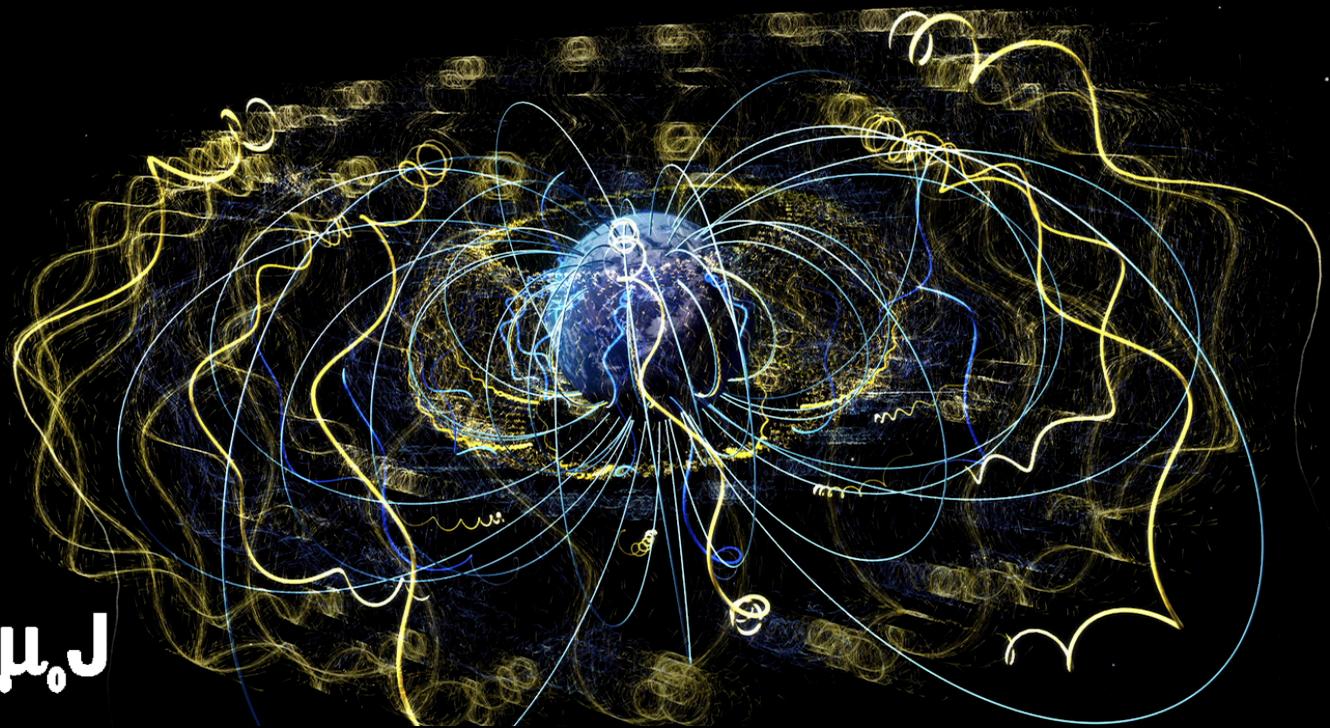


$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = - \frac{\partial \mathbf{B}}{\partial t}$$

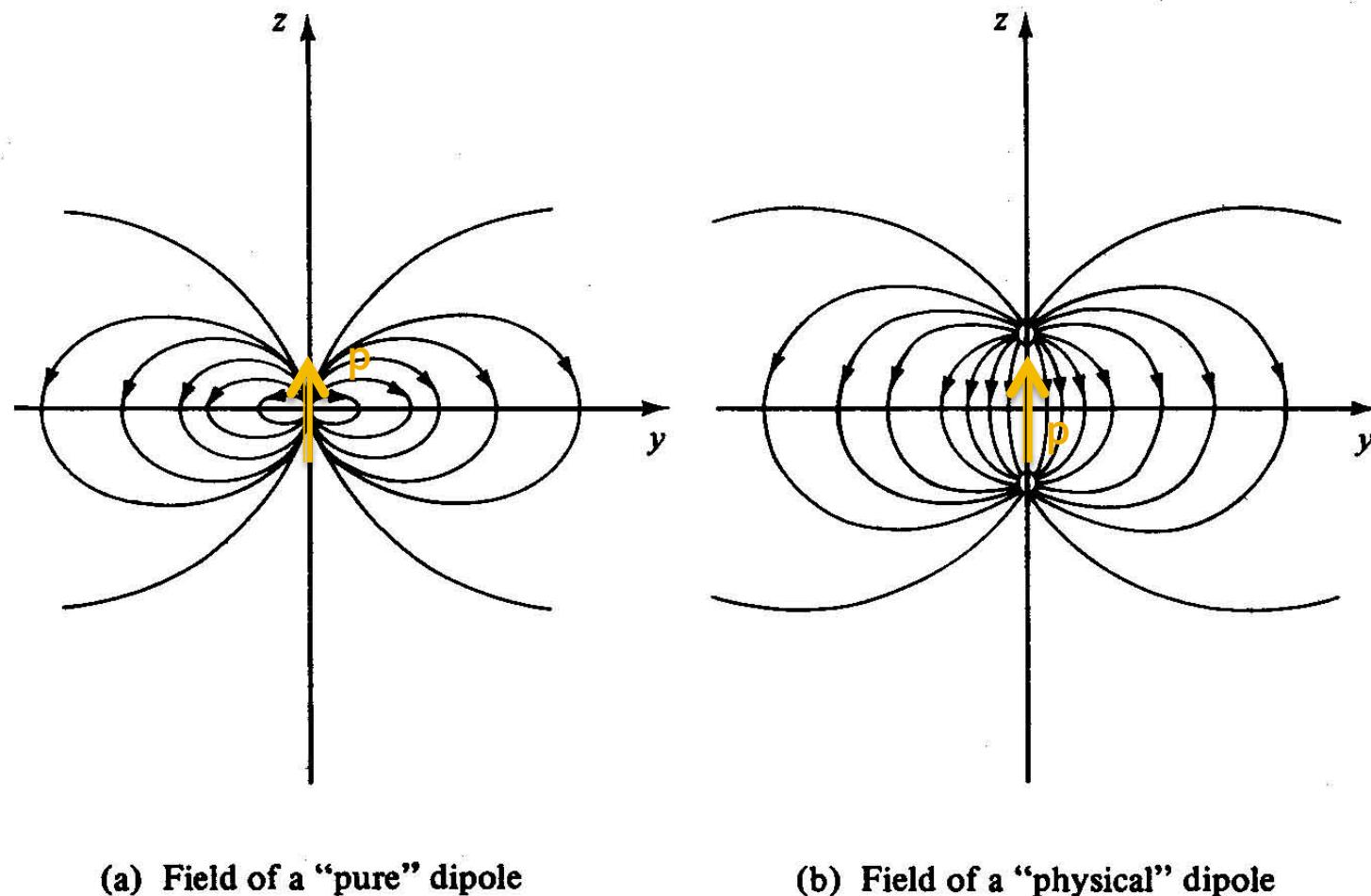
$$\nabla \times \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} + \mu_0 \mathbf{J}$$



