

# Dead Stars

- End of nuclear burning
- Supernova explosions
  - Formation of elements (R, S process)
- Chandrasekhar limit on white dwarf mass
- Neutron stars
- Pulsars

# Nuclear Burning in Stars

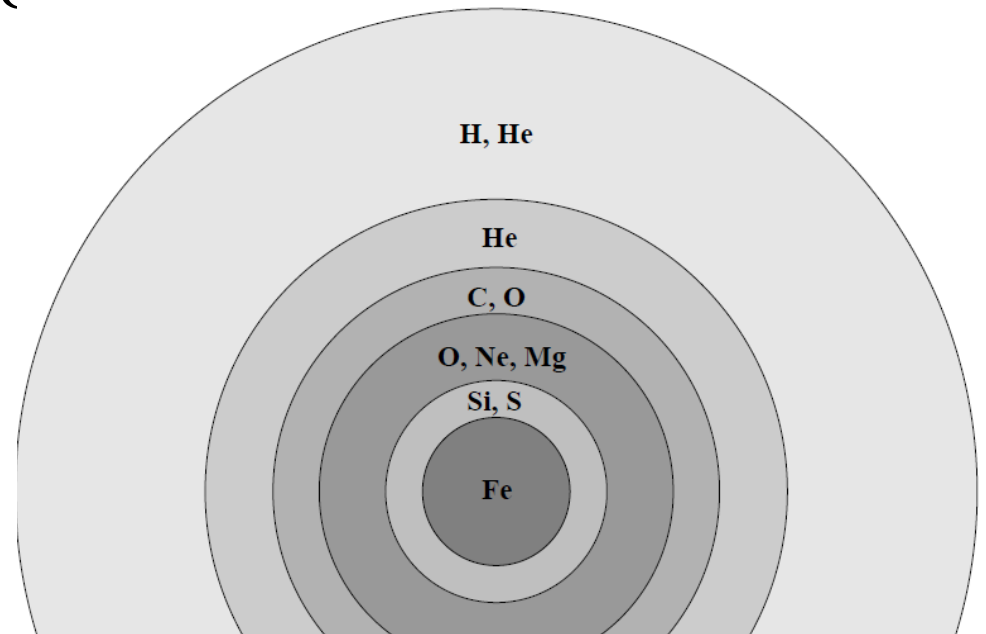
- Stars are powered by nuclear burning, which produces the energy needed to maintain pressure balance against gravity.
- Initial reaction is proton burning to He.
- At end of life for a solar mass star, final step is He burning.
- He burning:  $4\text{He} + 4\text{He} + 4\text{He} \rightarrow {}^{12}\text{C} + \gamma$  ( 7.275 MeV)
  - Reaction occurs through  $4\text{He} + 4\text{He} \rightarrow {}^8\text{Be} + \gamma$ , but  ${}^8\text{Be}$  has lifetime of only  $10^{-16}$  seconds.
- He burning also produces O and N
  - $4\text{He} + {}^{12}\text{C} \rightarrow {}^{16}\text{O} + \gamma$
  - $4\text{He} + {}^{16}\text{O} \rightarrow {}^{20}\text{N} + \gamma$
- When conditions no longer allow He burning, nuclear fusion stops. Have a white dwarf – pressure support is from “electron degeneracy pressure” instead of pressure due to thermal motions.

# End of Nuclear Burning

- There is a limit to the mass that electron degeneracy pressure can support, called the Chandrasekhar limit (1.4 solar masses).
- If gas is accreted onto a white dwarf, then this limit can be exceeded and the white dwarf collapses.
- This produces a supernova explosion.

