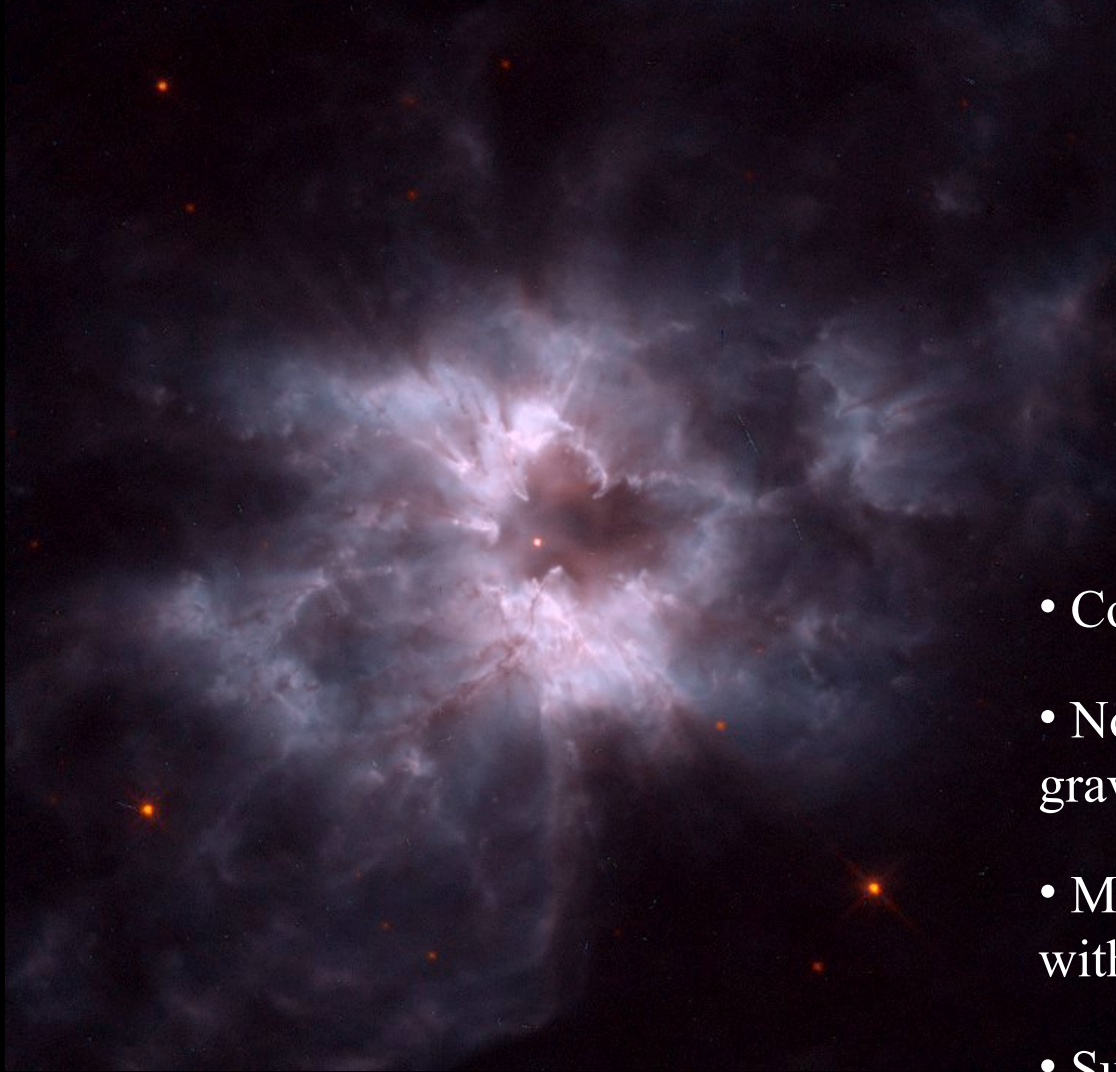


# Really, what universe do we live in?

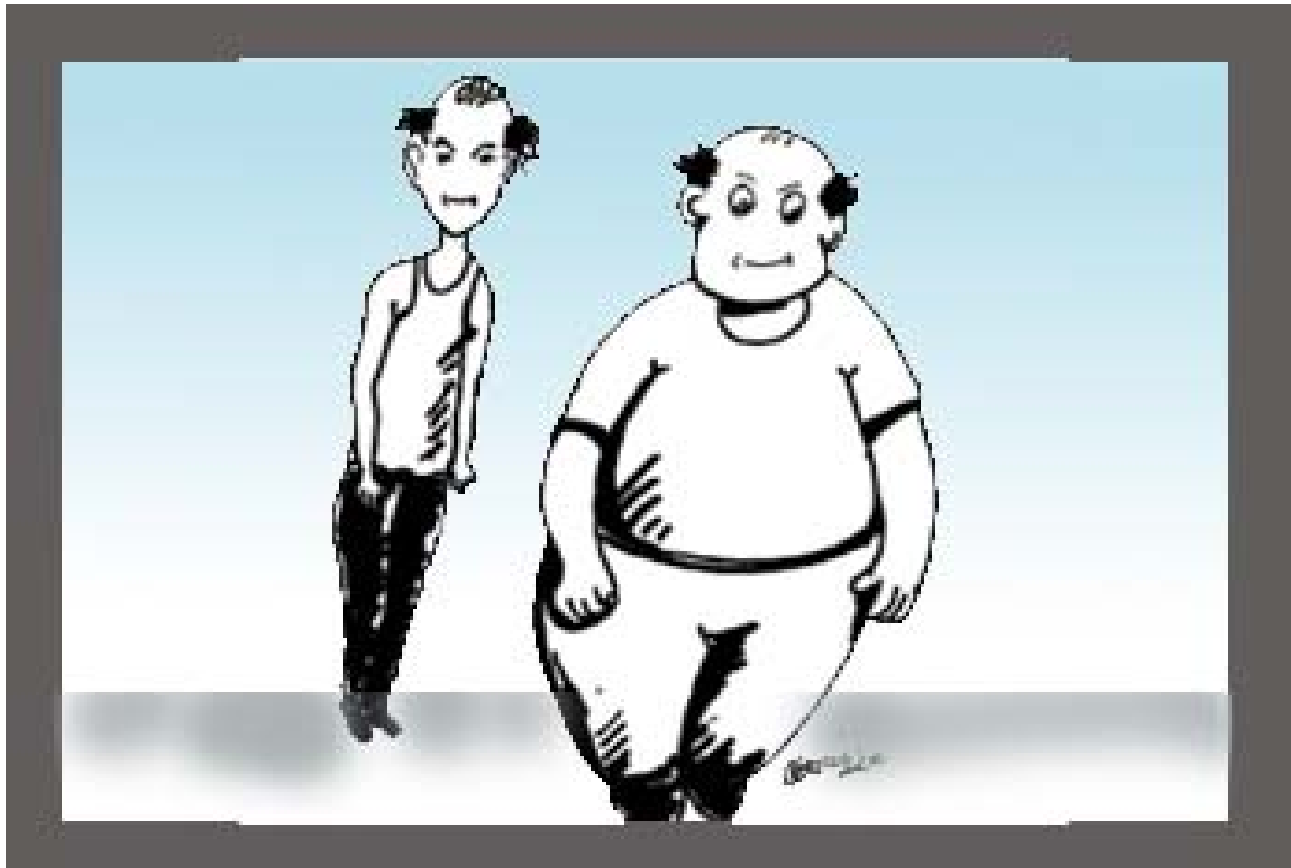
- White dwarfs
- Supernova type Ia
- Accelerating universe
- Cosmic shear
- Lyman  $\alpha$  forest

