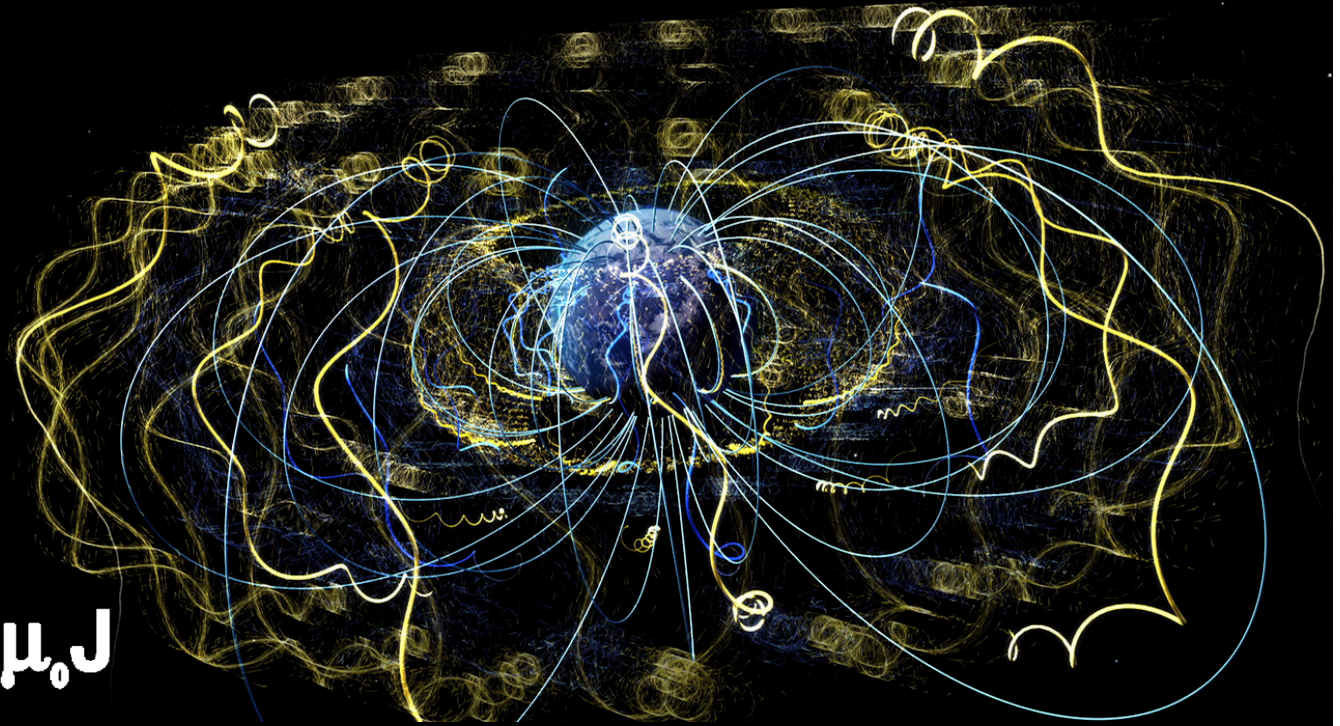


$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} + \mu_0 \mathbf{J}$$



Electricity and Magnetism I: 3811

Professor Jasper Halekas
Van Allen 301
MWF 9:30-10:20 Lecture

Continuous Charge Distribution

$$\frac{1}{2} \sum_i q_i V(\vec{r}_i) \rightarrow \boxed{W = \frac{1}{2} \int \rho V d\tau}$$

In terms of \vec{E} :

$$\rho = \epsilon_0 \nabla \cdot \vec{E}$$

$$W = \frac{\epsilon_0}{2} \int (\nabla \cdot \vec{E}) V d\tau$$

Integrate by parts

$$W = \frac{\epsilon_0}{2} \left[- \int \vec{E} \cdot \nabla V d\tau + \oint V \vec{E} \cdot d\vec{a} \right]$$

$$= \frac{\epsilon_0}{2} \int \vec{E} \cdot \vec{E} d\tau + \frac{\epsilon_0}{2} \oint V \vec{E} \cdot d\vec{a}$$

For all space $\rightarrow 0$

$$\Rightarrow \boxed{W = \frac{1}{2} \epsilon_0 \int E^2 d\tau}$$

- Electric field contains energy!

