

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

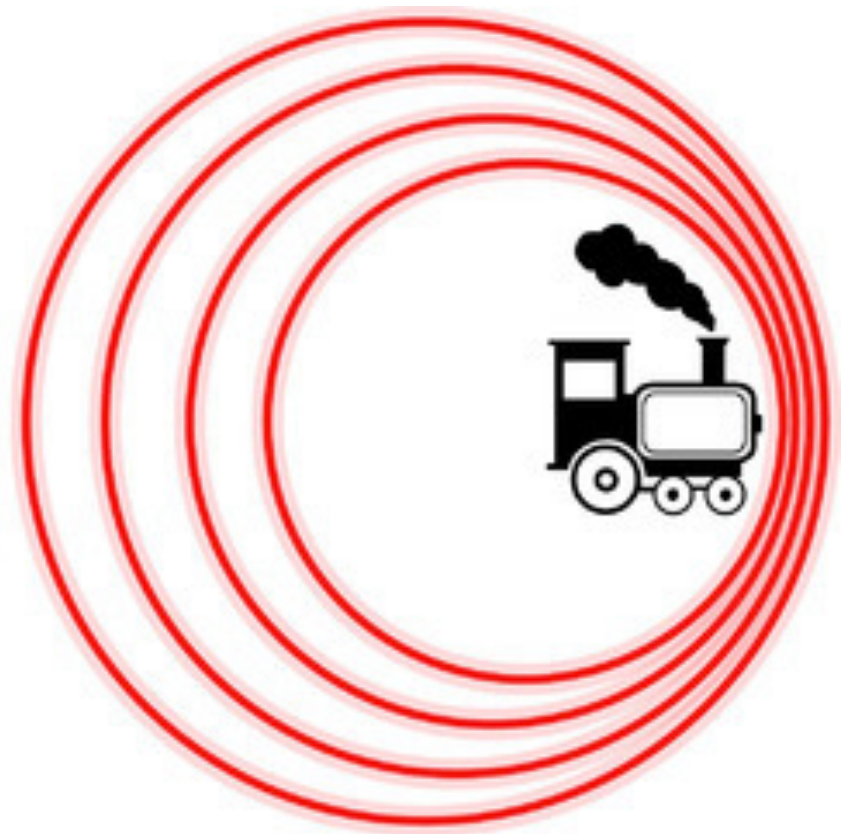
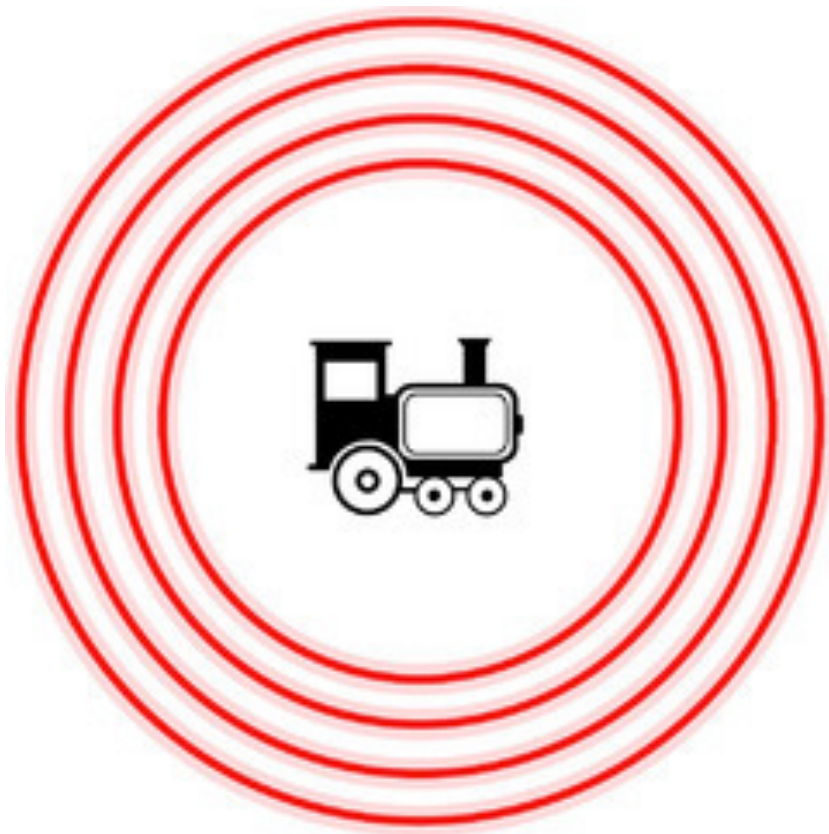
$$\nabla \times \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} + \mu_0 \mathbf{J}$$