# White dwarf



- Core of solar mass star
- No energy from fusion or gravitational contraction
- Most oxygen and carbon nuclei with degenerate gas of electrons
- Supported by quantum mechanical motions of electrons

# Mass versus radius relation



- For objects made of normal matter, radius tends to increase with mass

# Mass/radius relation for degenerate star

- Star mass =  $M$ , radius =  $R$

- Gravitational potential energy =  $-\frac{3GM^2}{5R}$

- Heisenberg uncertainty:  $\Delta x \Delta p \geq \hbar$

- Electron density  $n = \frac{3N}{4\pi R^3} \approx \frac{M}{m_p R^3}$

$$\Delta x \approx n^{-1/3} \quad \Delta p \approx \frac{\hbar}{\Delta x} \approx \hbar n^{1/3}$$

- Kinetic energy  $\varepsilon = \frac{p^2}{2m_e}$   $K = N\varepsilon = \frac{M}{m_p} \varepsilon \approx \frac{\hbar^2 M^{5/3}}{m_e m_p^{5/3} R^2}$

# Mass/radius relation for degenerate star

- Total energy  $E = K + U \approx \frac{\hbar^2 M^{5/3}}{m_e m_p^{5/3} R^2} - \frac{GM^2}{R}$

- Find  $R$  by minimizing  $E$

$$\frac{dE}{dR} \approx -\frac{\hbar^2 M^{5/3}}{m_e m_p^{5/3} R^3} + \frac{GM^2}{R^2} = 0$$

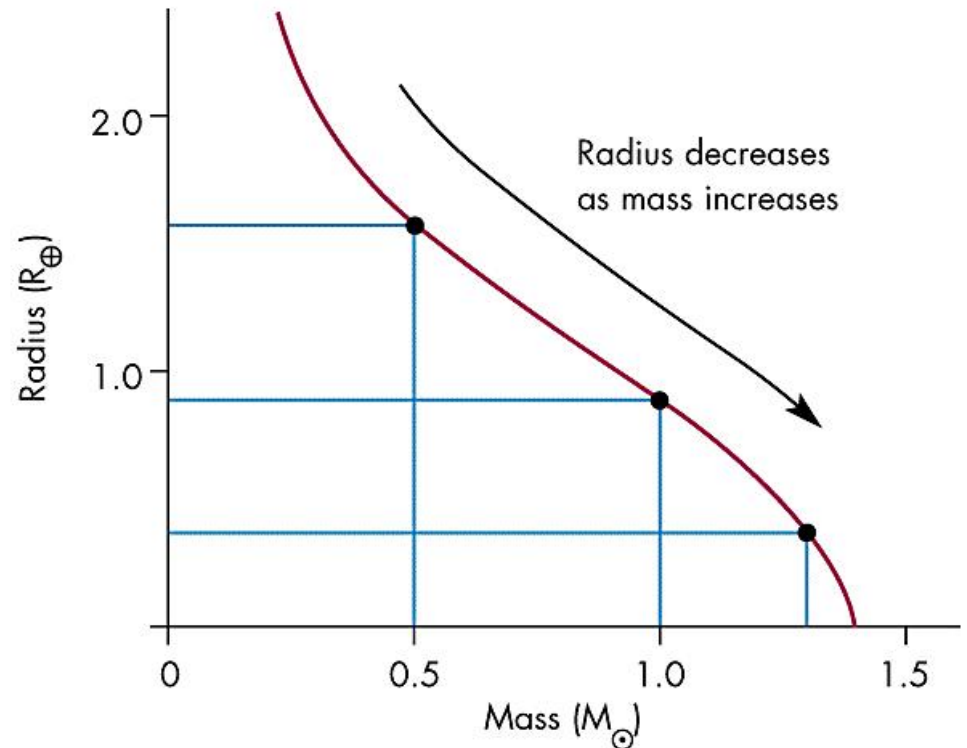
$$R \approx \frac{\hbar^2 M^{-1/3}}{Gm_e m_p^{5/3}}$$