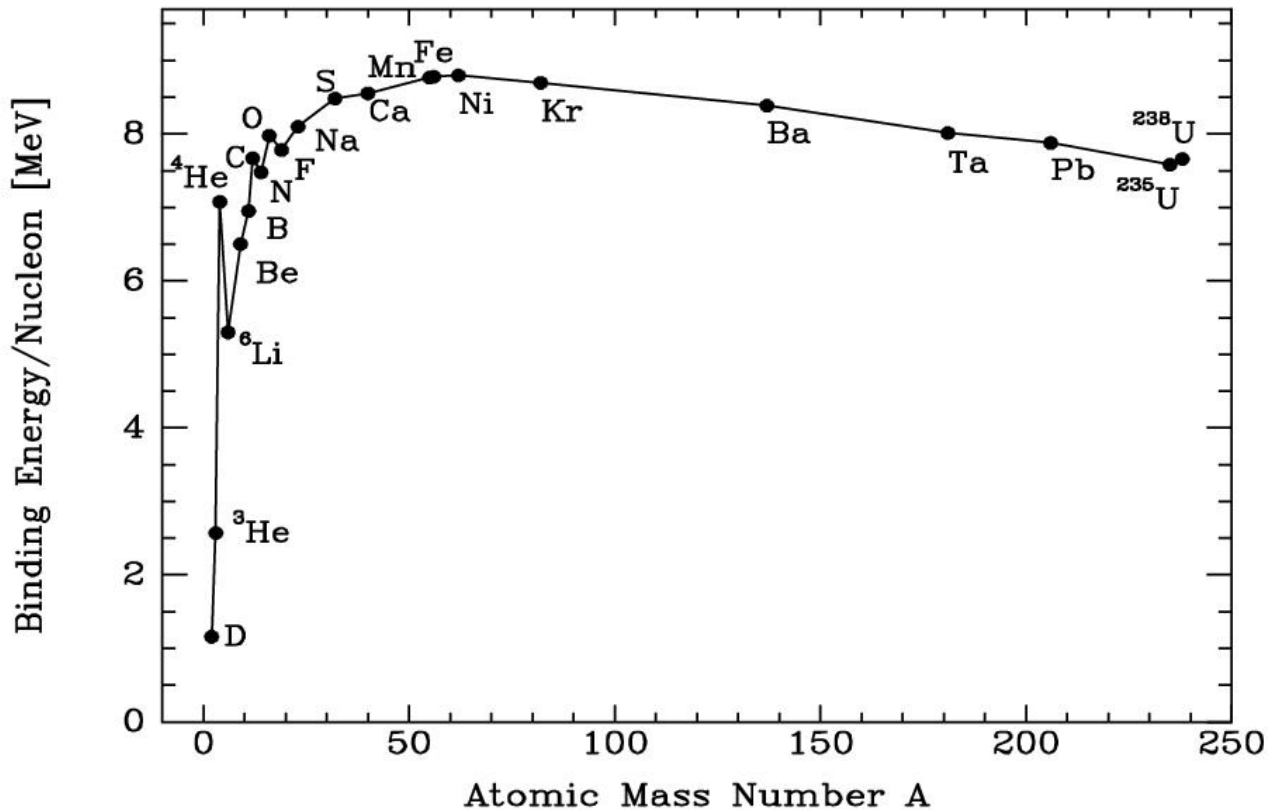
# Nuclear Burning in Massive Stars

- Core temperature and density is higher in high mass stars
  - Nuclear burning can progress past CNO.
- $^{16}\text{O} + ^4\text{He} \rightarrow ^{20}\text{Ne} + \gamma$
- $^{12}\text{C} + ^{12}\text{C}$  produces mostly  $^{20}\text{Ne}$  and  $^{24}\text{Mg}$ .
- $^{16}\text{C} + ^{16}\text{C}$  produces mostly  $^{28}\text{Si}$  and  $^{32}\text{S}$ .
- Si burns to  $^{56}\text{Fe}$  and some  $^{52}\text{Cr}$
- Leads to shell structure.
- Burning is very rapid
  - C  $\sim$  500 years
  - Si  $\sim$  1 day



# Fusion

- Burning stops releasing energy with Fe because it is the most tightly bound nucleus.



# Core Collapse

- Core can no longer support itself by nuclear burning.
- Core collapses, new energy source is gravitational contraction.
- High temperature and density lead to:
- Nuclear photodisintegration:
  - $\gamma + {}^{56}\text{Fe} \rightarrow 13({}^4\text{He}) + 4n$  (-124 MeV)
  - $\gamma + {}^4\text{He} \rightarrow 2p + 2n$  (-28.3 MeV)
- Neutronization:
  - $e^- + p \rightarrow n + \nu_e$
  - $e^- + {}^{56}\text{Fe} \rightarrow {}^{56}\text{Mn} + \nu_e$
  - $e^- + {}^{56}\text{Mn} \rightarrow {}^{56}\text{Cr} + \nu_e$