Electricity and Magnetism I: 3811

Professor Jasper Halekas
Van Allen 301
MWF 9:30-10:20 Lecture

Pure Vs. Physical Dipole

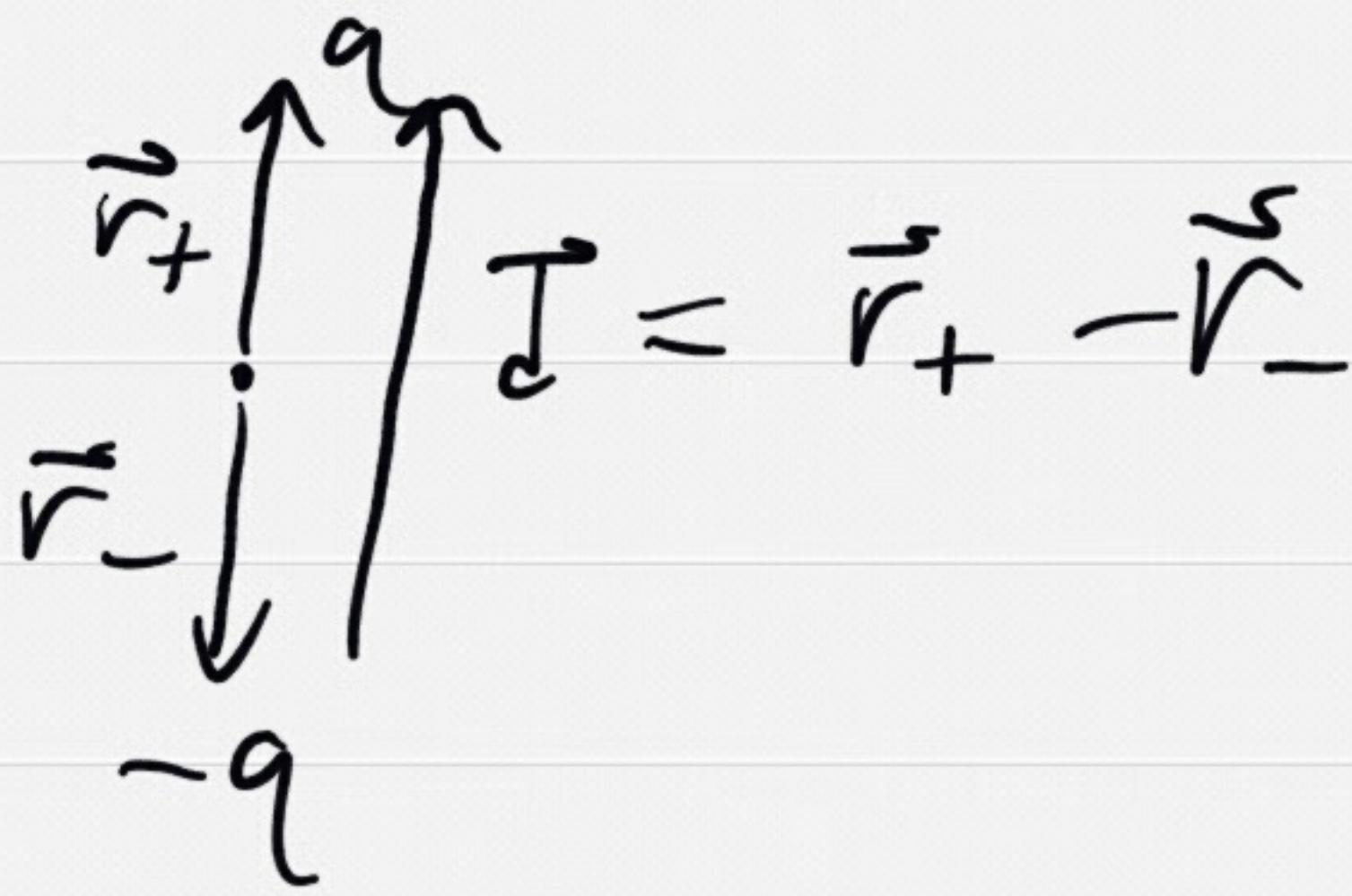


(a) Field of a “pure” dipole

(b) Field of a “physical” dipole

Figure 3.32

Physical Dipole



$$\begin{aligned} \text{Monopole moment } Q &= \int \rho(\vec{r}') d\tau' \\ &= 0 \end{aligned}$$

$$\begin{aligned} \text{Dipole moment } \vec{p} &= \int \vec{r}' \rho(\vec{r}') d\tau' \\ \rho(\vec{r}') &= q \delta^3(\vec{r}' - \vec{r}_+) - q \delta^3(\vec{r}' - \vec{r}_-) \\ \Rightarrow \vec{p} &= q \left(\int (\vec{r}' - \delta^3(\vec{r}' - \vec{r}_+) - \vec{r}_-) \delta^3(\vec{r}' - \vec{r}_-) d\tau' \right) \\ &= q (\vec{r}_+ - \vec{r}_-) \\ &= q \vec{d} \end{aligned}$$

$$V_{\text{dip}} = \frac{\vec{p} \cdot \hat{r}}{4\pi\epsilon_0 r^2} = \frac{p \cos\alpha}{4\pi\epsilon_0 r^2}$$

