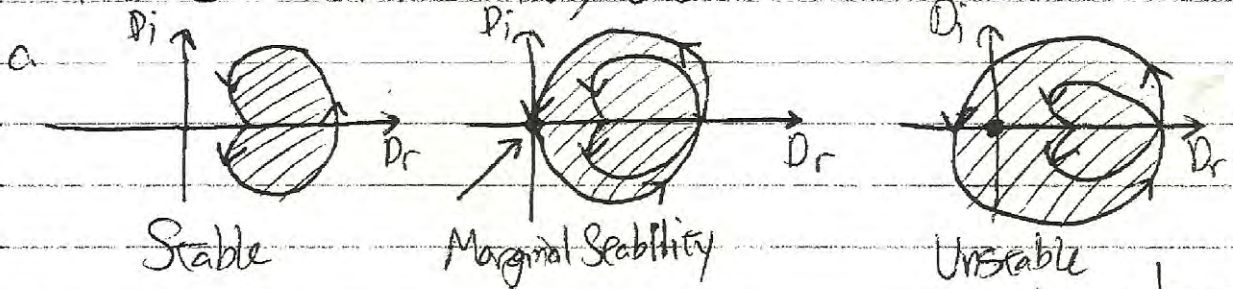
2. Substitute into Dispersion Relation and solve for complex ω !

B. Kinetic Stability

1. Gardner's Thm: One-humped distribution is always stable!

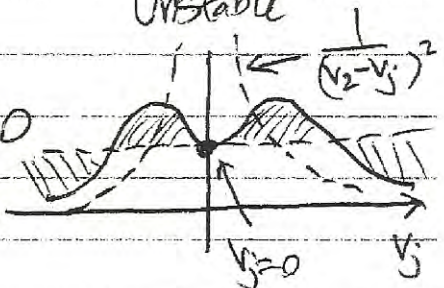
2. Nyquist Criteria: Map Complex- p to Complex- D space.

b. If $D=0$ within contour, unstable



3. Pearse Condition: $\int_{-\infty}^{\infty} dv_z \frac{F_0(v_z) - F_0(v_j)}{(v_z - v_j)^2} > 0$

\uparrow where $\frac{dF_0}{dv_z} = 0$!



VII. Quasilinear Theory: Not on Exam

VIII. Plasma Sheaths & Langmuir Probes: Not on Exam