Example

- Shell w/ charge Q
radius R

$$\vec{E} = \begin{cases} \frac{Q}{(4\pi\epsilon_0 r^2)} \hat{r} & r > R \\ 0 & r < R \end{cases}$$

$$V = \begin{cases} \frac{Q}{(4\pi\epsilon_0 r)} & r > R \\ \frac{Q}{(4\pi\epsilon_0 R)} & r < R \end{cases}$$

$$W = \frac{1}{2} \int \rho V d\tau = \frac{1}{2} \int \sigma V da$$

$$= \frac{1}{2} \int \frac{Q}{4\pi R^2} \frac{Q}{4\pi\epsilon_0 R} \cdot da$$

$$= \frac{1}{2} \frac{Q}{4\pi R^2} \cdot \frac{Q}{4\pi\epsilon_0 R} \cdot 4\pi R^2$$

$$= \frac{Q^2}{8\pi\epsilon_0 R}$$

$$W = \frac{1}{2} \epsilon_0 \int E^2 d\tau$$

$$\frac{1}{2} \epsilon_0 \int_0^{2\pi} \int_0^\pi \int_R^\infty \frac{Q^2}{(4\pi\epsilon_0 r^2)^2} \cdot r^2 dr \sin\theta d\theta d\phi$$

$$= \frac{Q^2}{32\pi^2\epsilon_0} \cdot \int_R^\infty \frac{1}{r^2} dr \int_0^{2\pi} \int_0^\pi \sin\theta d\theta d\phi$$

$$= \frac{Q^2}{32\pi^2\epsilon_0} \cdot \frac{1}{R} \cdot 4\pi = \frac{Q^2}{8\pi\epsilon_0 R} //$$

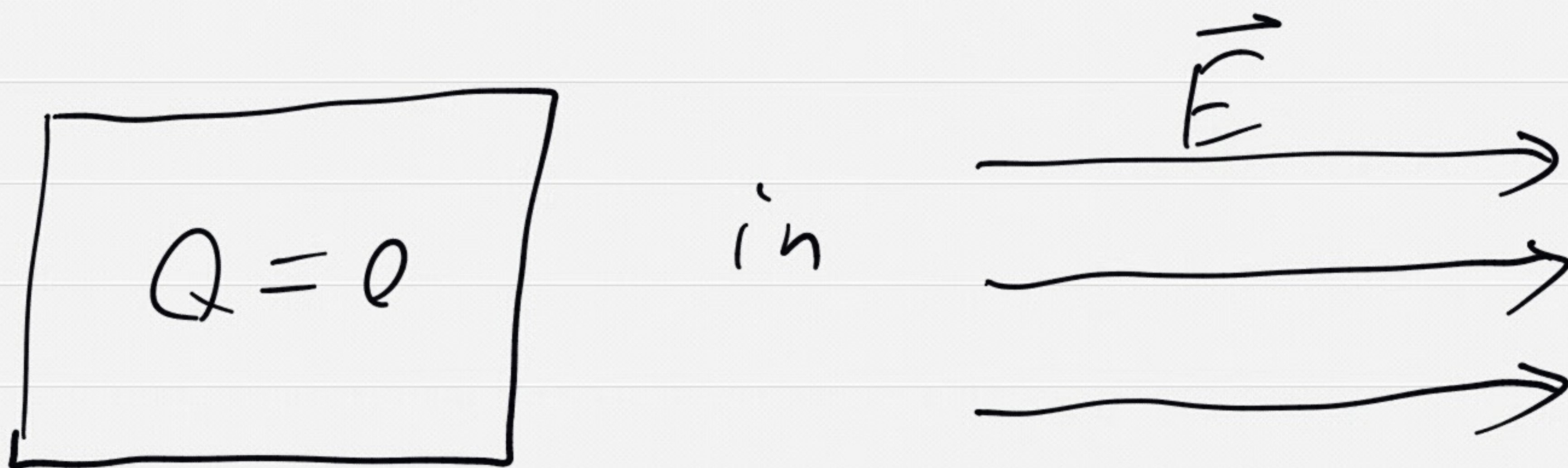
Conductors

- Materials w/ many unbound "free electrons" which can quickly and easily move in response to any electric field

