## 29:293 Homework \#6

Due at the beginning of class, Thurs, April 2, 2015.

1. Here we will apply a simplified version of Multiple-Timescale Analysis to the problem of particle motion in constant, uniform $\mathbf{E}$ and $\mathbf{B}$ fields.
We assume a right-handed, orthonormal basis aligned with the direction of the magnetic field ( $\left.\hat{\mathbf{e}}_{1}, \hat{\mathbf{e}}_{2}, \hat{\mathbf{b}}\right)$ such that $\hat{\mathbf{e}}_{1} \times \hat{\mathbf{e}}_{2}=\hat{\mathbf{b}}$. The Lorentz Force Law is

$$
m \frac{d \mathbf{v}}{d t}=q(\mathbf{E}+\mathbf{v} \times \mathbf{B})
$$

for an electric field $\mathbf{E}=E_{1} \hat{\mathbf{e}}_{1}+E_{2} \hat{\mathbf{e}}_{2}+E_{\|} \hat{\mathbf{b}}$ and a magnetic field $\mathbf{B}=B_{0} \hat{\mathbf{b}}$. For this problem, we will take the case $E_{\|}=0$.
(a) First, let us convert the dimensional form of the Lorentz Force Law above to a dimensionless equation. Derive the dimensionless form

$$
\begin{equation*}
\frac{d \mathbf{v}^{\prime}}{d t^{\prime}}=\mathbf{E}^{\prime}+\mathbf{v}^{\prime} \times \hat{\mathbf{b}} \tag{1}
\end{equation*}
$$

for dimensionless quantities $t^{\prime}=\omega_{c} t, \mathbf{v}^{\prime}=\mathbf{v} / v_{\perp}$, and $\mathbf{E}^{\prime}=\frac{\mathbf{E}}{B_{0} v_{\perp}}$ where $v_{\perp}=\sqrt{v_{1}^{2}+v_{2}^{2}}$.
(b) Verify that the quantity $E^{\prime}=\left|\mathbf{E}^{\prime}\right|$ is dimensionless (in the SI system of units).
(c) Show that the condition $E^{\prime} \ll 1$ means that the $\mathbf{E} \times \mathbf{B}$ drift is slow compared to the perpendicular velocity, $\left|\mathbf{v}_{E}\right| \ll v_{\perp}$.
(d) Assuming $E^{\prime} \ll 1$, the timescales of the Larmor motion and the $\mathbf{E} \times \mathbf{B}$ drift are well separated. For the expansion parameter, take $\epsilon=E^{\prime} \ll 1$. As an aid in the bookkeeping for the order of magnitude of each term, we can add an $\epsilon$ to the electric field term in our equation to remind us of its order,

$$
\begin{equation*}
\frac{d \mathbf{v}^{\prime}}{d t^{\prime}}=\epsilon \mathbf{E}^{\prime}+\mathbf{v}^{\prime} \times \hat{\mathbf{b}} \tag{2}
\end{equation*}
$$

We'll assume a fast timescale $t^{\prime}$ and a slow timescale $\tau^{\prime}=\epsilon t^{\prime}$. Decompose the total velocity into rapidly varying piece $\mathbf{v}_{1}^{\prime}$ and a smaller slowly varying piece $\mathbf{v}_{2}^{\prime}, \mathbf{v}^{\prime}=\mathbf{v}_{1}^{\prime}\left(t^{\prime}\right)+\epsilon \mathbf{v}_{2}^{\prime}\left(\tau^{\prime}\right)$.
Write down the expansion of $d / d t^{\prime}$ assuming two timescales.
(e) Derive the equation at $\mathcal{O}(1)$ and solve for $\mathbf{v}_{1}^{\prime}\left(t^{\prime}\right)$ given the (dimensional) initial conditions at $t=0$ of $\mathbf{v}=$ $v_{\perp} \hat{\mathbf{e}}_{1}+v_{\| 0} \hat{\mathbf{b}}$.
(f) Derive the equation at $\mathcal{O}(\epsilon)$. Solve for $\mathbf{v}_{2}^{\prime}\left(\tau^{\prime}\right)$. HINT: Do not forget to treat $t^{\prime}$ and $\tau^{\prime}$ as independent variables.
(g) Sum the solution for each order to get the total solution $\mathbf{v}^{\prime}\left(t^{\prime}, \tau^{\prime}\right)$. Convert back to dimensional form to yield the final, complete solution $\mathbf{v}(t)$.
2. Recall the Child-Langmuir Law for a 1-D electrostatic plasma of hydrogen with isothermal electrons with temperature $T_{e}$ and cold ions,

$$
\begin{equation*}
j_{i}=\frac{4}{9} \epsilon_{0}\left(\frac{2 e}{m_{i}}\right)^{1 / 2} \frac{\phi_{w}^{3 / 2}}{d^{2}} \tag{3}
\end{equation*}
$$

which expresses the space-charge limited ion current across the sheath in the limit $-e \phi_{w} / T_{e} \gg 1$ as a function of sheath width $d$ and the potential difference $\phi_{w}$ between the wall and the potential at the sheath edge $x=d$. Recall that the potential at the sheath edge is chosen to define a potential of zero, $\phi(d)=0$. Note that we absorb Boltzmann's constant to give temperature in units of energy.
(a) Taking the ion current to be given by $j_{i}=e n_{d} c_{s}$, where $c_{s}=\sqrt{T_{e} / m_{i}}$ and the $n_{d}$ is ion density at the sheath edge, compute an expression for the sheath width $d$ as a function of the wall potential $\phi_{w}$, the electron temperature $T_{e}$, and the Debye length computed using the plasma conditions at the sheath edge $x=d$.
(b) For typical laboratory plasma parameters of $T_{e}=5 \mathrm{eV}$ and $n_{d}=10^{18} \mathrm{~m}^{3}$, compute the width of the sheath for a wall voltage of $\phi_{w}=-300 \mathrm{~V}$.
3. For a Langmuir probe trace using a cylindrical probe (for which the electron saturation current does not become constant), (a) compute the electron temperature (in eV ) using the data in the table below, (b) estimate the plasma potential $\phi_{p}$, and (c) estimate floating potential $\phi_{f}$.

| Probe Bias (V) | - Probe Current (A) |
| :--- | :--- |
| -65.00 | -0.0001290 |
| -60.00 | -0.0001290 |
| -55.00 | -0.0000860 |
| -50.00 | -0.0000860 |
| -45.00 | -0.0000430 |
| -40.00 | -0.0000430 |
| -35.00 | 0.0000000 |
| -30.00 | 0.0000430 |
| -26.00 | 0.0001730 |
| -24.00 | 0.0003020 |
| -22.00 | 0.0004740 |
| -20.00 | 0.0009060 |
| -18.00 | 0.0015960 |
| -16.00 | 0.0032350 |
| -14.00 | 0.0041410 |
| -12.00 | 0.0046580 |
| -10.00 | 0.0051330 |
| -8.00 | 0.0055210 |
| -6.00 | 0.0058230 |
| -4.00 | 0.0062540 |
| -2.00 | 0.0064270 |
| 0.00 | 0.0068150 |