Note : - physical dipole also has non-zero quadrupole, octupole, etc.
- perfect dipole if $d \rightarrow 0$

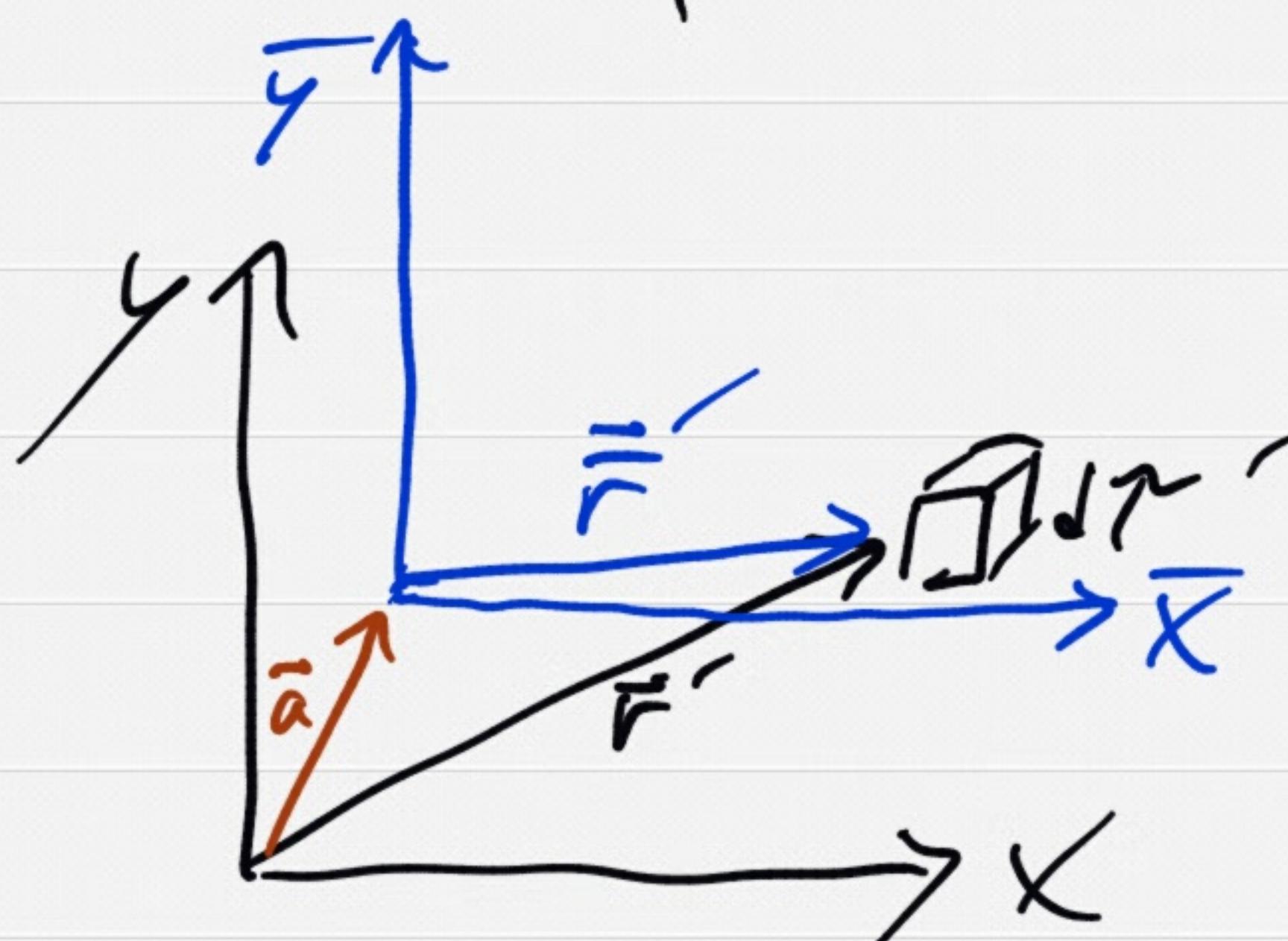
Moments & Origin of Coordinates

- Monopole moment

$$\int \rho(\vec{r}') d\tau' = Q \text{ ind. of origin}$$

- Dipole moment

$$\vec{P} = \int \vec{r}' \rho(\vec{r}') d\tau'$$



$$\bar{\vec{P}} = \int \bar{\vec{r}'} \rho(\bar{\vec{r}'}) d\tau'$$

$$= \int (\vec{r}' - \vec{a}) \rho(\vec{r}') d\tau'$$

$$= \int \vec{r}' \rho(\vec{r}') d\tau' - \int \vec{a} \rho(\vec{r}') d\tau'$$

$$= \vec{P} - \vec{a} Q$$

$$\bar{\vec{P}} \neq \vec{P} \text{ unless } Q = 0$$

Dipole Electric Field

$$V_{dip}(r, \theta) = \frac{\hat{r} \cdot \vec{p}}{4\pi\epsilon_0 r^2}$$
$$= \frac{p \cos\theta}{4\pi\epsilon_0 r^2}$$

