#### Supernova Remnants

- Shell-type versus Crab-like
- Phases of shell-type SNR

# Shell-type SNR



#### Shell-type SNR:

X-ray, radio, and optical emission come from a shell. Xrays are usually thermal, but can have non-thermal components.

Shell is expanding.

Power source is inertia left from initial supernova. No current input of energy.

#### RCW 86, SN in 185 AD

# Pulsar Wind Nebulae (Plerions)



#### Crab, SN in 1054 AD

Center filled or Crab-like SNR, or pulsar wind nebulae:

X-ray, radio, and optical emission come from a filled, central region. X-rays are nonthermal.

Motions can be detected internal to the nebula.

Continuously powered by relativistic wind from pulsar at center of nebula.

## Mixed Morphology



Plerionic composite: shell-type on the outside, Crab-like at the center.

Thermal composite: Radio shell, center-filled X-ray emission, but X-rays are thermal. Thought to occur in denser ISM than shelltype SNR. X-rays may be due to evaporation of clouds ISM after shock front has passed.

W28: red = radio, green =  $H\alpha$ , blue = X-ray,

## Phases of Shell-type SNRs

- Supernova explosion ejecta  $v \sim 10^4$  km/s
- Free expansion ejecta mass > swept up mass
- Adiabatic or Sedov swept-up mass > eject mass
- **Snow-plow or Cooling** shock front cools, interior also cools
- **Disappearence** remnant slows to speed of the random velocities in the surrounding medium, merges with ISM

#### Shock Formation

At time t=0, mass  $m_0$  of gas is ejected with velocity  $v_0$  and total kinetic energy  $E_0$ . This interacts with surrounding interstellar material with density  $\rho_0$  and low T.



Shock front, ahead of 'heated' material

> The shell velocity much higher than the sound speed in ISM, so a shock front of radius R forms.

### Shock Jump

 $v_1$   $v_0$  Look at reference frame where downstream upstream shock is stationary,  $v_0$  = shock speed

Mass flux:  $\rho_1 v_1 = \rho_0 v_0$  Momentum flux:  $P_1 + \rho_1 v_1^2 = P_0 + \rho_0 v_0^2$ Energy flux:  $\frac{1}{2}\rho_1 v_1^3 + Pv_1 \gamma/(\gamma - 1) = \frac{1}{2}\rho_0 v_0^3 + Pv_0 \gamma/(\gamma - 1)$ 

Where  $\rho$  is density, *P* is pressure,  $\gamma$  is the adiabatic index.

Introduce the Mach number  $M = v_0/c_0$  where  $c_0 = \operatorname{sqrt}(\gamma P_0/\rho_0)$  is the sound speed upstream, and find in the limit of large M

$$\rho_1 / \rho_0 = (\gamma + 1) / (\gamma - 1)$$
 and  $T_1 / T_0 = 2\gamma(\gamma - 1) M^2 / (\gamma + 1)^2$ 

For  $\gamma = 5/3$ , find  $\rho_1 / \rho_0 = 4$  and  $T_1 / T_0 = (5/16)M^2$ 

Get large increase in temperature for large M.

## Free Expansion

- Shell of swept-up material in front of shock does not represent a significant increase in mass of the system.
- ISM mass previously within the swept-up sphere of radius R is still small compared to the ejecta mass:  $(4\pi/3)\rho R^3 \ll m_0$

• Since momentum is conserved:  $m_0 v_0 = (m_0 + (4\pi/3)\rho_0 R^3) v$ 

- As long as swept-up mass << ejecta mass, the velocity of the shock front remains constant and  $R_s(t) \sim v_0 t$
- The temperature decreases due to adiabatic expansion,  $T \propto R^{-3(\gamma-1)}$

• Adiabatic process has  $TV^{\gamma-1} = \text{constant}$ 

where  $\gamma = C_p / C_v$  = ratio of specific heats

For ideal gas,  $\gamma = (f+2)/f$ 

- f = degrees of freedom
  - = 3 for monatomic gas
  - = 5 for 5 for diatomic gas
- Note  $V \propto R^3$ , so  $TV^{\gamma-1} = \text{constant} = TR^{3(\gamma-1)}$
- Hence  $T \propto R^{-3(\gamma-1)}$

#### Free Expansion



## **Dimensional Analysis**

What is the free fall time scale of the Sun?