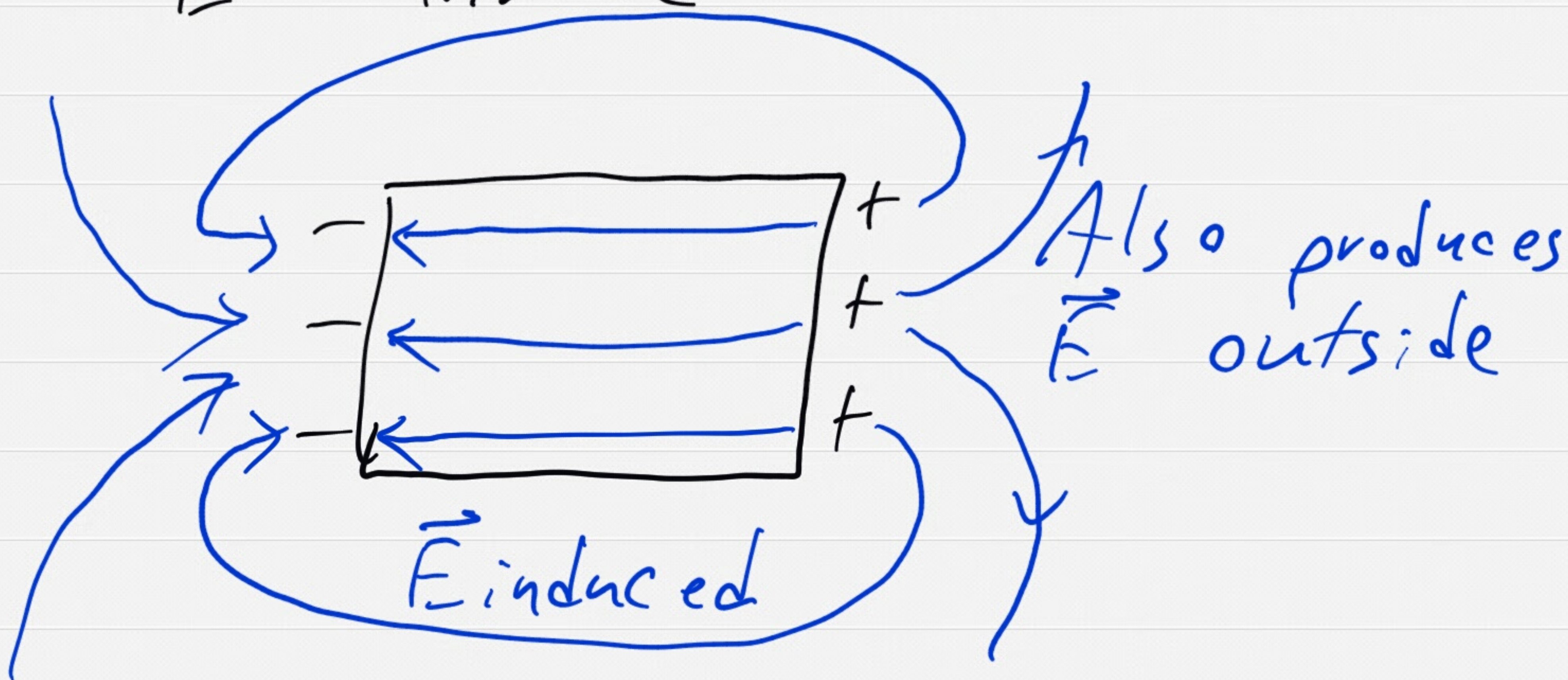
Steady-State Conductors

1. $\vec{E} = 0$ inside
2. $\rho = 0$ inside
($\nabla \cdot \vec{E} = 0$)
3. Net charge only on surface
4. Conductor is equipotential
 $\Delta V = -\int \vec{E} \cdot d\vec{l} = 0$
5. \vec{E} normal to surface
 $\Delta \vec{E} = \frac{\sigma}{\epsilon_0} \hat{n} = \vec{E}_{\text{outside}}$