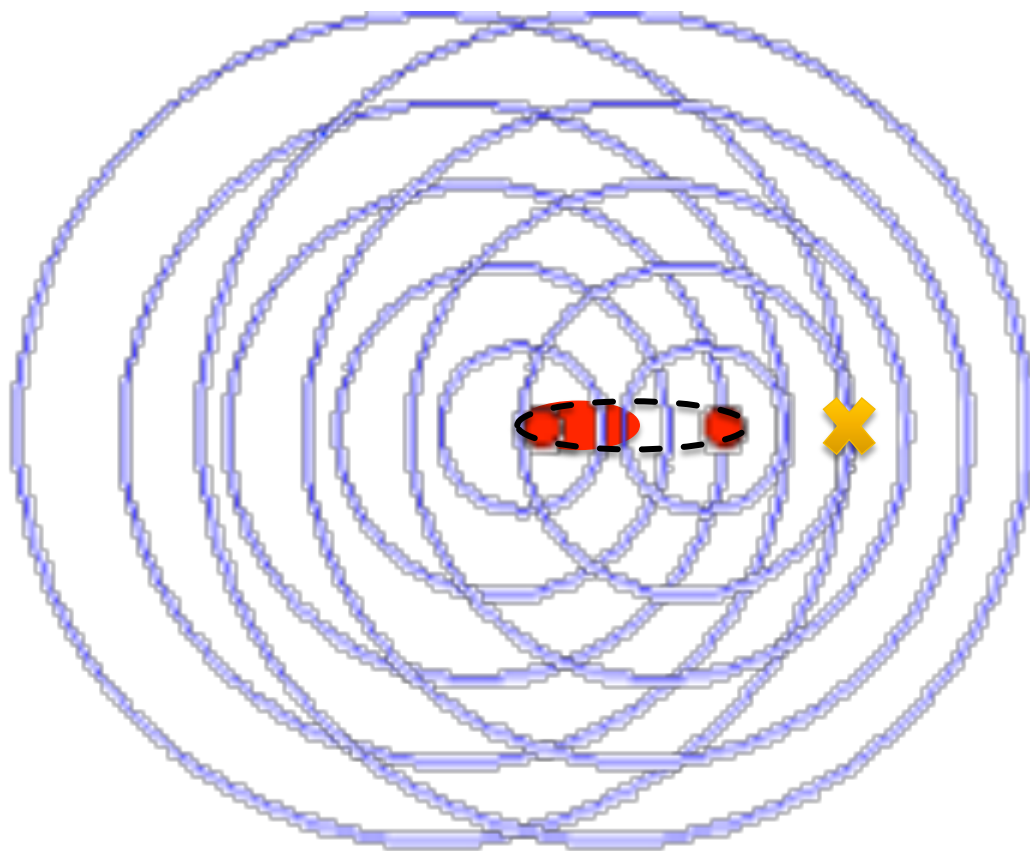
Electricity and Magnetism II: 3812

Professor Jasper Halekas
Van Allen 70
MWF 9:30-10:20 Lecture

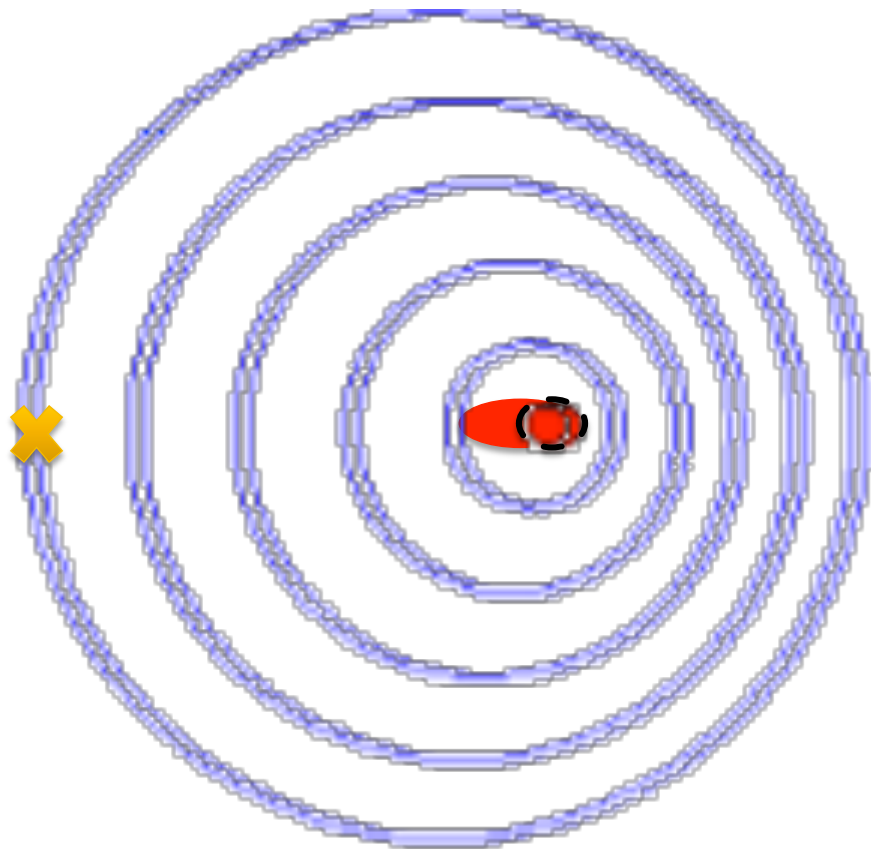
Classical Doppler Shift



Object Moving Toward