for dipole @ origin
aligned w/ z-axis

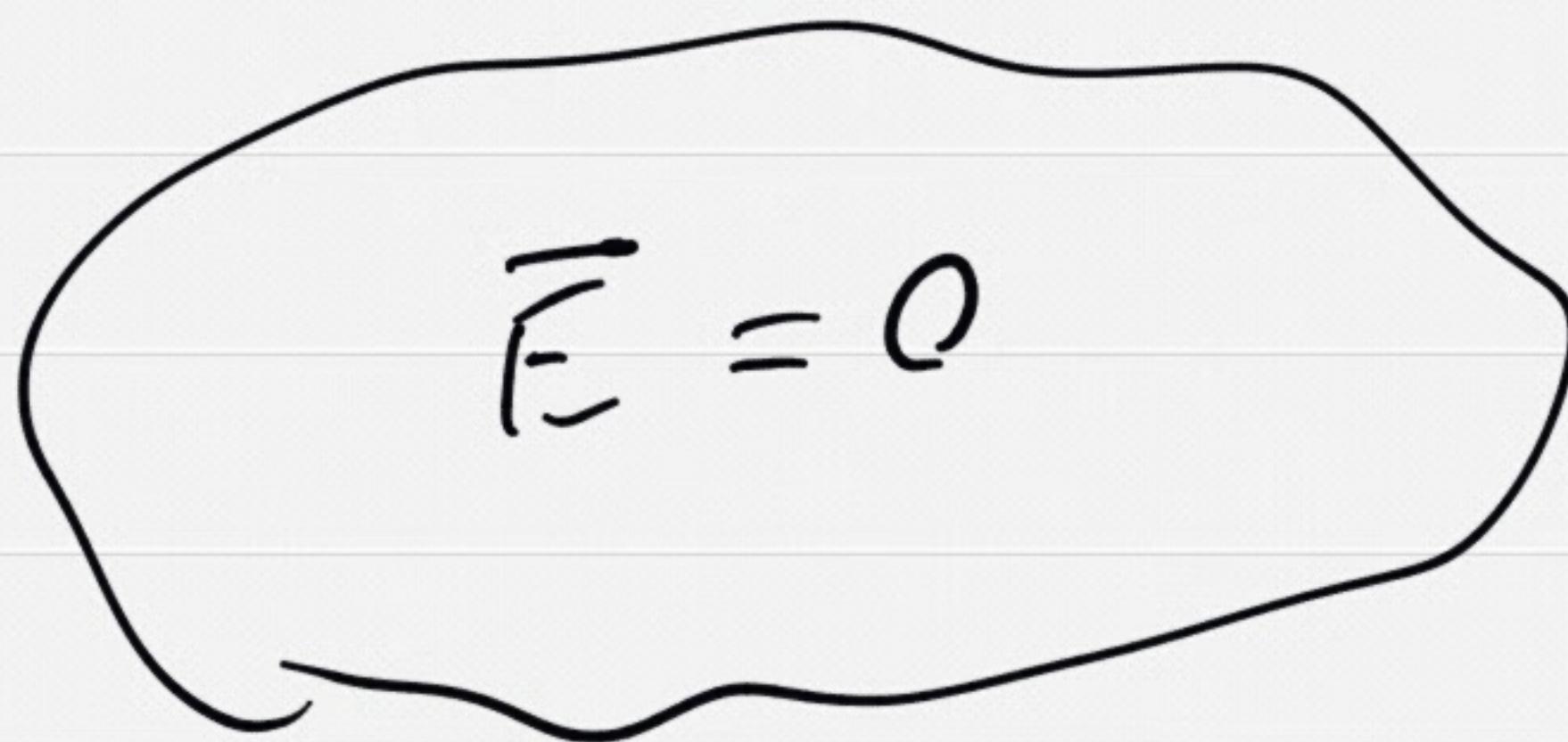
$$E_r = -\frac{\partial V}{\partial r} = \frac{2p \cos\theta}{4\pi\epsilon_0 r^3}$$

$$E_\theta = -\frac{1}{r} \frac{\partial V}{\partial \theta} = \frac{p \sin\theta}{4\pi\epsilon_0 r^3}$$