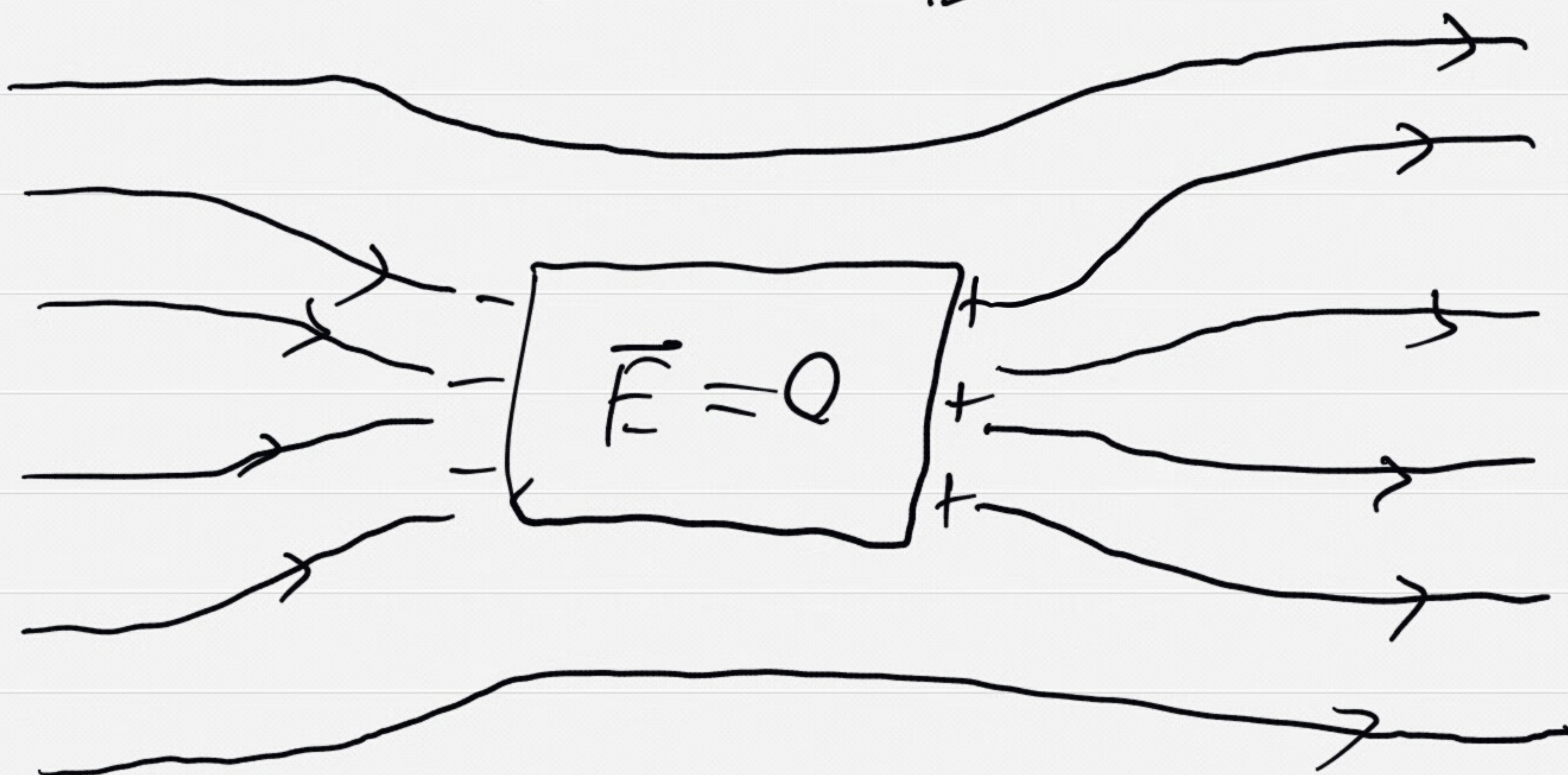
Induced Charge and Field