Object Moving Away

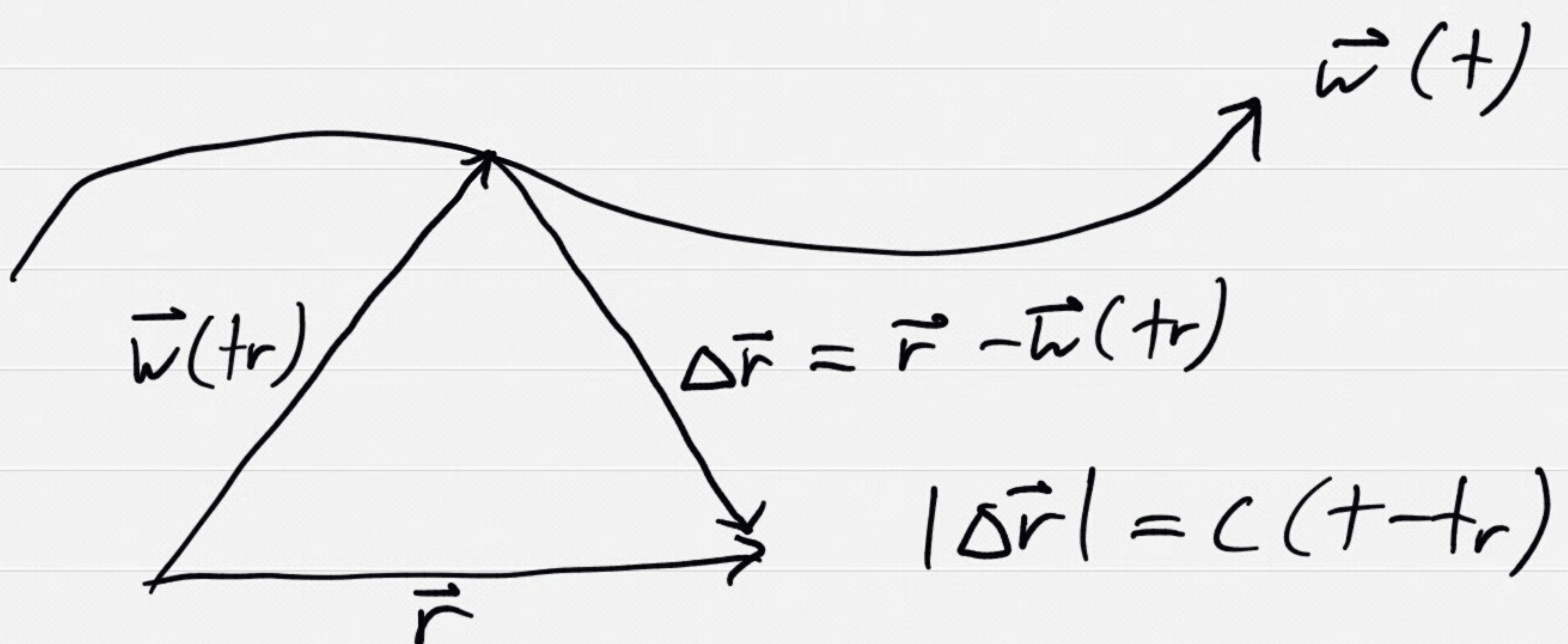


10.3 | Point Charges

Liénard-Wiechert Potentials

- Point charge moving along specified trajectory $\vec{w}(t)$

- What are $V(\vec{r}, t)$, $\vec{A}(\vec{r}, t)$?



