

I. Why Do We Study Plasma Physics?

A. The Fundamental Physics of Plasmas

1. We would like to understand the fundamental behavior of naturally occurring plasmas and their effect on their environment.
2. Examples: Astrophysical plasmas (intracluster medium, interstellar medium, accretion disks, stellar interiors), solar corona and solar wind, Earth's magnetosphere and ionosphere.

B. Applications

1. Exploit the properties of plasmas to achieve desirable results.
2. Examples: a. FUSION: cheap energy!
b. Industrial applications: plasma processing of semiconductors, plasma sterilization, efficient lighting, electronics

II. Aim of this Lecture

A. The Basics of Plasma Physics

1. In 29:194 & 29:293, we have learned all of the fundamentals of the study of plasma physics.
2. We should now feel confident to attack any problem using some combination of these fundamental concepts and approaches.

B. The Framework

1. Today, I will attempt to provide an overview of plasma physics highlighting:
 - a. Relationships between the topics we have studied
 - b. The context placing what we have learned in terms of what we have not studied.
⇒ "the known unknowns."

III. Characteristic Scales in a Plasma

A. Summary

Length	Time/Frequency	Dimensionless
Particle spacing, $n_0^{-1/3}$	Plasma Frequency, ω_p	Plasma Parameter, N_0
Debye Length, λ_D	Cyclotron Frequency, ω_c	Plasma Beta, β
Larmor Radius, r_L	Collision Frequency, ν	Magnetization v_L/L
Mean Free Path, λ_m	Obsecurian "Frequency", $\frac{1}{\tau}$	Collisionality $\nu\tau$
System size, L		

B. Choosing an appropriate plasma description

1. Use the characteristic scales of the system of interest to determine the system of equations to use
 2. Generally, we want to use the most simple system that contains the relevant physical effects
 - a. Separation of timescales often allows us to ignore fast, small-scale physics if we are interested in larger times and larger scales.
- $\frac{1}{2}$ Fast $\mathbf{E} \times \mathbf{B}$ drift: The details of the fast Larmor motion are often not of interest — rather, we want to know the net drift of the particles.
2. Fast charge density oscillations at the plasma frequency have negligible effect on the slow, large scale MHD flow. \Rightarrow These fast motions average out.

IV. How do we choose an appropriate description of the plasma?

We'll begin with the most simple description, then generalize as necessary to the most complicated description.

0. Look at what others have done \rightarrow but be wary!

A. Single Particle Motion:

1. Motion of Charged particles in prescribed fields $\underline{E}(x, t)$ & $\underline{B}(x, t)$.

a. This is only useful if you can guess \underline{E} & \underline{B} well, and if the effect of the charge density ρ and current density \mathbf{j} is small.

b. Gives \underline{x}_s and \underline{v}_s using the Lorentz Force Law, $m_s \frac{d\underline{v}_s}{dt} = q_s (\underline{E} + \underline{v}_s \times \underline{B})$

c. Can be a useful description of very low density plasma behavior in strong external fields. For example, laser plasma interactions.

d. Otherwise, useful for building intuition about plasma behavior.

2. Guiding Center Approximation

a. Useful for fields that vary slowly in time and smoothly in space.

$$\text{Slow: } \frac{1}{|B|} \left| \frac{dB}{dt} \right| \ll \omega_{ci} \qquad \frac{1}{|E|} \left| \frac{dE}{dt} \right| \ll \omega_{ci}$$

$$\text{Smooth: } \frac{|\nabla B|}{|B|} \ll \frac{1}{r_{Li}} \qquad \frac{|\nabla E|}{|E|} \ll \frac{1}{r_{Li}}$$

b. Adiabatic Invariants can be useful to describe the motion of particle in prescribed fields.