- Radius decreases as mass increases

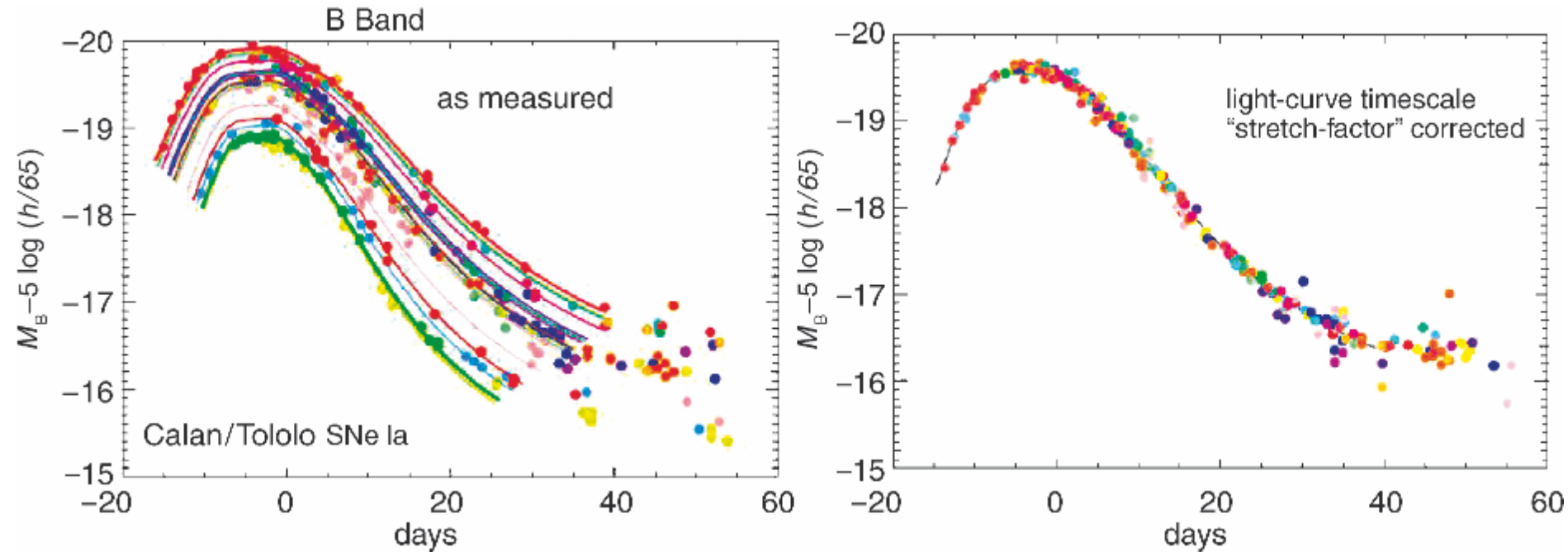
# Maximum white dwarf mass

- As mass increases, electron speed  $\rightarrow c$ , kinetic energy  $\rightarrow K=pc$
- Electron degeneracy cannot support a white dwarf heavier than 1.4 solar masses, the “Chandrasekhar limit”.

$$M = K \left( \frac{\hbar^3 c^3}{G^3 m_p^4} \right)^{1/2}$$



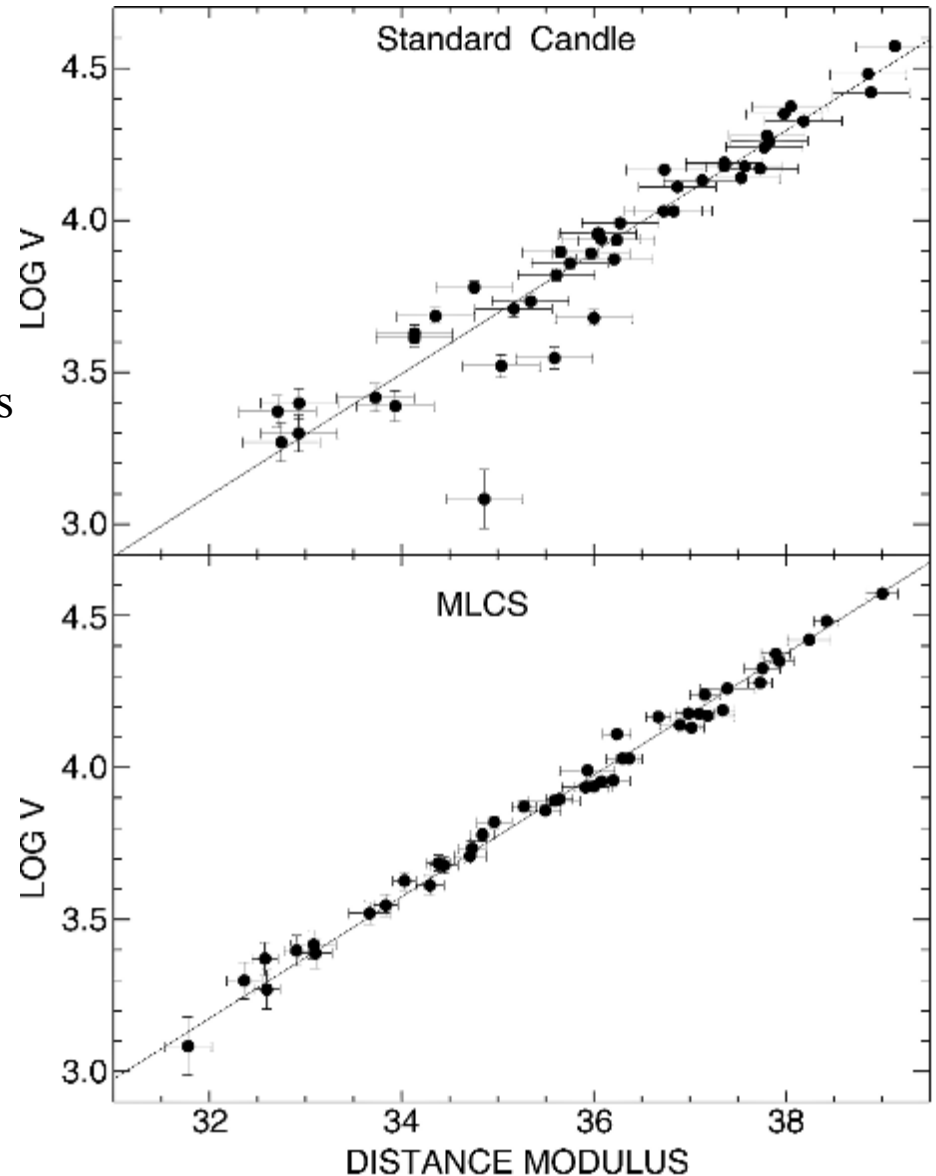
# Supernovae type Ia



- Type Ia supernovae occur when the mass of a white dwarf exceeds the Chandrasakar limit (1.4  $M_{\text{sun}}$ ). May occur via accretion from a companion or collision with another star.
- Because starting condition is the same, reasonable that all Ia have same luminosity.
- Maximum luminosity does vary, but is well correlated with time scale of light curve.
- Can correct using measured light curves and recover standard luminosity. Note need to first correct for cosmological time dilation =  $1+z$ .
- Also need to correct for extinction in host and MW galaxy.

# SN Ia as distance indicators

- Uncorrected maximum luminosity of SNIa has scatter of  $\sim 0.42$  mag.
- Correction improves this to 0.15 mag.
- Luminosity estimate is good to 15%.
- Distance estimates good to 8% if extinction is measured accurately.
- Distance modulus  $\mu = 5 \log(d/\text{pc}) - 5$