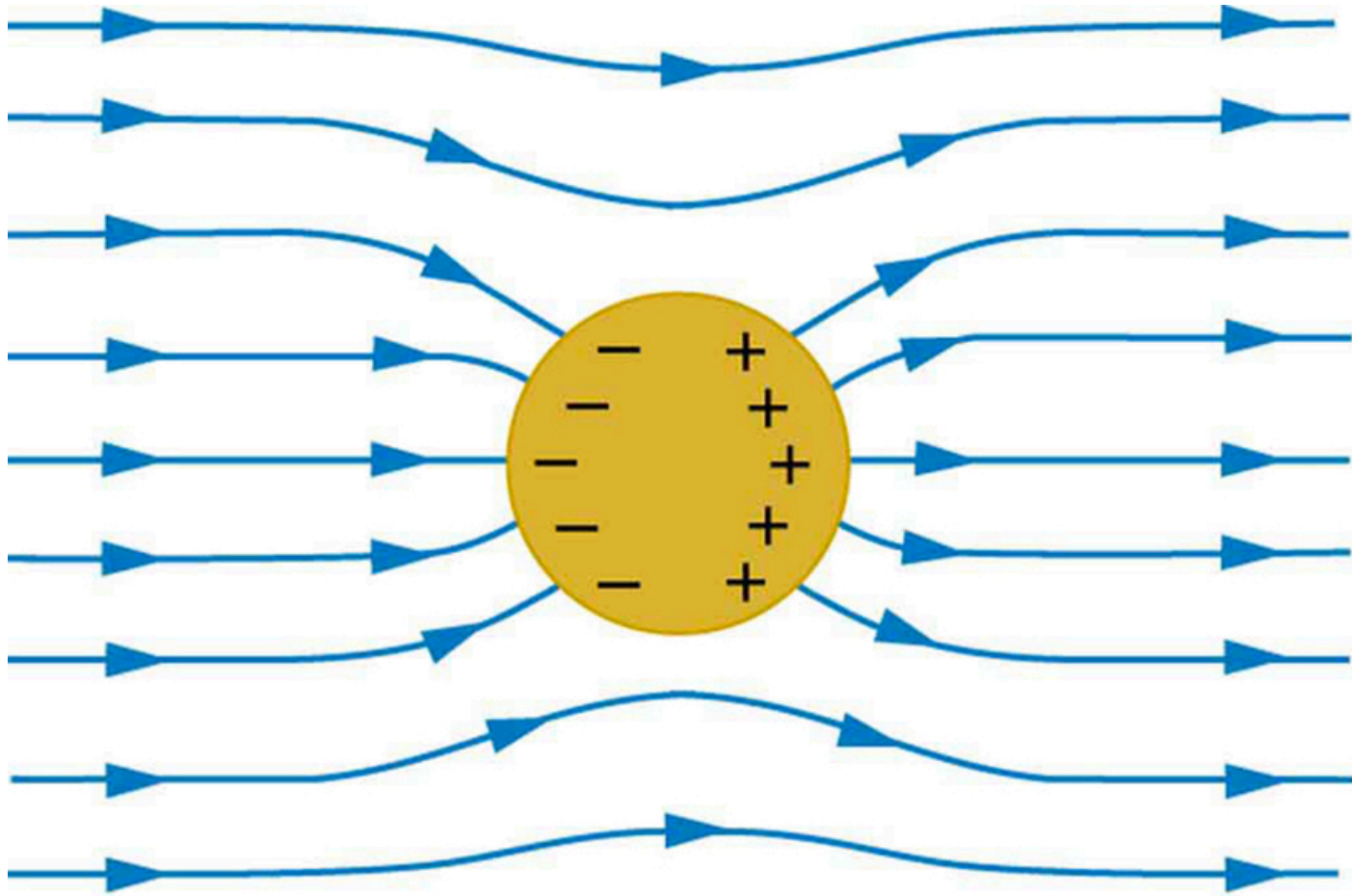
Charge flows to cancel \vec{E} inside



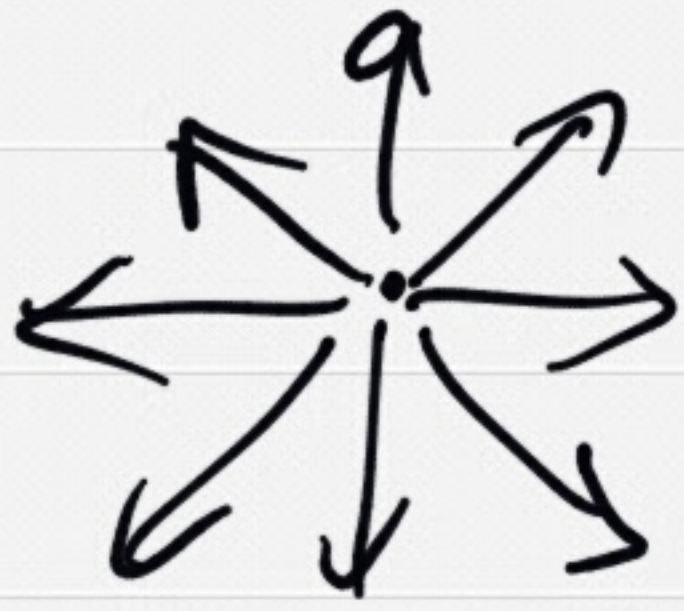