$$E_\phi = -\frac{1}{rs\sin\theta} \frac{\partial V}{\partial \phi} = 0$$

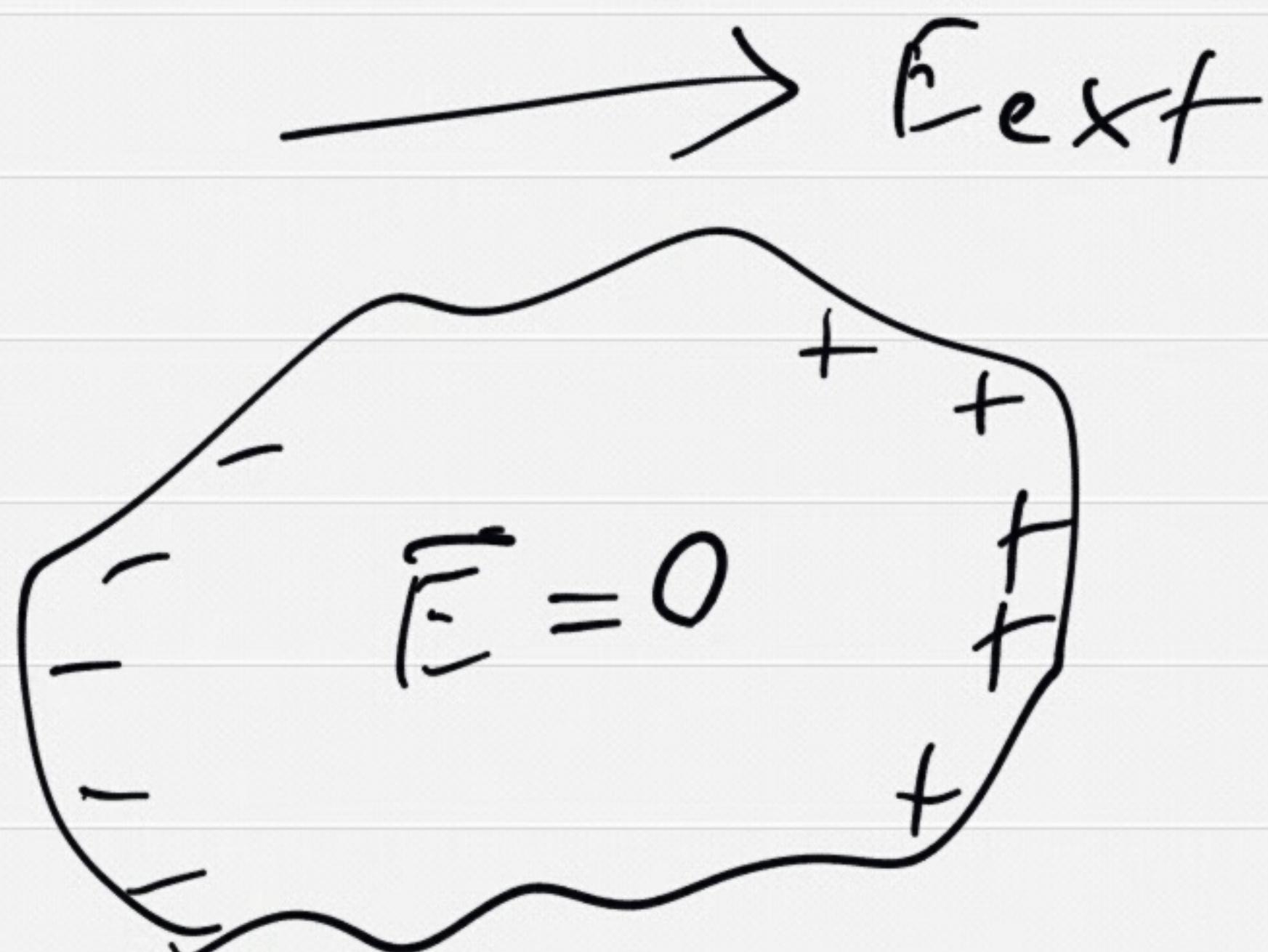
$$\vec{E}_{dip} = \frac{p}{4\pi\epsilon_0 r^3} [2 \cos\theta \hat{r} + \sin\theta \hat{\theta}]$$