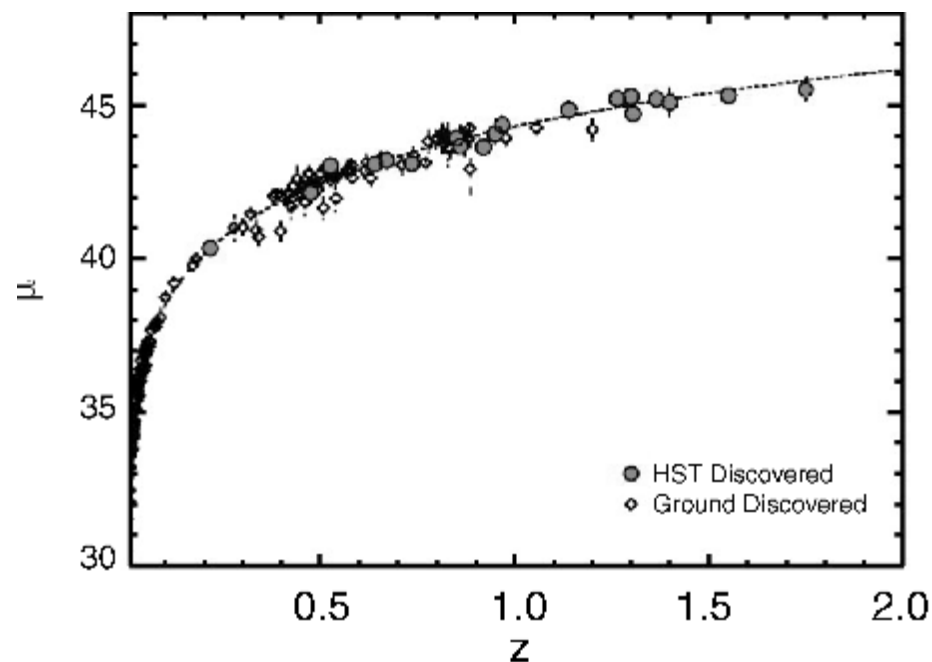
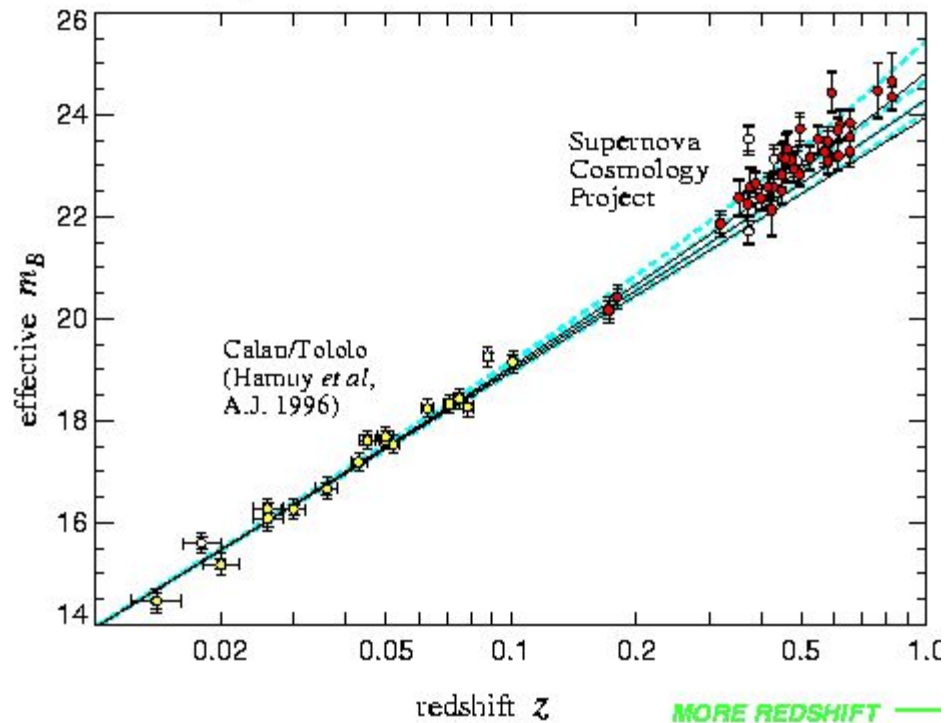


# Searching for supernovae

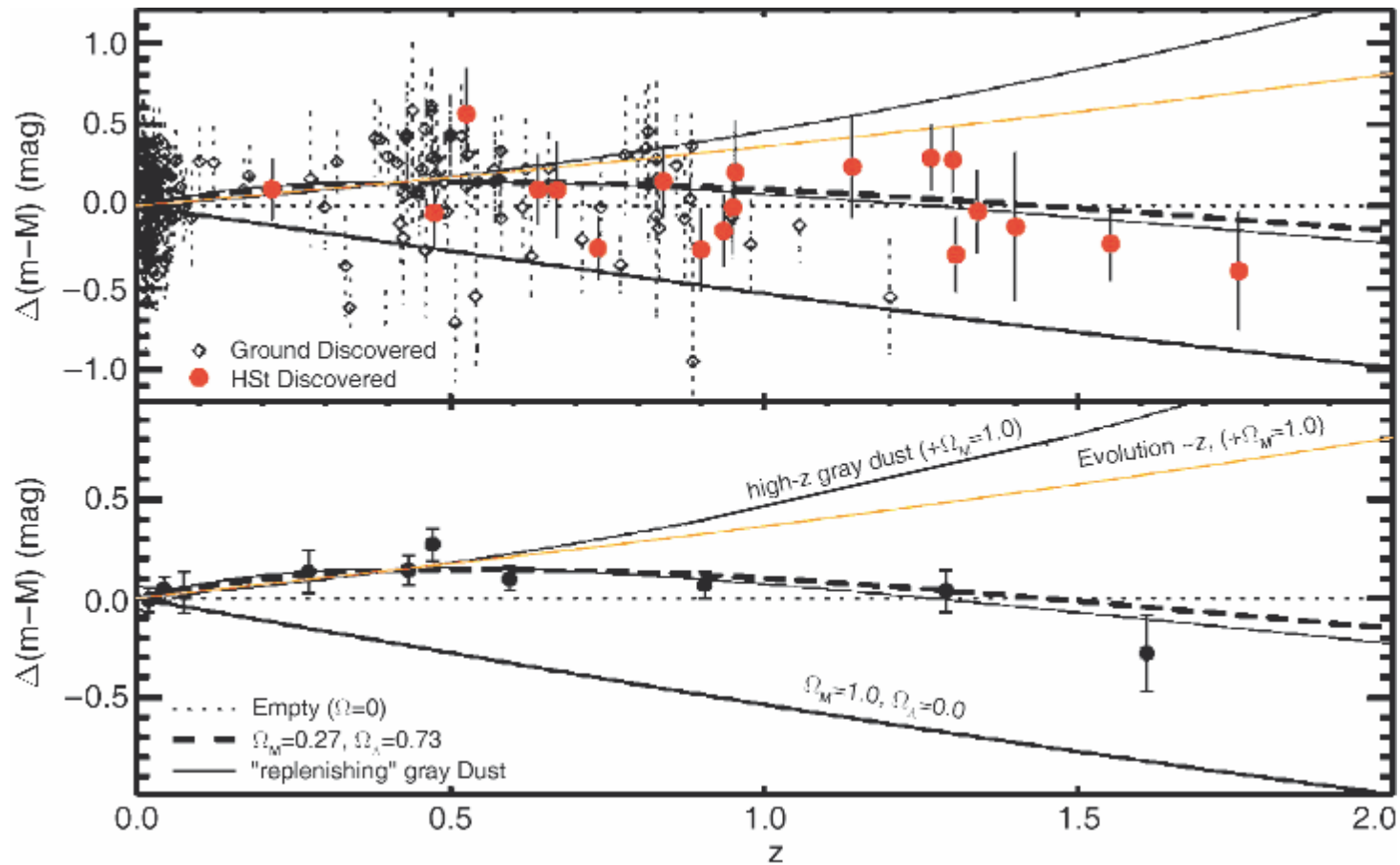
- Need to repeatedly obtain images of target galaxies or fields.
- Search for new sources appearing.
- Cadence of monitoring is set by time scale of SNIa light curves,  $\sim 10$  days.
- Then need to:
  - follow each SN with  $\sim$ daily photometry, preferably in several bands
  - Obtain spectrum to confirm Ia identification and measure redshift
- Originally, SN searches done by the Berkeley group had a very hard time in arranging telescope time for follow-up because they would need to bump observers off big telescopes (target of opportunity observations = TOOs).
- Important threshold was when search observations could discover a new SN every few nights, so there was always at least one SN to observe on any given night.
- Searches have also been done with the Hubble Space Telescope that found SN at  $z \sim 1.7$ .