# Core Collapse

- Energy loss  $\sim (124 + 13 \times 28.3) \text{ MeV}/56 \sim 8.8 \text{ MeV/nucleon}$ .
- For Chandrasekhar mass of  $10^{56}$  nucleons, release  $10^{52}$  erg.
  - Compare with solar luminosity of  $4 \times 10^{33} \text{ erg/s}$ .
- Energy is released as neutrinos, leaves the star.
- Core collapses on close to free-fall time scale
  - Core density at end of Si burning  $\rho \sim 10^9 \text{ g/cm}^3$ .
  - Free-fall time scale  $(3\pi/32G\rho)^{1/2} \sim 0.1 \text{ s}$
- Most of nucleons (99.5%) are converted to neutrons.
  - Note free neutron lifetime  $\sim 886$  seconds.
- Tremendous release of energy creates a **supernova explosion**.

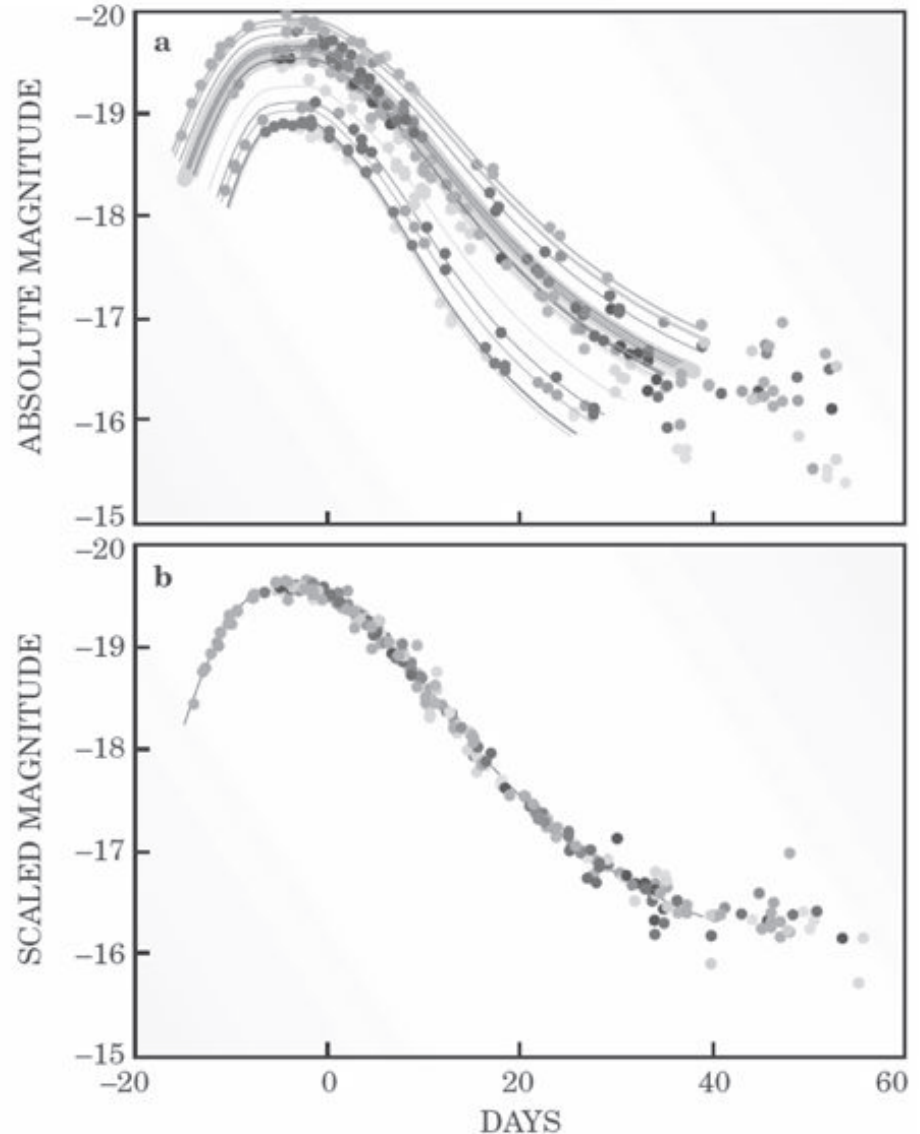
# Types of Supernovae

- Type I – no hydrogen absorption lines
  - Ia – no hydrogen lines, no helium lines, late in decay strongest lines are iron
  - Ib, Ic – helium lines, still no hydrogen
- Type II – hydrogen absorption lines
- Collapse of massive stars leads to type II and Ib, only difference is whether or not star sheds outer hydrogen layer before exploding
- Type Ia is thought to be white dwarf exceeding Chandrasekhar limit either via accretion or merger. Star is evaporated and no remnant remains.