- Only one point $\vec{w}(t_r)$ contributes to $V(\vec{r}, t)$

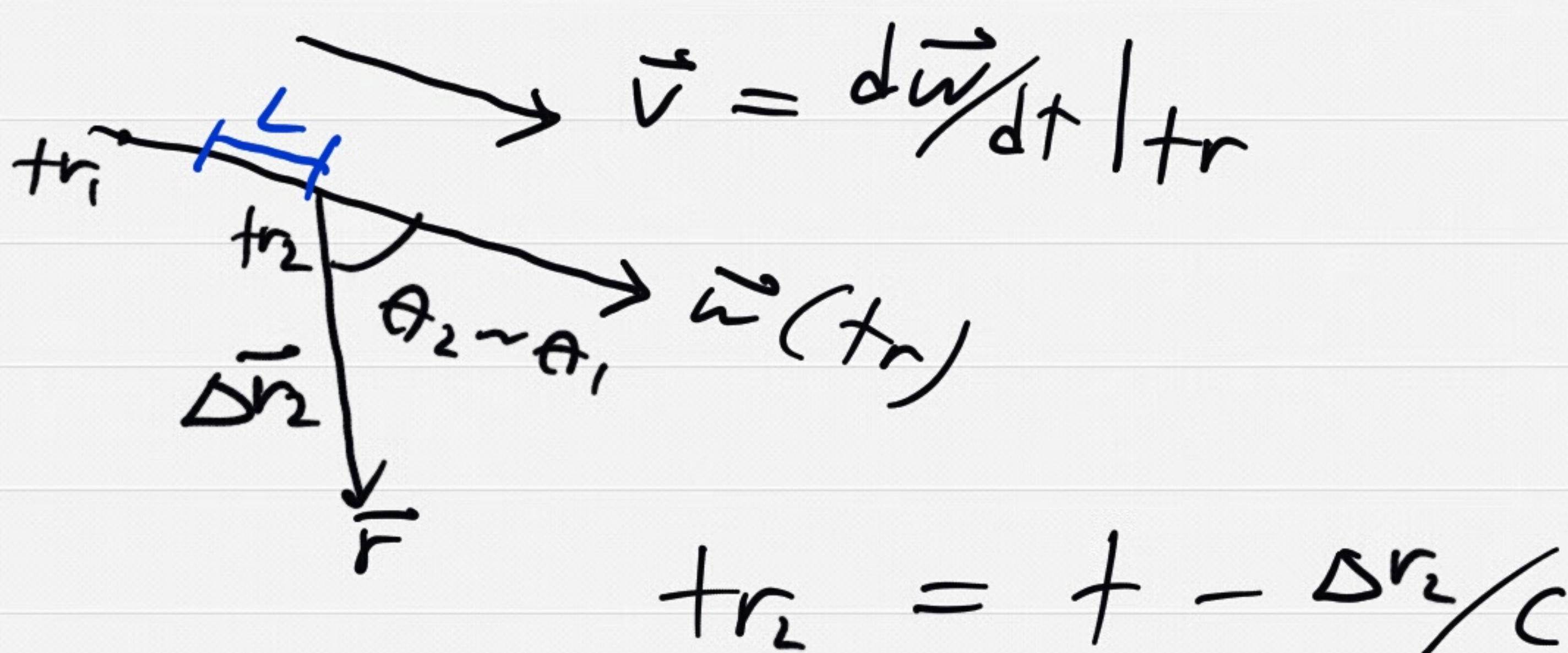
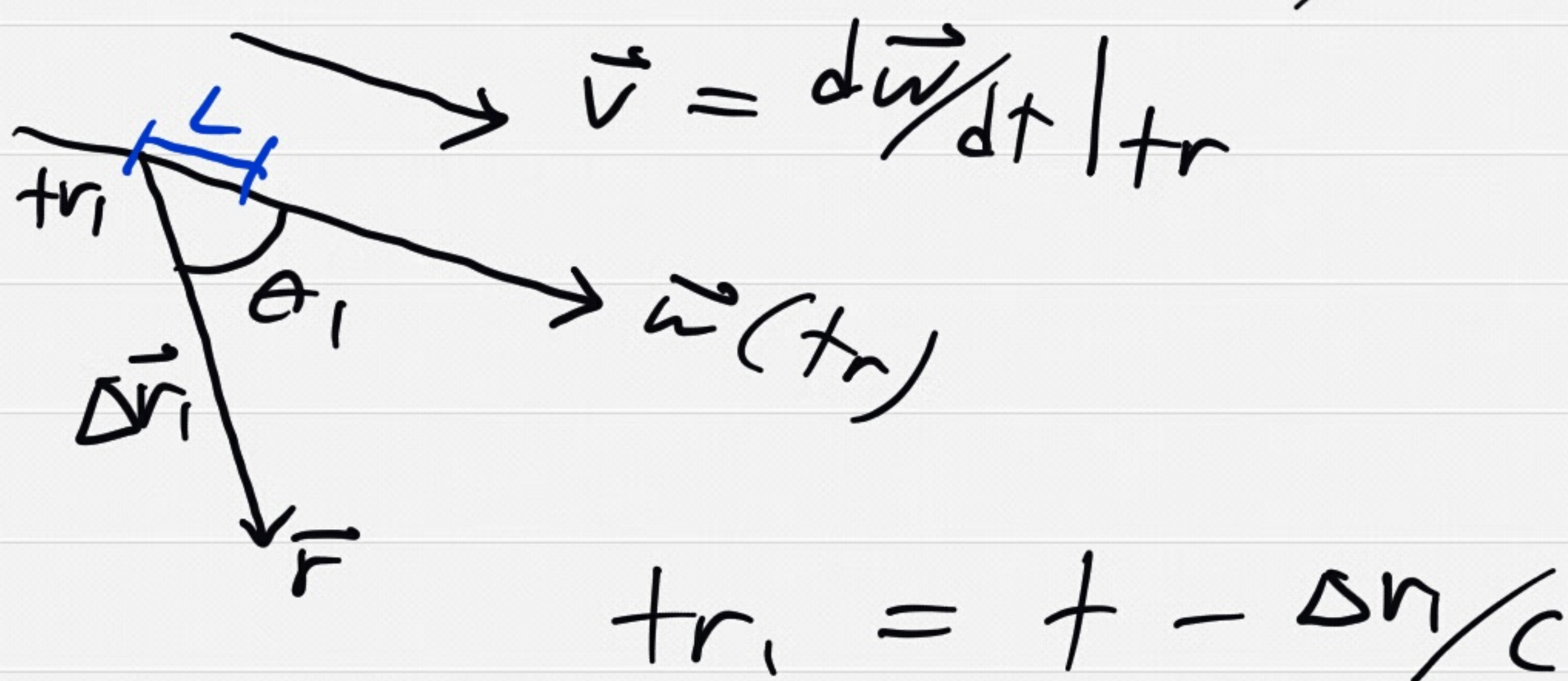
$$\begin{aligned} V(\vec{r}, t) &= \frac{1}{4\pi\epsilon_0} \int \frac{e(\vec{r}', t_r)}{\Delta r} d\tau' \\ &= \frac{1}{4\pi\epsilon_0 \Delta r} \int e(\vec{r}', t_r) d\tau' \end{aligned}$$