# Supernova cosmology

- Plot of distance vs redshift for SNIa is a Hubble diagram, extending to high  $z$ .
- Diagrams probe expansion history of the universe.



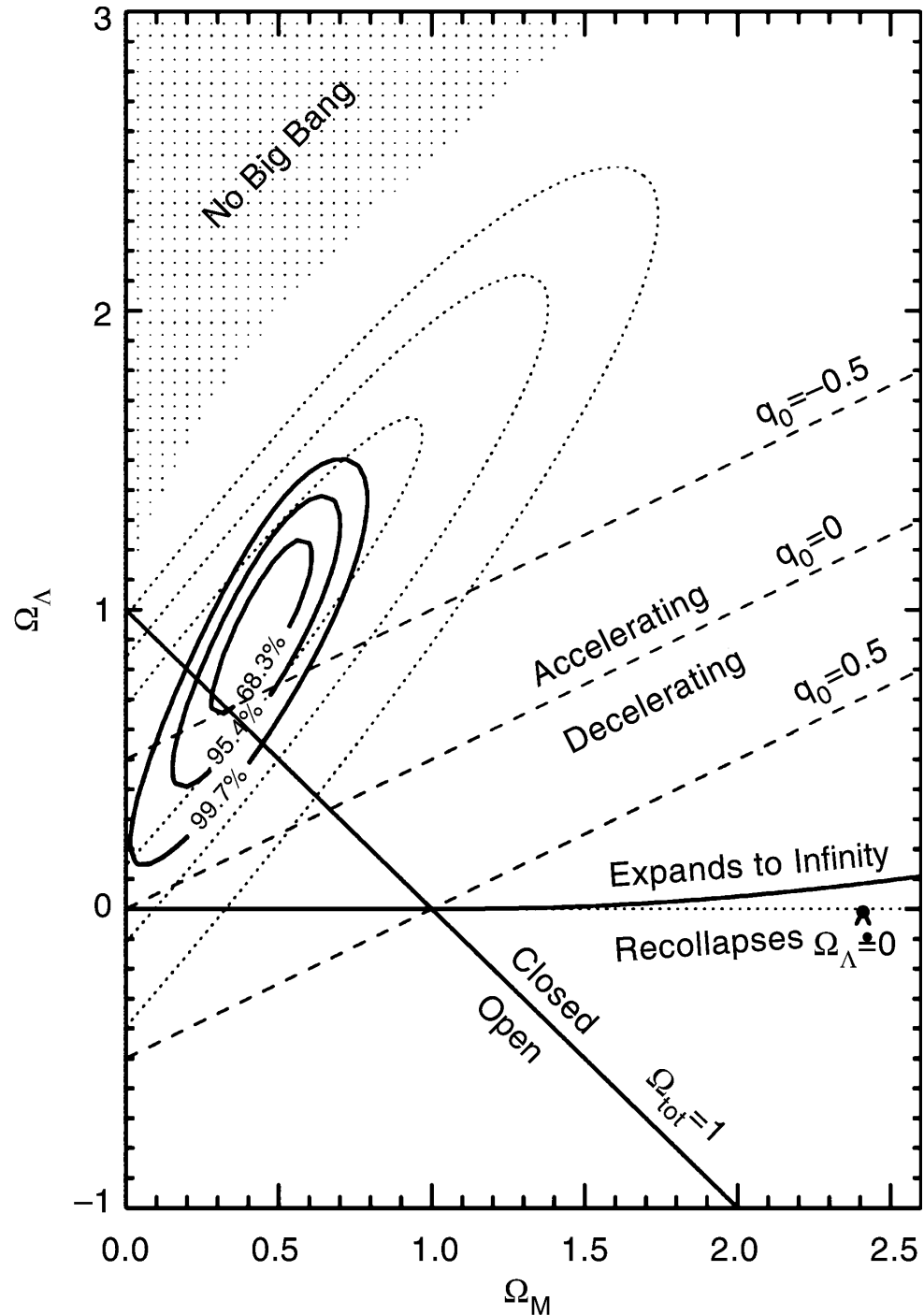
# Supernova cosmology



- Plot of mag – mag in empty universe. Empty universe means constant expansion (dotted).
- Points above dotted line mean that expansion has accelerated, e.g.  $\Omega_\Lambda > 0$ .
- Thick dashed curve  $\Omega_m = 0.27, \Omega_\Lambda = 0.73$ . Lower solid curve  $\Omega_m = 1, \Omega_\Lambda = 0$ .
- Two upper curves, evolution of SN or grey dust with  $z$ .
- High  $z$  points show turn over, important for confirming cosmological nature.