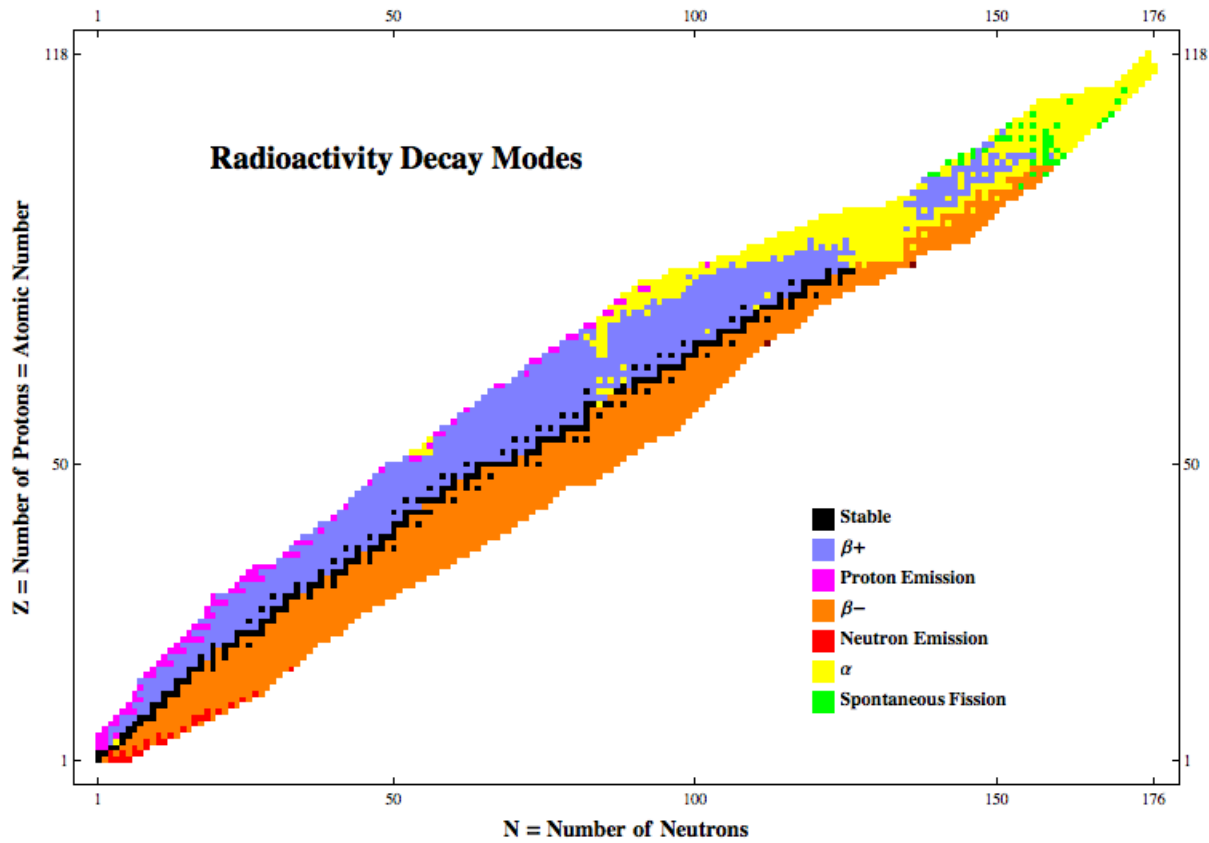
# Type Ia Supernovae

- Type Ia light curves reach the same maximum luminosity if scaled by width in time.
- Makes them useful for cosmology.