Inputs:

- Mass =  $M = 2 \times 10^{33}$  g
- Radius =  $r = 7 \times 10^{10}$  cm
- $G = 6.7 \times 10^{-8}$  erg cm g-2 =  $6.7 \times 10^{-8}$  cm3 s-2 g-1.

Output is time = t with units of seconds,

 $T = M^{a} r^{b} G^{c}$ 

find a, b, c to make units work.

Have an explosion with energy, *E*, that goes off in the ISM of density,  $\rho_0$ . These are the only two variables in the problem.

Dynamics can be described by location of shock front versus time. Look for a solution,  $R = E^a \rho_0^b t^c$ .

Have students work out solution.

Have an explosion with energy, *E*, that goes off in the ISM of density,  $\rho_0$ . These are the only two variables in the problem.

Dynamics can be described by location of shock front versus time. From dimensional analysis,

$$R = E^{1/5} \rho_0^{-1/5} t^{2/5} = (E/\rho_0)^{1/5} t^{2/5}$$

Another way to time about this is that  $E/\rho_0$  has units of  $(\text{length})^5(\text{time})^{-2}$ . Therefore,  $(E/\rho_0)(t^2/R^5)$  is a dimensionless quantity which describes the dynamics of the expansion.

The solution requires  $R(t) = k(E/\rho_0)^{1/5} t^{2/5}$  and v(t) = 2R/5t

This solution describes the expansion of SNR pretty well.

Can you estimate the energy of this atomic bomb explosion from the photograph?



### Sedov Solution

In Sedov solution, find for downstream material:

pressure =  $(3/4) \rho_0 v^2$ 

temperature =  $(3m/16k) v^2$  where *m* is the mean mass per particle downstream (including electrons) and k is Boltzmann's constant.

Temperature ~  $(10 \text{ K})v^2$  for v in km/s,

For  $v \sim 1000$  km/s, have  $T \sim 10^7$  K which means gas is heated to X-ray producing temperatures.

Also have internal shock that propagates back into the SNR.



### N132D in the LMC

Shock speed  $\sim 2,000$  km/s.

Gas is heated by shock to X-ray emitting temperatures.

Although gas glows in X-rays, the loss of energy due to radiation is relatively unimportant to the dynamics of the expansion, i.e. cooling time is longer than age of SNR.



## Radiative Cooling

- Eventually, the shock slows down, gas is heated less. Define end of adiabatic phase as when half of energy has been radiated away. Typically, shock speed is then about 200 km/s (with dependence on initial energy and ISM density). Most material swept-up into dense, cool shell. Residual hot gas in interior emits weak X-rays.
- Matter behind shock cools quickly, pressure is no longer important, shell moves with constant momentum

 $(4\pi/3)R^{3}\rho_{0}v = \text{constant.}$ 

$$R = R_{rad} \left( \frac{8}{5} \frac{t}{t_{rad}} - \frac{3}{5} \right)^{1/4}$$

## Disappearance

- When shock velocity drop to ~20 km/s, the expansion becomes subsonic and the SNR merges with the ISM.
- However, the SNR leaves magnetic fields and cosmic rays which can still persist with observable imprints for millions of years.

## Phases of Shell-type SNRs

- Supernova explosion Fast
- Free expansion Hundreds of years
- Adiabatic or Sedov 10,000-20,000 years
- Snow-plow or Cooling Few 100,000 years
- **Disappearence** Up to millions of years