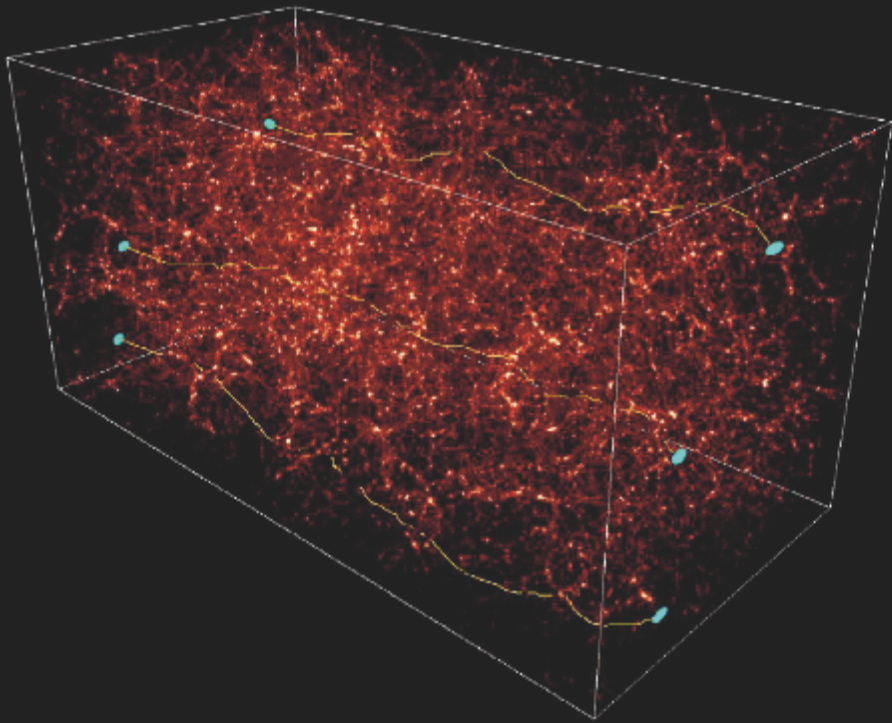
# Supernova cosmology

- Solid contours from data on previous slide. Dotted contours from 1998 data. expansion has accelerated, e.g.  $\Omega_\Lambda > 0$ .
- If we require  $\Omega_{\text{tot}} = 1$  (from CMB), then best fit is  $\Omega_m = 0.27$ ,  $\Omega_\Lambda = 0.73$ .
- $\Omega_\Lambda$  may be true cosmological constant or could be a quantum field with equation of state  $P = w\rho c^2$ ,  $w = -1$  is same as CC.
- Primary goal of future dark energy surveys is to measure  $w$ .