Conductor



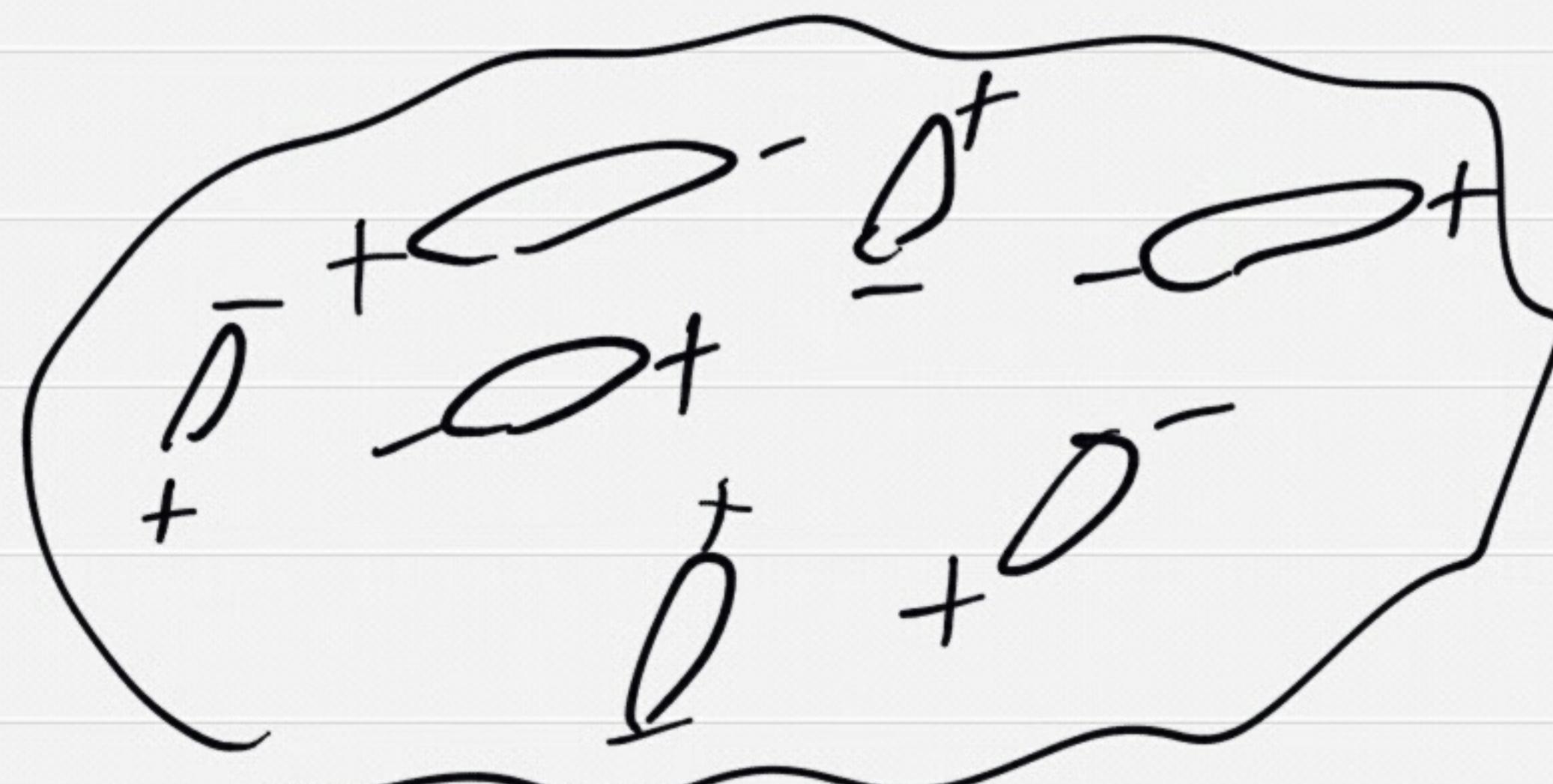
All charge
on surface

Conductor in external \vec{E}



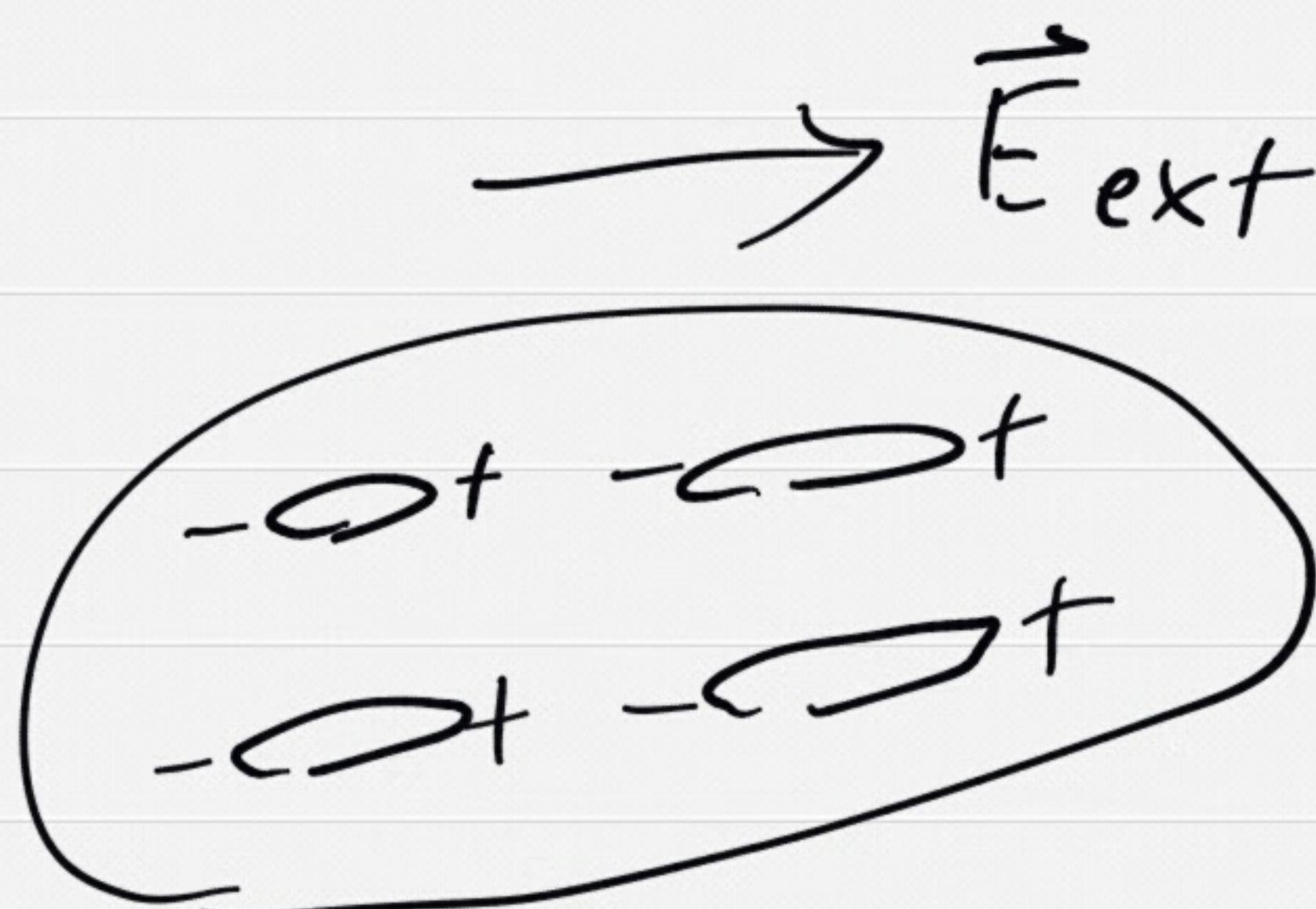
charge cancels
out \vec{E}_{ext}
in conductor