Ex: The motion of energetic particles in Earth's magnetosphere

i. First Adiabatic Invariant: Magnetic Moment μ (Chamber Motion)

ii. Second Adiabatic Invariant: Parallel Bounce Motion, J_z

iii. Third Adiabatic Invariant: Azimuthal Drift Motion

3. Inconsistent Model: Fields are not consistently generated.

IV. (Continued)

B. Magnetohydrodynamics (Single Fluid Theory), MHD

1. This is the most simple, consistent model of plasma evolution.
2. The MHD Approximation

a. Strong Collisions: $\nu_{ei} \tau \gg 1$ (or $\lambda_{m}/L \ll 1$)

b. Non-relativistic $v_0^2/c^2 \ll 1$ (where $v_0 \sim \frac{L}{\tau}$)

c. Magnetized $\mu_0/L \ll 1$

3. In MHD Approximation, plasma fluctuations are quasi-neutral, $\sum_s n_s q_s = 0$

4. When Magnetic Reynolds Number $R_{em} = \frac{\mu_0 L v_0}{\eta} \gg 1$,

a condition often satisfied in most plasmas, the Magnetic Flux is Frozen-In to the plasma.

5. The high collisionality enables:

a. Isotropic pressure $\left\{ \begin{array}{l} \nabla \cdot \mathbf{p} \approx -\nabla p_s - \nabla \cdot \mathbf{p}_s \approx -\nabla p_s \end{array} \right.$

b. Negligible viscosity

c. No Heat Flow along field lines \leftarrow Probably the approximation most often violated in real plasmas of interest.

\Rightarrow Adiabatic $\frac{d}{dt} \left(\frac{p}{\rho^{\gamma}} \right) = 0$ \leftarrow no heat flow

6. MHD, due to its simplicity, is the most widely used plasma description

a. WARNING! MHD is often applied even in situations where it is not formally valid \Rightarrow "Enter at your own risk."

7. MHD is a valid, low-frequency limit of plasma behavior, $\omega \ll \omega_{ci}$

IV. (Continued)

C. Two Fluid Equations, Cold Plasma Approximation

i. Cold Plasma Approximation: $\frac{\omega}{k} \gg v_{te}$

- Often a good description of high-frequency plasma behavior
- At high frequencies, electron and ion responses are different
 \Rightarrow requires a Two-Fluid Approach
- Since $\tau \sim \frac{1}{\omega}$ is small, $\tau v_{te} \ll 1$ often, so collisional effects may be neglected.

2. In this limit, thermal spread of velocities is small, so ^{all} particles respond to fields in the same way, regardless of velocity
 \Rightarrow No velocity space effects, no wave-particle interactions.

D. Two Fluid Equations, Warm Plasma

1. One may extend the validity of the Two-Fluid Approach to lower frequencies by including thermal effects (pressure).

2. Still accounts for differences in ion and electron responses.

This occurs when: $\left. \begin{array}{l} a. r_{Li} \gtrsim L \\ b. \omega \gtrsim \omega_{ci} \end{array} \right\}$ Single-fluid MHD breaks down in these conditions.

3. Valid for ω High collisionality $\tau v_{te} \gg 1$

In this case, we have

- Isotropic pressure $\nabla \cdot \mathbf{P}_s \approx \nabla P_s$
- Adiabatic Equation of State $\frac{d}{dt} \left(\frac{P_s}{\rho_s} \right) = 0$

b. Moderate Collisionality $\tau v_{te} \gtrsim 1$

i. Anisotropic pressure $\mathbf{P}_s = \begin{bmatrix} P_{\perp s} & 0 & 0 \\ 0 & P_{\perp s} & 0 \\ 0 & 0 & P_{\parallel s} \end{bmatrix}$

ii. Double Adiabatic Equation of State (Chew-Goldberger-Low, CGL) $\frac{d}{dt} \left(\frac{P_{\perp s}}{n_s B} \right) = 0, \frac{d}{dt} \left(\frac{P_{\parallel s} B^2}{n_s^3} \right) = 0$

IV. D. (Continued)

4. Because of moderate to high collisionality, velocity-space effects are unimportant.
- Collisions enforce Maxwellian, or Bi-Maxwellian, distributions.
 \Rightarrow In this case, fluid approximations are a good description.
 (Local Thermodynamic Equilibrium)