Total \vec{E}



Conductors

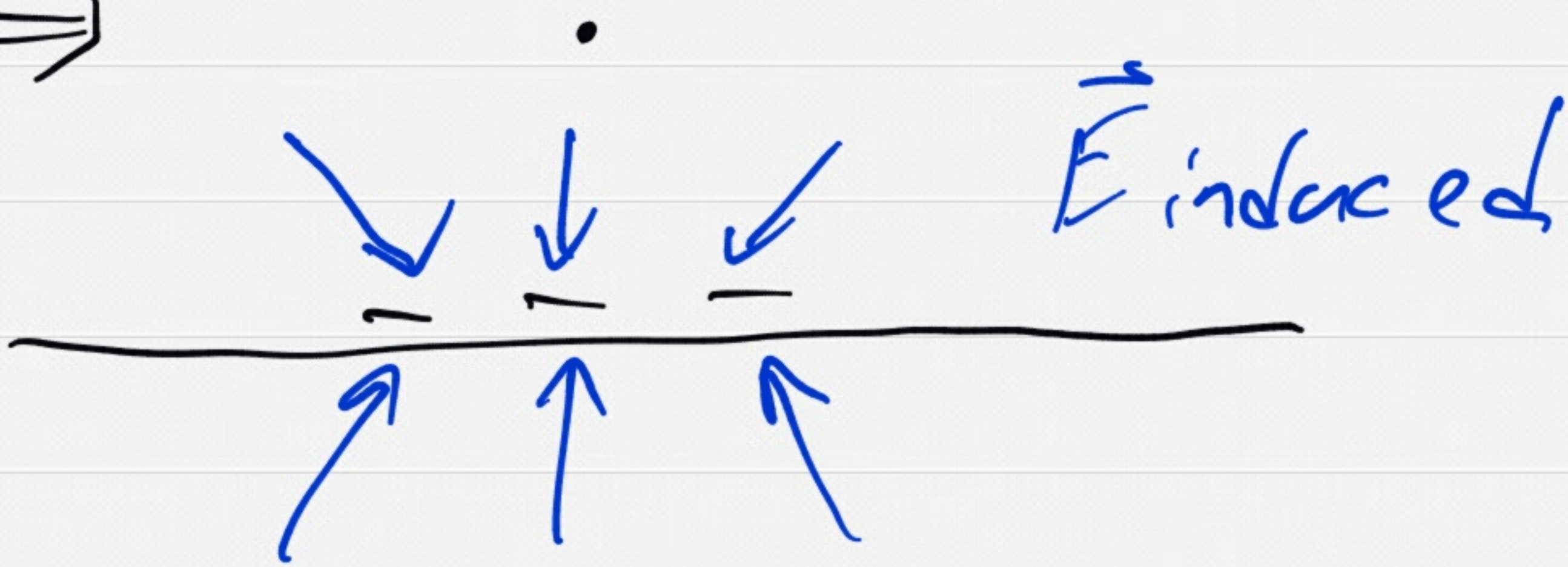