Insulator



- Made up of dipoles
- Usually randomly oriented so net field \neq zero

Insulator in External \vec{E}

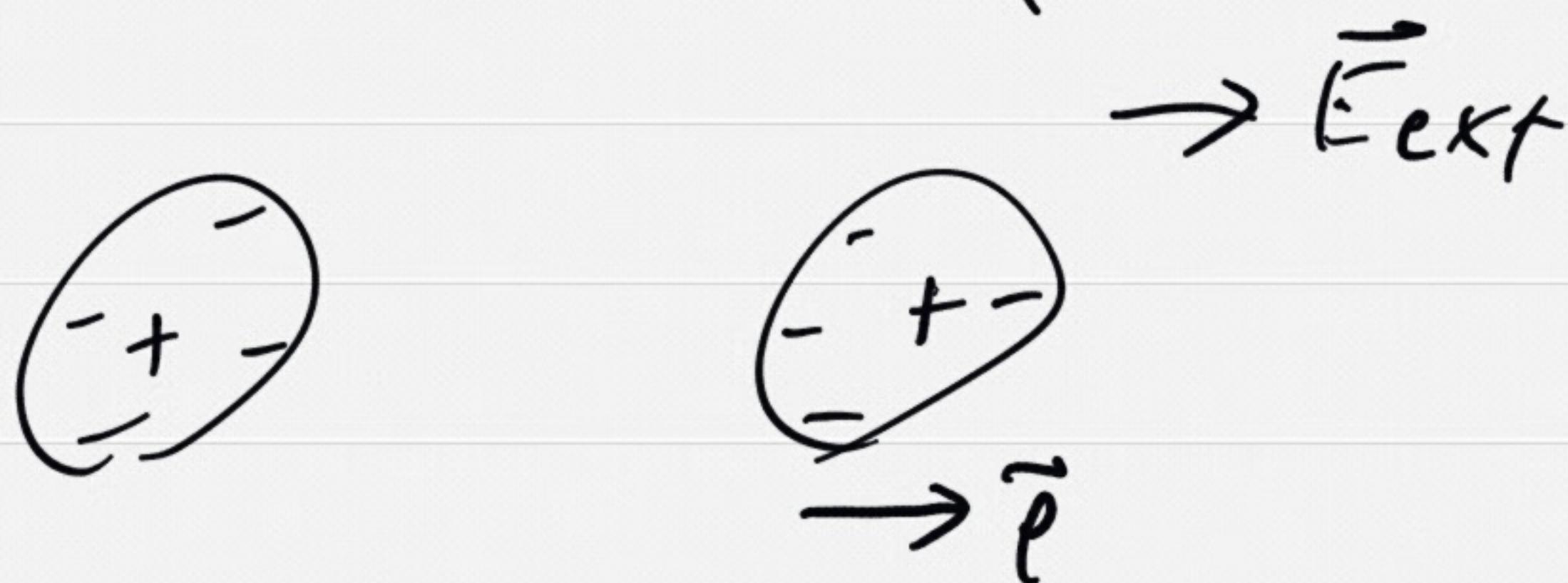


- Dipoles align w/ \vec{E}_{ext}

- weaken but don't cancel out \vec{E}_{ext}

Two kinds of dipoles in material

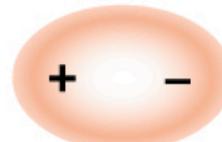
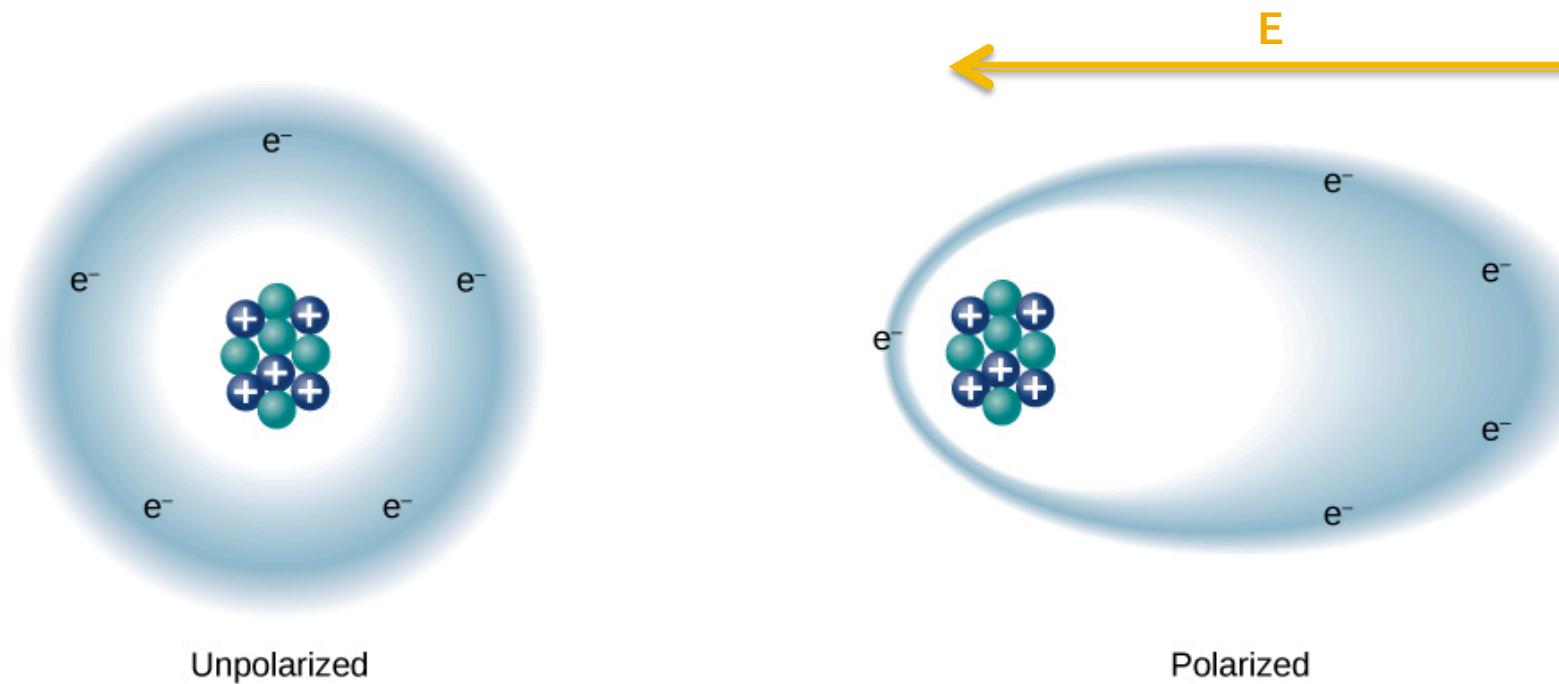
Induced Dipoles



Intrinsic Dipoles



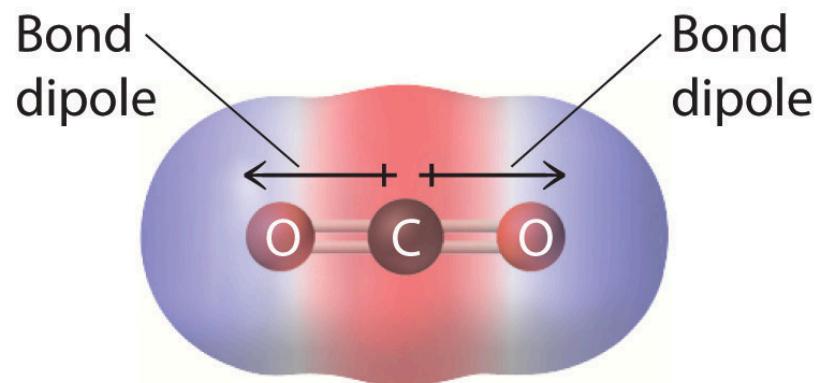
Induced Dipoles (Polarizability)