E. Plasma Kinetic Equation (Boltzmann Equation)

- Required when velocity-space effects are important
 - Collisionless plasma $v_{ei} \tau \ll 1$
 - Finite plasma temperature (non-negligible pressure)
- The conditions above lead to strong wave-particle interaction at resonant velocities, $v = \frac{\omega}{k}$.
 - Landau damping (collisionless damping)
 - Cyclotron damping (collisionless damping)
- For a kinetic description, no Equation of State need be assumed.
 - Equation of State is needed for any fluid description to close the hierarchy of moment equations.

IV. Waves in Plasmas

- Why do we spend so much time studying waves in plasmas?
 - Waves are the natural plasma response to an applied perturbation
 - Small perturbations lead to linear wave behavior
 - Unstable waves may lead to growth to nonlinear amplitudes.

V. B. Map of Waves in Plasmas

Geometry

Homogeneous (Magnetized) Plasmas

Inhomogeneous Plasmas (Often Magnetized)

Plasma Model:

Closure

Vlasov, Fokker-Planck

Fluid Theory

Kinetic Theory

Fluid Theory

Kinetic Theory

Cold Plasma

Two-Fluid

MHD

Hot Plasma

Cold Plasma

Two-Fluid

MHD

Hot Plasma

WKB

WKB

Linear Waves:

Fluid Dispersion Relation $D(k, \omega)$ Real

Kinetic Dispersion Relation $D(k, \omega)$ Complex \Rightarrow Landau Damping

Eigenmode Calculations

Linear Wave Propagation

Instabilities:

Real Space (Fluid) Instabilities: MHD: Interchange, Kink, Sausage, Ion Temperature Gradient, Driftwave

Velocity Space (Kinetic) Instabilities: Two-stream, Bump-in-tail, etc.

Nonlinearity:

Solitons Nonlinear Waves BGK Modes

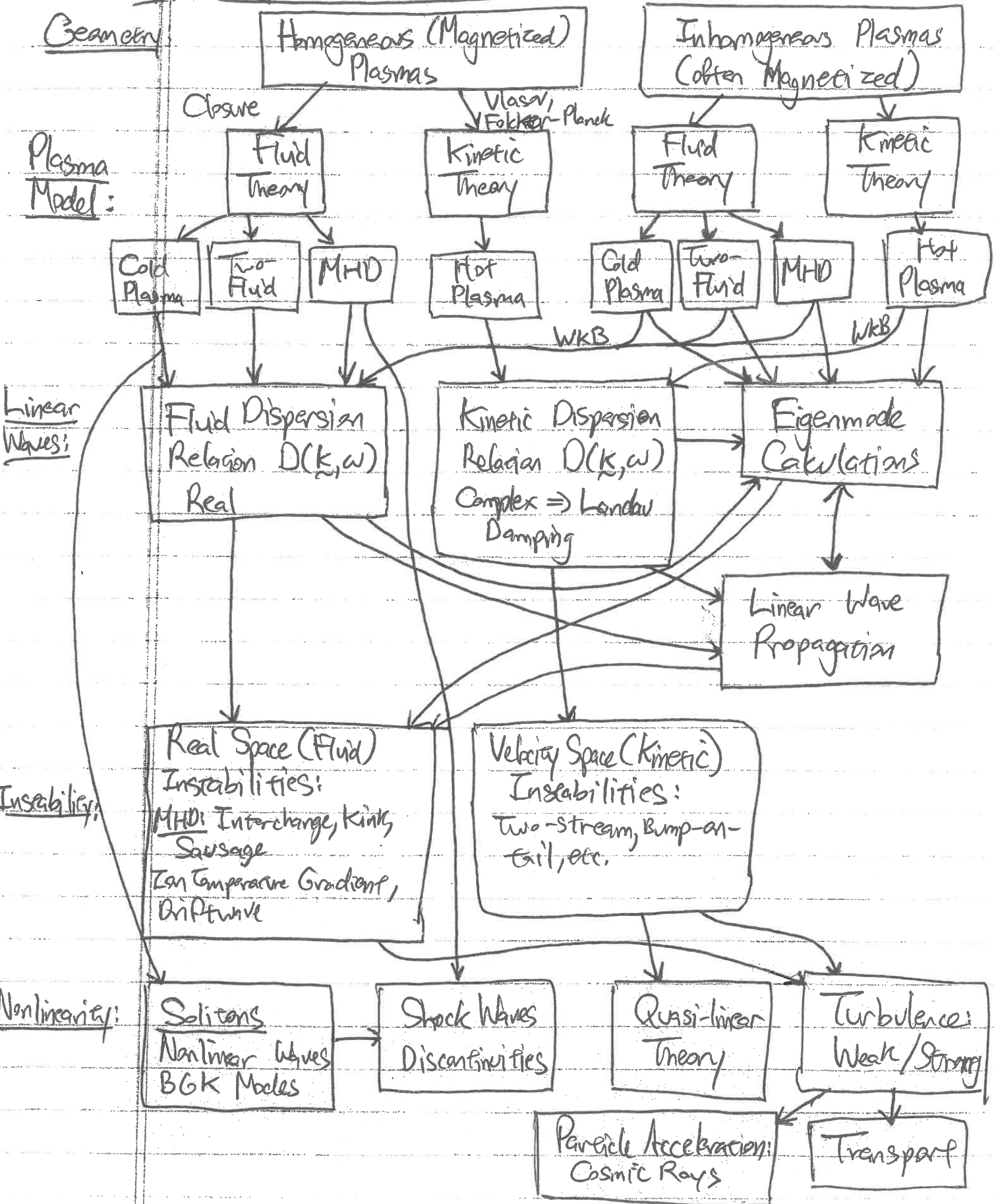
Shock Waves Discontinuities

Quasi-linear Theory

Turbulence: Weak/Strong

Particle Acceleration: Cosmic Rays

Transport



Lecture #24 (Continued)

VI. A Map of Plasma Physics

Hines (8)

Single Particle Motion
 Prescribed E & B
 Solve for X & v

Guiding Center Theory
 Adiabatic Invariance

Full Problem: Kinomom Eq.

Particles S, X →
 Lorentz Force Law

Fields E, B ←
 Maxwell's Equations

Plasma Equilibrium
 Only relevant for highly magnetized plasma
 $\nabla p = \mathbf{j} \times \mathbf{B} + \rho \mathbf{g}$

Stable: Waves

Cold Plasma Waves
MHD Waves

Hot Plasma Waves

- Landau Damping
- Quasi-linear theory

Unstable: Instabilities

Fluid: MHD - Interchange
 Kink, Sausage, Rayleigh-Taylor

Drift waves
 Growth to M₀ amplitudes

Kinetic: Beam, Two-Stream
 Temperature Anisotropy

Firchak, Alfvén, Zon
 cyclotron, whistler

- Often saturate at moderate amplitudes
- Zoo of kinetic instabilities

Statistical Approach Boltzmann Equation

$\frac{dN}{dt} + \mathbf{v} \cdot \nabla N + \frac{qB}{m\Omega} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \frac{\partial N}{\partial \mathbf{v}} = \left(\frac{dN}{dt} \right)_{coll}$

$\rho_i = \int \mathbf{v} q_i N_i \mathbf{v} d\mathbf{v}$ $\mathbf{j} = \int \mathbf{v} q_i N_i \mathbf{v} d\mathbf{v}$

Maxwell's Equations

↓ Closure

Two Fluid Equations:
 Local Thermodynamic Equilibrium
 ⇒ Maxwellian Distribution
 Useful for cold, dense plasmas

MHD: Single Fluid Theory
 Frozen-in flux
 Isotropic pressure
 No heat flow

Nonlinear EPICs:
 Quasilinear Theory
 Shocks
 Turbulence
 Solitons

perturbations

Closure