- You might think

$$\int \rho(\vec{r}', t_r) d\tau' = q$$

- But because the charge is moving it is effectively stretched (or compressed) in terms of its effects on an observer

Volume Stretching



Apparent length $L' = |\Delta \vec{r}_2 - \Delta \vec{r}_1|$

Note: This is a non-relativistic effect!

$$\begin{aligned}
 L' &= L + v \cos \theta (tr_2 - tr_1) \\
 &= L + v \cos \theta \left(\frac{|\Delta \vec{r}_2 - \Delta \vec{r}_1|}{c} \right) \\
 &= L + \frac{v}{c} \cos \theta L'
 \end{aligned}$$

$$\Rightarrow L' = \frac{L}{\left(1 - \frac{v}{c} \cos \theta\right)}$$

since $\cos \theta = \hat{\Delta r} \cdot \hat{v}$

$$L' = \frac{L}{\left(1 - \frac{\hat{\Delta r} \cdot \vec{v}}{c}\right)}$$

Transverse dimensions not stretched

$$\text{So } d\tau' = \frac{d\tau}{\left(1 - \frac{\hat{\Delta r} \cdot \vec{v}}{c}\right)}$$

$$\Rightarrow \int \rho(\vec{r}', tr) d\tau'$$

$$= \frac{q}{\left(1 - \frac{\hat{\Delta r} \cdot \vec{v}}{c}\right)} = \frac{q c}{c - \hat{\Delta r} \cdot \vec{v}}$$

$$\Rightarrow V(\vec{r}, t) = \frac{1}{4\pi\epsilon_0 \Delta r} \cdot \frac{q c}{c - \hat{\Delta r} \cdot \vec{v}}$$

$$\text{or } \boxed{V(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \frac{q c}{c \Delta r - \hat{\Delta r} \cdot \vec{v}}}$$

$$\vec{v} = \frac{d\vec{w}}{dt} \Big|_{tr}$$

$$\Delta \vec{r} = \vec{r} - \vec{w}(tr)$$

$$\vec{A}(\vec{r}, t) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}', t')}{\Delta r} d\tau'$$

$$= \frac{\mu_0}{4\pi \Delta r} \int \rho(\vec{r}', t') \vec{v}(t') d\tau'$$

$$\vec{A}(\vec{r}, t) = \frac{\mu_0}{4\pi} \frac{qc\vec{v}}{c\Delta r - \Delta r \cdot \vec{v}}$$

$$\frac{\mu_0}{4\pi} / \frac{1}{4\pi\epsilon_0} = \mu_0\epsilon_0 = 1/c^2$$

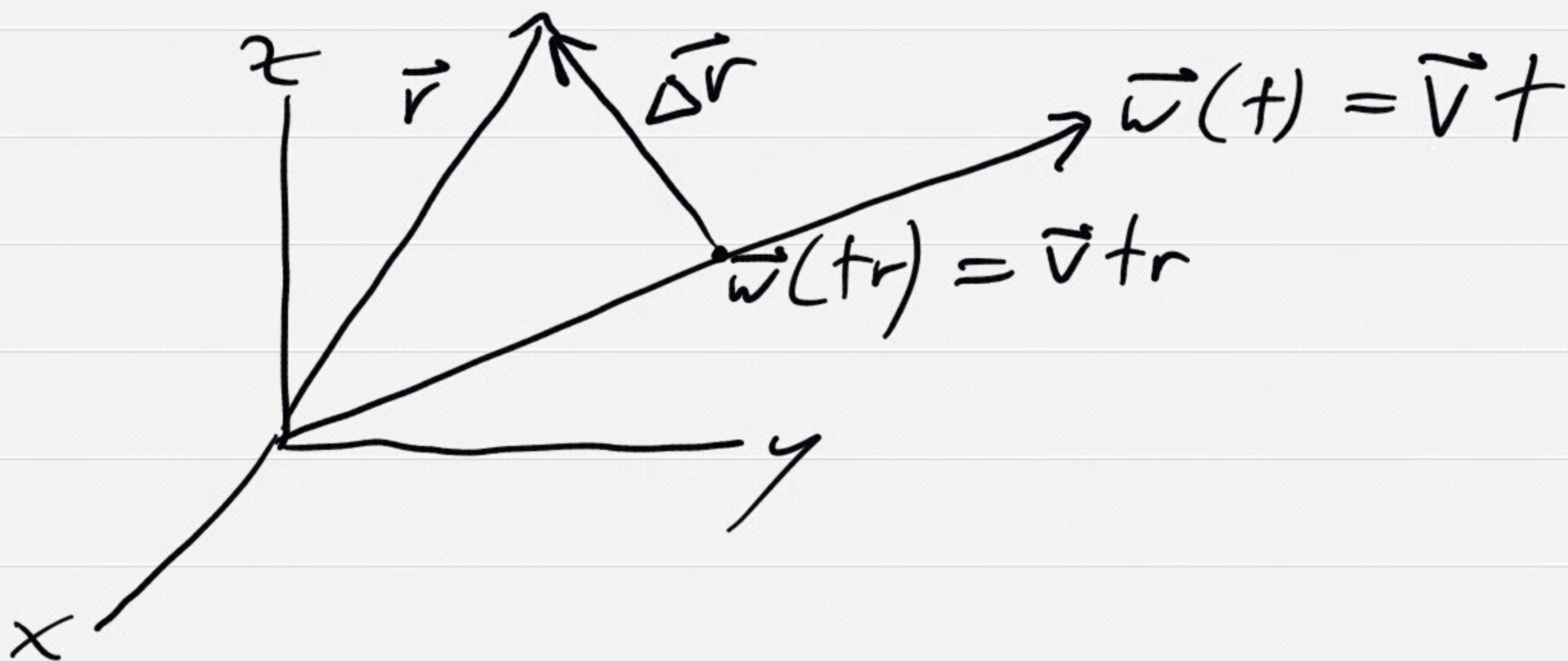
So $\vec{A}(\vec{r}, t) = \frac{1}{c^2} V(\vec{r}, t)$