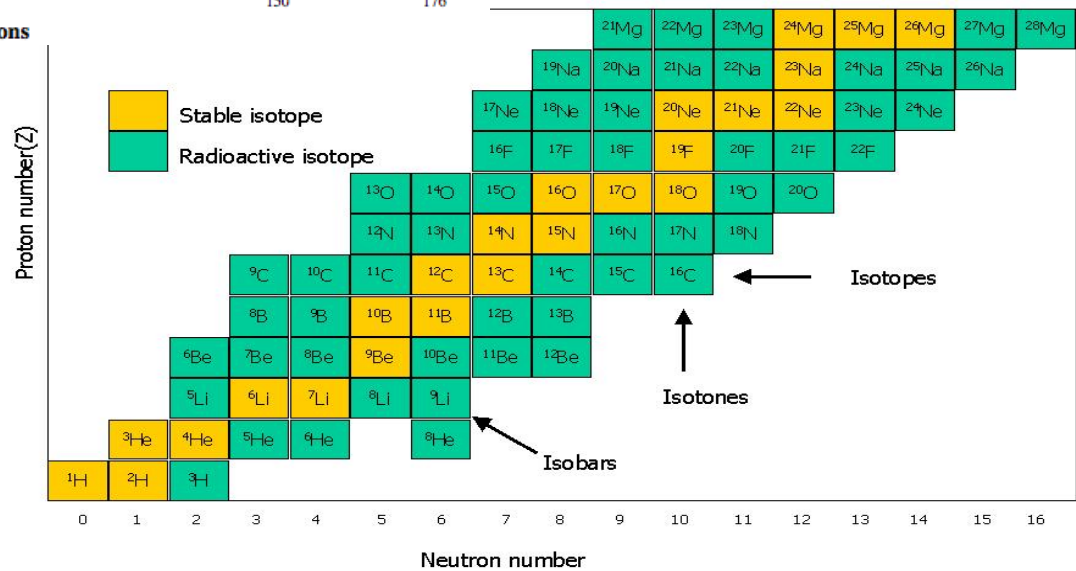
# Nucleosynthesis

- Fusion (slow) – Nuclear burning in stars that is energetically favorable. Produces the abundant elements up to Fe.
- S-process (slow) - Rate of neutron capture by nuclei is slower than beta decay rate. Produces stable isotopes by moving along the valley of stability. Occurs in massive stars, particular AGB stars.
- R-process (Rapid) – Rate of neutron capture fast compared to beta decay. Forms unstable neutron rich nuclei which decay to stable nuclei.