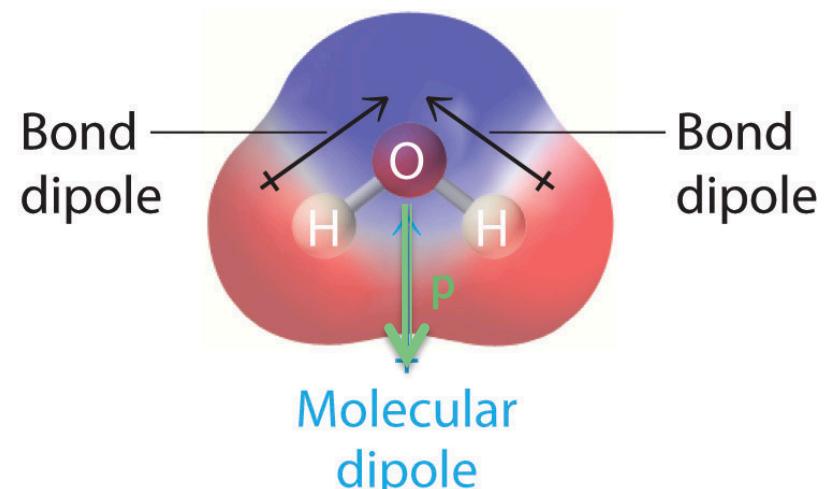
Large-scale view of polarized atom

Intrinsic Dipoles (Polar Molecules)

Electronegativity: H 2.1, C 2.5, O 3.5 => electrons closer to O than C, H

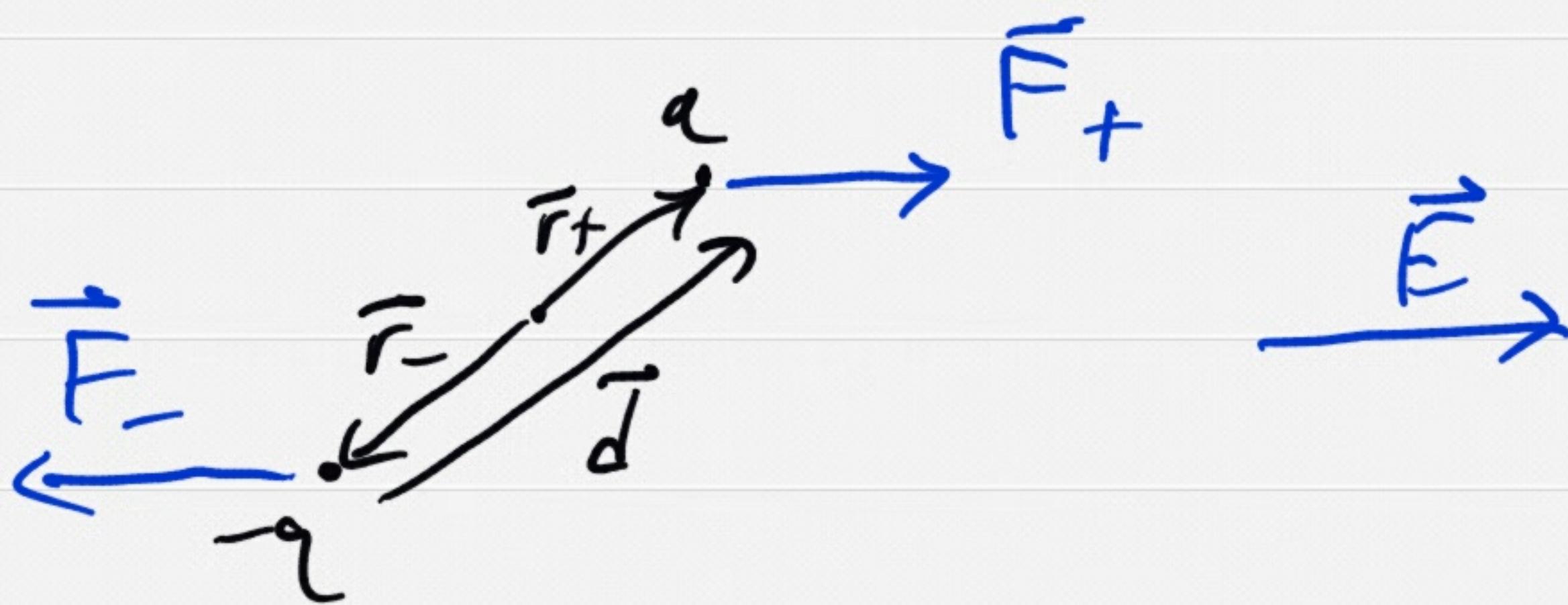


(a) No net dipole moment



(b) Net dipole moment

Force on Dipole (Tinker-Toy Model)



$$\begin{aligned}\vec{F} &= \vec{F}_+ + \vec{F}_- \\ &= q \vec{E}_+ - q \vec{E}_- \\ &= q \Delta \vec{E}\end{aligned}$$

$$\begin{aligned}&= q (\vec{J} \cdot \nabla) \vec{E} \\ &= \boxed{(\vec{P} \cdot \nabla) \vec{E}}\end{aligned}$$

$$\begin{aligned}\vec{\tau} &= \vec{r}_+ \times \vec{F}_+ + \vec{r}_- \times \vec{F}_- \\ &= \vec{J}/2 \times q \vec{E} + -\vec{J}/2 \times -q \vec{E} \\ &= q \vec{J} \times \vec{E} \\ &= \boxed{\vec{P} \times \vec{E}}\end{aligned}$$