For a moving point charge

Charge in Uniform Motion

$$\vec{w}(t) = \vec{v}t \quad w/ \quad \vec{v} \text{ constant}$$

$$\text{Say } \vec{w}(0) = [a, 0, 0]$$



$$|\Delta \vec{r}| = |\vec{r} - \vec{v}t_r|$$
$$= c(t - t_r)$$

$$\Rightarrow (\vec{r} - \vec{v}t_r) \cdot (\vec{r} - \vec{v}t_r) = c^2(t - t_r)^2$$

$$\Rightarrow r^2 + v^2 t_r^2 - 2\vec{r} \cdot \vec{v}t_r = c^2 t^2 + c^2 t_r^2 - 2c^2 t t_r$$

quadratic formula

$$\Rightarrow t_r = \frac{(c^2 t - \vec{r} \cdot \vec{v}) - \sqrt{(c^2 t - \vec{r} \cdot \vec{v})^2 + (c^2 - v^2)(r^2 - c^2 t^2)}}{c^2 - v^2}$$

$$\Delta r = |\Delta \vec{r}| = c(t - t_r)$$

$$\hat{\Delta \vec{r}} = \frac{\Delta \vec{r}}{\Delta r} = \frac{\vec{r} - \vec{v}t_r}{c(t - t_r)}$$

$$\Rightarrow \Delta r \left(1 - \frac{\Delta \vec{r} \cdot \vec{v}}{c}\right)$$

$$= c(t - t_r) \left[1 - \frac{v}{c} \cdot \frac{\vec{r} - \vec{v} t_r}{c(t - t_r)}\right]$$

$$= c(t - t_r) - \frac{\vec{v} \cdot \vec{r}}{c} + \frac{v^2}{c} t_r$$

$$= \frac{1}{c} \left[(c^2 t - \vec{r} \cdot \vec{v}) - (c^2 - v^2) t_r \right]$$

$$= \frac{1}{c} \sqrt{(c^2 t - \vec{r} \cdot \vec{v})^2 + (c^2 - v^2)(r^2 - c^2 t^2)}$$

$$V(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \frac{q c}{c \left(\Delta r \left(1 - \frac{\Delta \vec{r} \cdot \vec{v}}{c}\right) \right)}$$

$$\Rightarrow V(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \frac{q c}{\sqrt{(c^2 t - \vec{r} \cdot \vec{v})^2 + (c^2 - v^2)(r^2 - c^2 t^2)}}$$

$$\& \vec{A}(\vec{r}, t) = \frac{\mu_0}{4\pi} \frac{q c \vec{v}}{\sqrt{(c^2 t - \vec{r} \cdot \vec{v})^2 + (c^2 - v^2)(r^2 - c^2 t^2)}}$$

Yikes!

Low-Velocity Limit

$$(c^2 + \vec{r} \cdot \vec{v})^2 + (c^2 - v^2)(r^2 - c^2 + 2)$$

$$\sim c^4 + 2 - 2c^2 + \vec{r} \cdot \vec{v} + c^2 r^2 - c^4 + 2 + v^2 c^2 + 2$$

for $v \ll c$

$$= c^2 (v^2 + 2 + r^2 - 2 + \vec{r} \cdot \vec{v})$$

$$= c^2 |\vec{r} - \vec{v}t|^2$$

$$\Rightarrow V(\vec{r}, t) \sim \frac{qc}{4\pi\epsilon_0} \frac{1}{c|\vec{r} - \vec{v}t|}$$

$$= \frac{q}{4\pi\epsilon_0 R}$$

w/ R distance from
current (not retarded)
position of charge