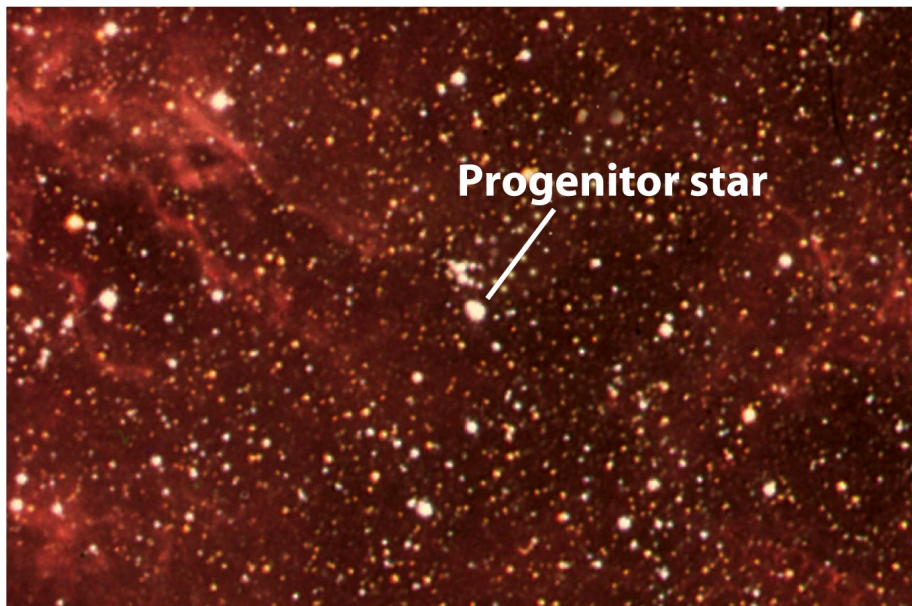


# Valley of Stability

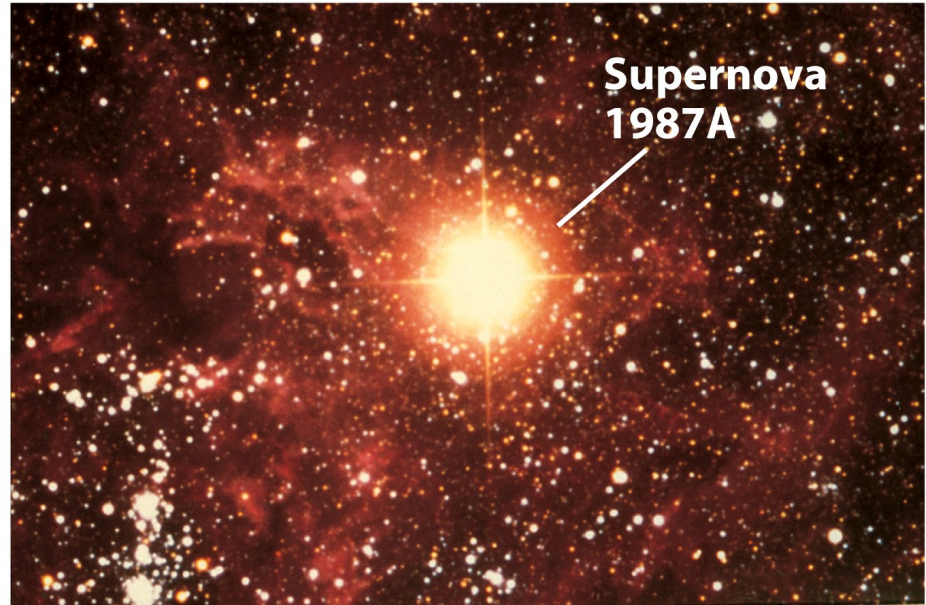
**N = Number of Neutrons**



In 1987 a nearby supernova gave us  
a close-up look at the death of a  
massive star