Charge Near Conductor

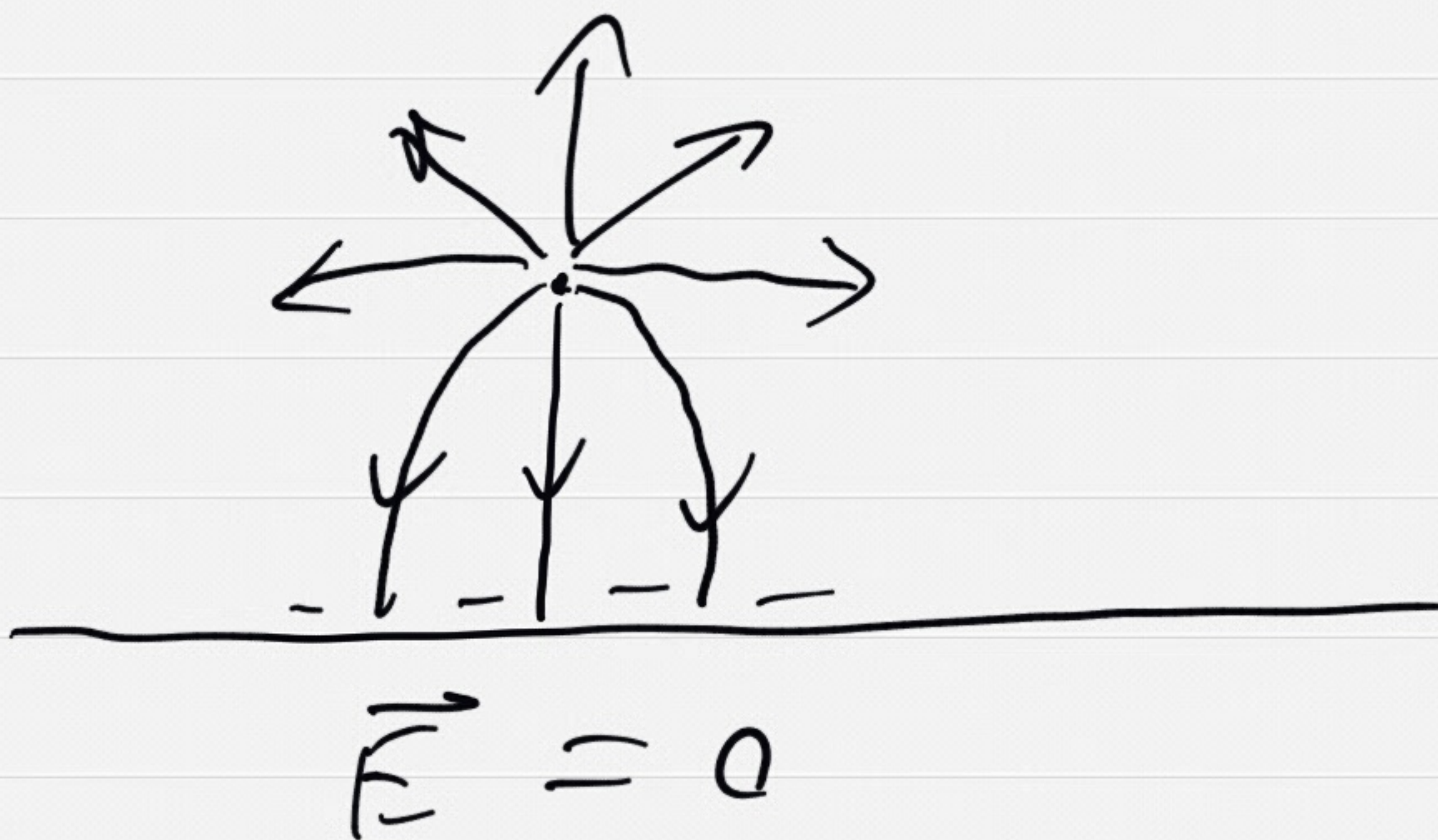


conductor