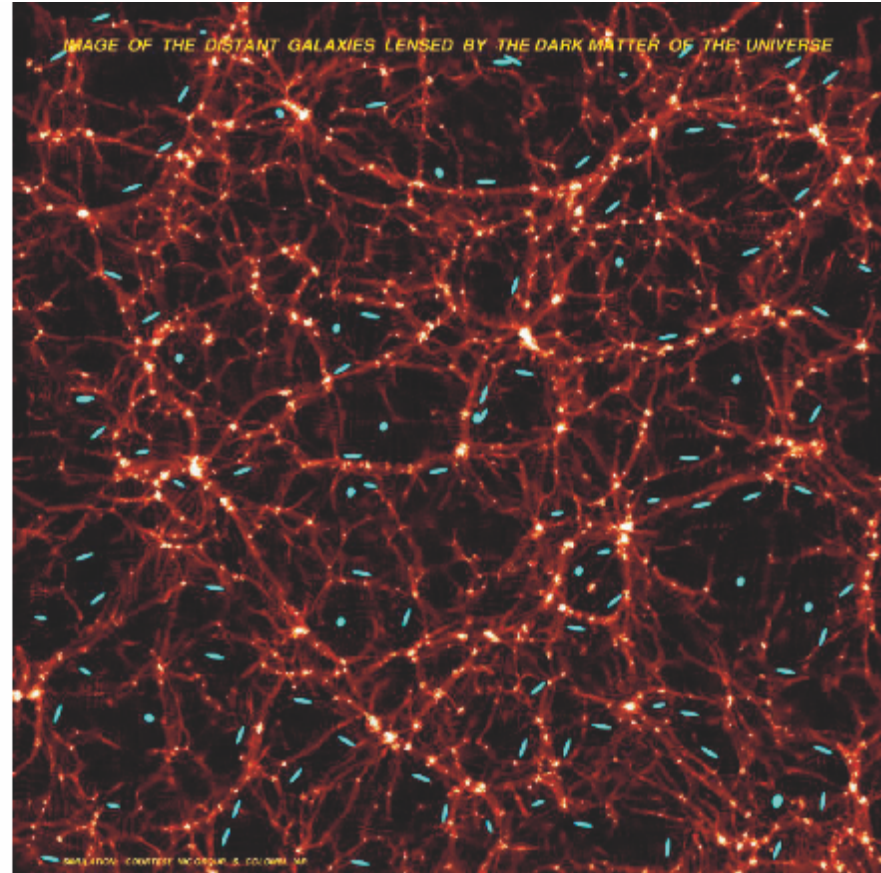
# Cosmic shear

DEFLECTION OF LIGHT RAYS CROSSING THE UNIVERSE, EMITTED BY DISTANT GALAXIES



SIMULATION: COURTESY WAC GROUP, B. COLOMBER, IAP

IMAGE OF THE DISTANT GALAXIES LENSED BY THE DARK MATTER OF THE UNIVERSE

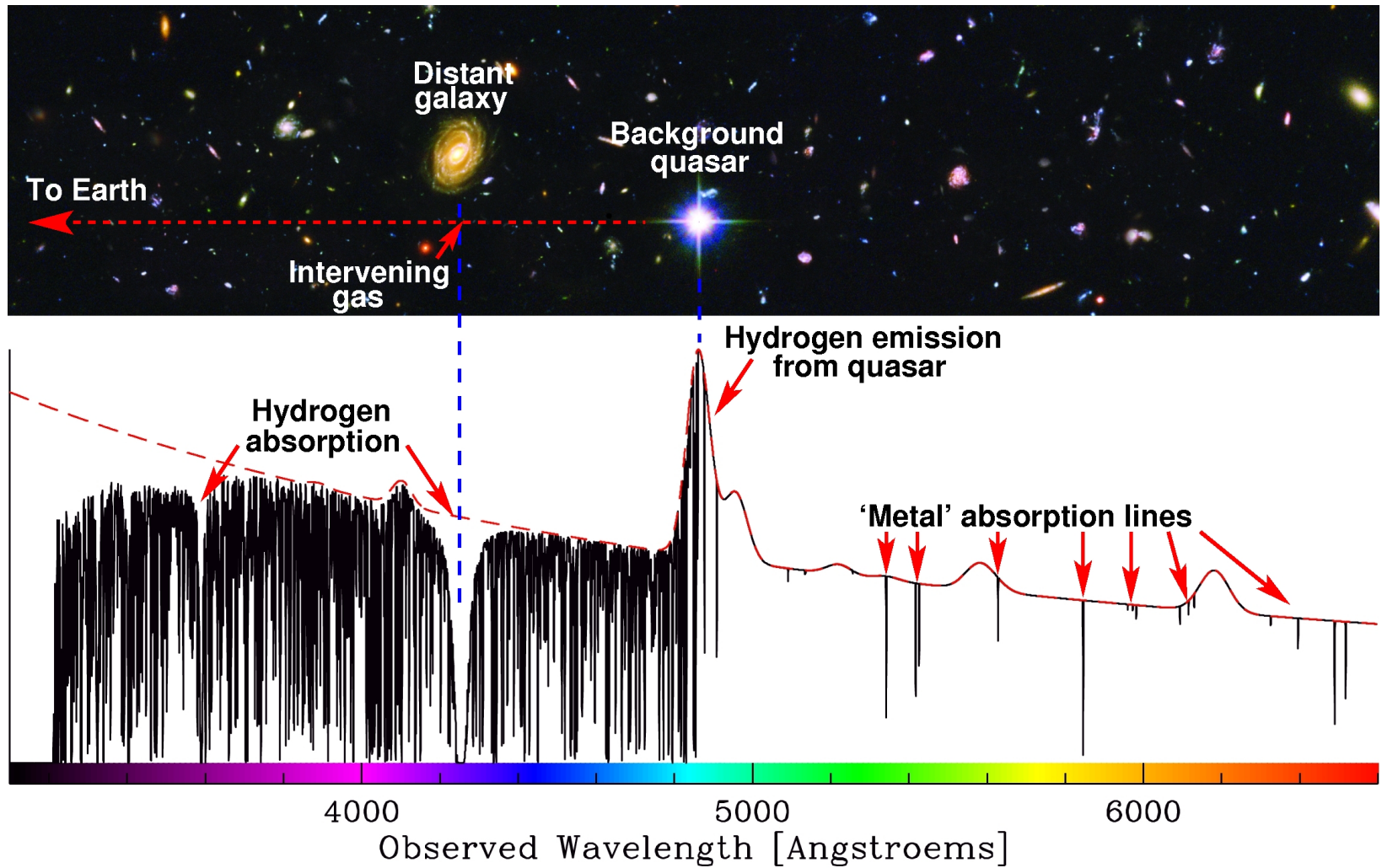


SIMULATION: COURTESY WAC GROUP, B. COLOMBER, IAP

- Shapes of galaxies are distorted as light moves through gravitational potentials to us.
- Similar to weak lensing by clusters, one can use the distortions (averaged over many galaxies) to directly probe the total matter distribution.
- Hard to do, really want redshifts for all galaxies.



# QSO absorption lines

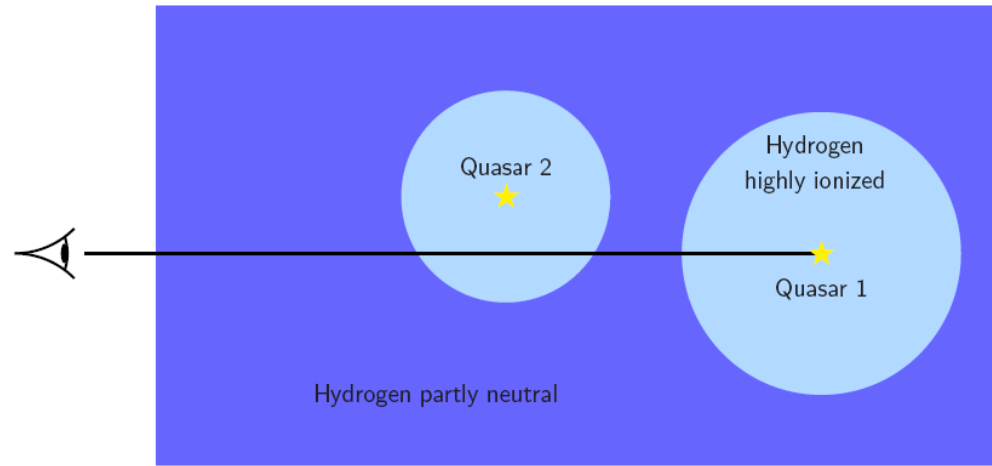


- Gas within quasars or in gas clouds or galaxies between us the and quasar can produce absorption lines.

# Diffuse intergalactic medium

- If diffuse IGM is neutral, it should absorb light shortward of Ly $\alpha$ ,  $\lambda < 1216 \text{ \AA}$ .
- Expect jump in continuum radiation from quasars between the red and blue sides of their Ly $\alpha$  line,  $S(\text{blue})/S(\text{red}) = e^{-\tau}$  where  $\tau \sim 6 \times 10^{10} (n_{\text{HI}}/\text{cm}^{-3})$  at  $z = 0$ .
- If  $\Omega_{\text{HI}} \geq 10^{-6}$  then  $\tau > 1$  and universe is opaque.
- For  $z < 3$ , find  $\tau < 0.05$ ; for  $z \sim 5$ , find  $\tau < 0.1$ .
- From these limits, find  $n_{\text{HI}}(\text{comoving}) < 3 \times 10^{-13} \text{ cm}^{-3}$  or  $\Omega_{\text{HI}} < 3 \times 10^{-8}$ .
- Compare with baryon density from BBN,  $\Omega_{\text{b}} = 0.04$ .
- Hydrogen in a diffuse component of the intergalactic medium must be essentially fully ionized.
- Searching for this gas in the nearby universe ( $z \sim 0$ ), figuring out when the ionization occurred (“reionization”), and explaining what caused reionization are major topics of current interest.