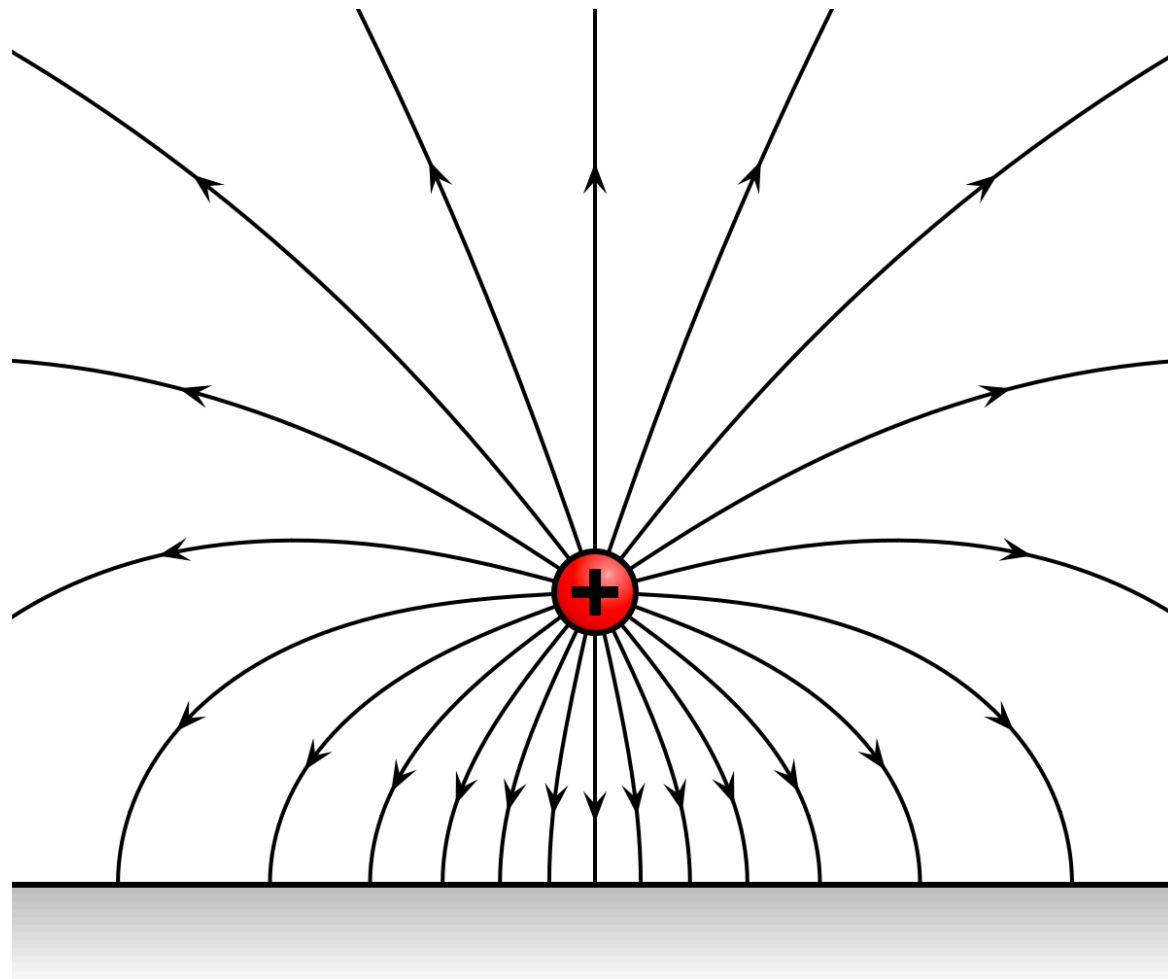
⇒



Total Field



Charge Near Conductor



Conductors