# Lyman $\alpha$ forest

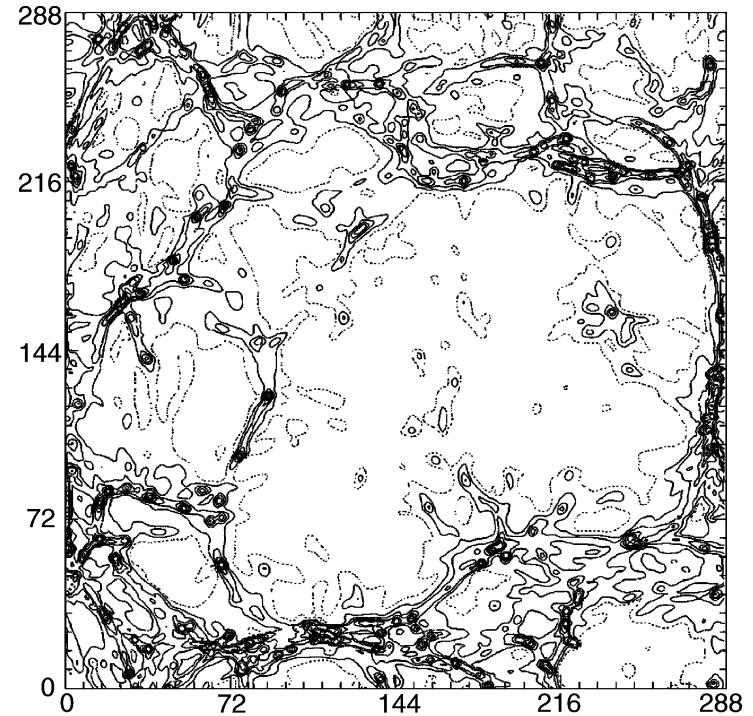


- We do see Ly $\alpha$  absorption features in spectra of distant quasars.
- Number density of Ly $\alpha$  lines for  $z > 2$  follows  $dN/dz \sim k(1+z)^\gamma$  with  $\gamma \sim 2.5$  and  $k \sim 4$ .
  - Increases strongly with  $z$ .
- Can determine HI column density from line strength.
- Widths of lines often  $\sim 10,000$  km/s, origin?
  - Massive cluster, outflow, or thermal ( $\sim 10^4$  K)
- Proximity effect: quasars ionize nearby gas, suppressing Ly $\alpha$  forest.
  - Ionization rate =  $\Gamma \cdot n_{\text{HI}} =$  recombination rate =  $\alpha n_p n_e$ , where  $\alpha = \alpha(T)$
  - In highly ionized medium  $n_{\text{HI}} \ll n_p = n_e$
  - Thus  $n_{\text{HI}} = \alpha n_p^2 / \Gamma$



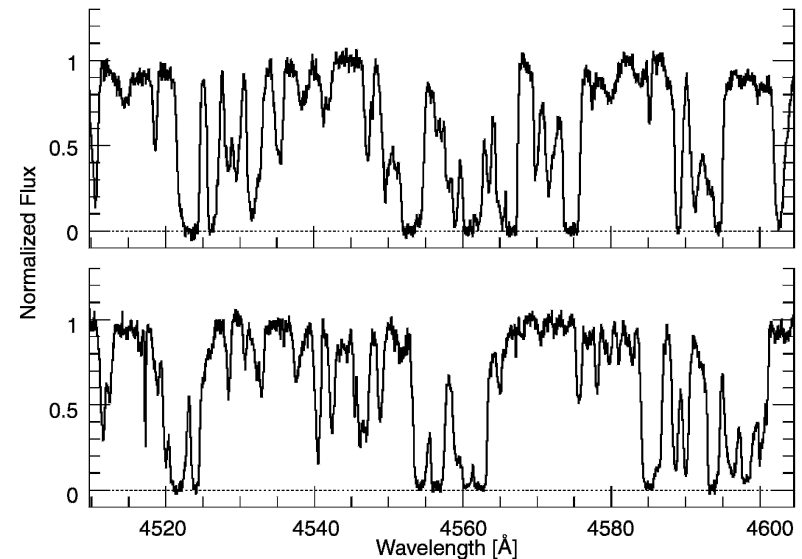
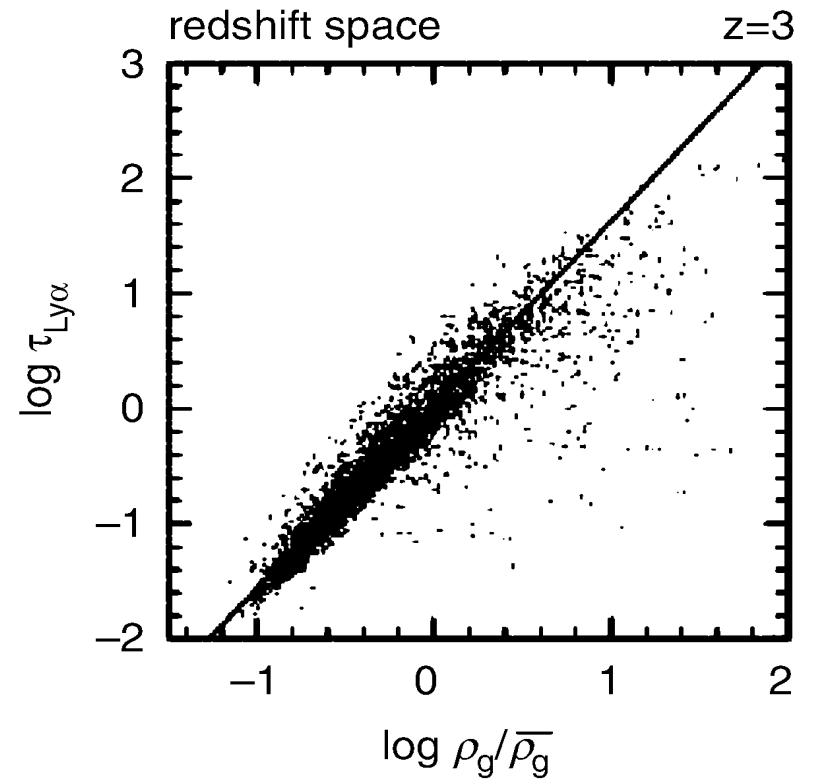
# Lyman $\alpha$ forest

- How did Ly $\alpha$  forest arise?
- Galaxies and clusters formed when baryons fell into dark matter gravitational potential wells after recombination.
- Gas clouds in the Ly $\alpha$  forest formed the same way. However, smaller density perturbations produced smaller baryon concentrations.
- These smaller perturbations are better described by linear theory.
- Ly $\alpha$  forest can be probed to higher  $z$  than galaxy and cluster distributions, up to  $z \sim 4$ .
- Can do same sorts of statistics (power spectrum, correlation function) as with galaxies but have disadvantage that probes are along a limit number of lines of sight.



# Lyman $\alpha$ forest

- Simulations follow evolution of dark matter. Must convert dark matter perturbations into gas density perturbations and then into the observable which is optical depth of Ly $\alpha$  absorption.
- As gas falls into potential well, it gains energy and heats up. Temperature depends on density.
- Use  $n_{\text{HI}} = \alpha n_p^2 / \Gamma$ , temperature dependence of  $\alpha$ .
- Find  $\tau = A(n_p / \bar{n}_p)^\beta$  with  $\beta \sim 1.6$  and  $A$  depends on  $z$ , intensity of ionizing radiation, average temperature of universe.
- This provides relation to convert simulation results into simulated spectra.
- Figure shows two spectra, one simulated, one observed.



# For next class

- Read 8.6-8.7.
- HW #8 due on November 5.
- Project presentations on November 12 and 14. Should be full presentations with all results.