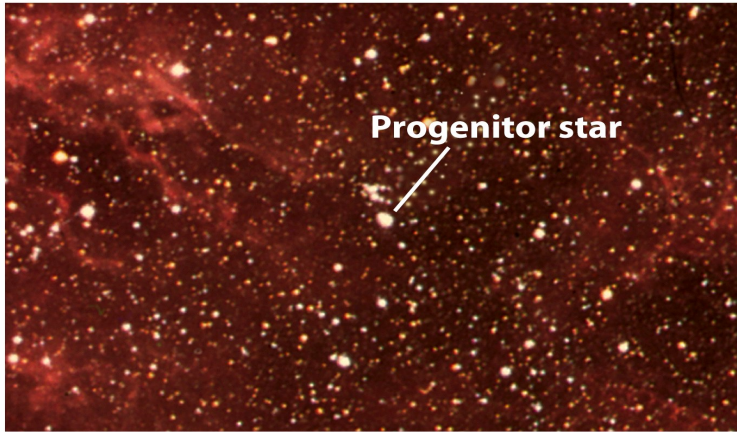


**Before the star exploded**

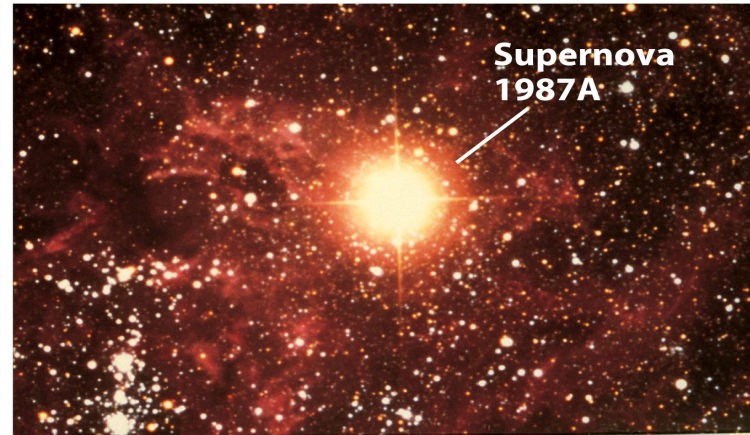


**After the star exploded**

# Neutrinos from SN1987A

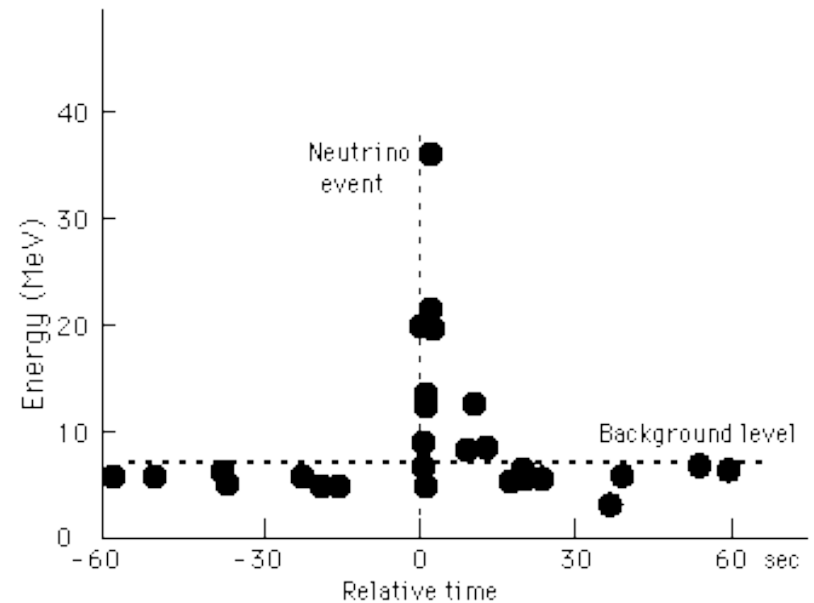


Before the star exploded



After the star exploded

- Supernova 1987A went off in the Magellanic Clouds.
- Neutrinos were detected on Earth.

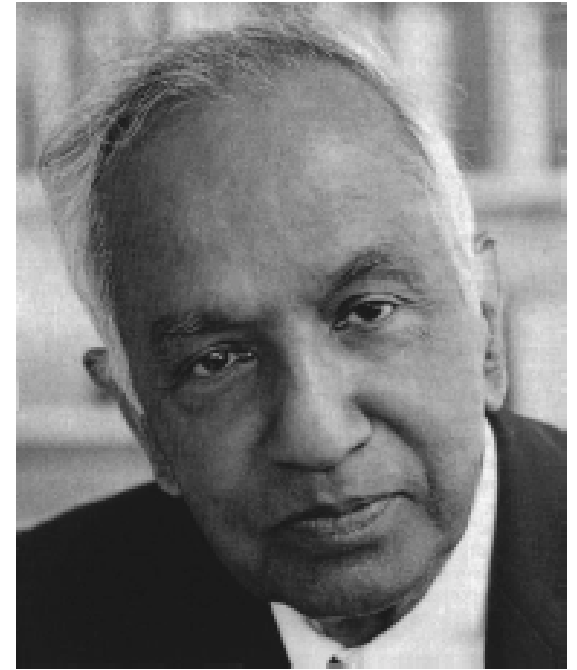


# Maximum white dwarf mass

Electron degeneracy cannot support a white dwarf heavier than 1.4 solar masses

This is the “Chandrasekhar limit”

Won Chandrasekhar the 1983 Nobel prize in Physics



# How Dense is Quantum?

- Need to use quantum statistics when average separation between particles is comparable to the de Broglie wavelength  $\lambda = h/p$ .
- Thermal energy  $E = p^2/2m = (3/2) kT$ , find  $\lambda = h(3mkT)^{-1/2}$ .
  - Note  $\lambda$  is larger for electrons, since  $m_e \ll m_p$ .
  - Quantum statistics will become important first for electrons.
  - Electrons will “become degenerate” first.
- Density  $\rho = m_p / (\text{volume per particle}) = m_p / (\lambda/2)^3 = 8m_p (3mkT/h^2)^{3/2}$ 
  - Note use  $m_p$  for density, but  $m_e$  for  $\lambda$ .
- As temperature decreases, gas becomes degenerate

# Relativistic degenerate gas

- As the density increases,  $\lambda$  becomes smaller, and  $p$  increases.
- Eventually electrons are relativistic  $E = pc = hc/\lambda^{-1/2}$ .
- Let's find the equation of state...
- Pressure  $P \sim$  energy density  $\sim E/\lambda^{-3} \sim hc/\lambda^{-4}$ .
- Still have mass density  $\rho \sim m_p/\lambda^3$ , therefore  $P \sim hc (\rho/m_p)^{-4/3}$
- Can work out the correct numerical coefficients using the Fermi-Dirac distribution (result is in the book, not the derivation).
- Find Chandrasekhar mass on board.

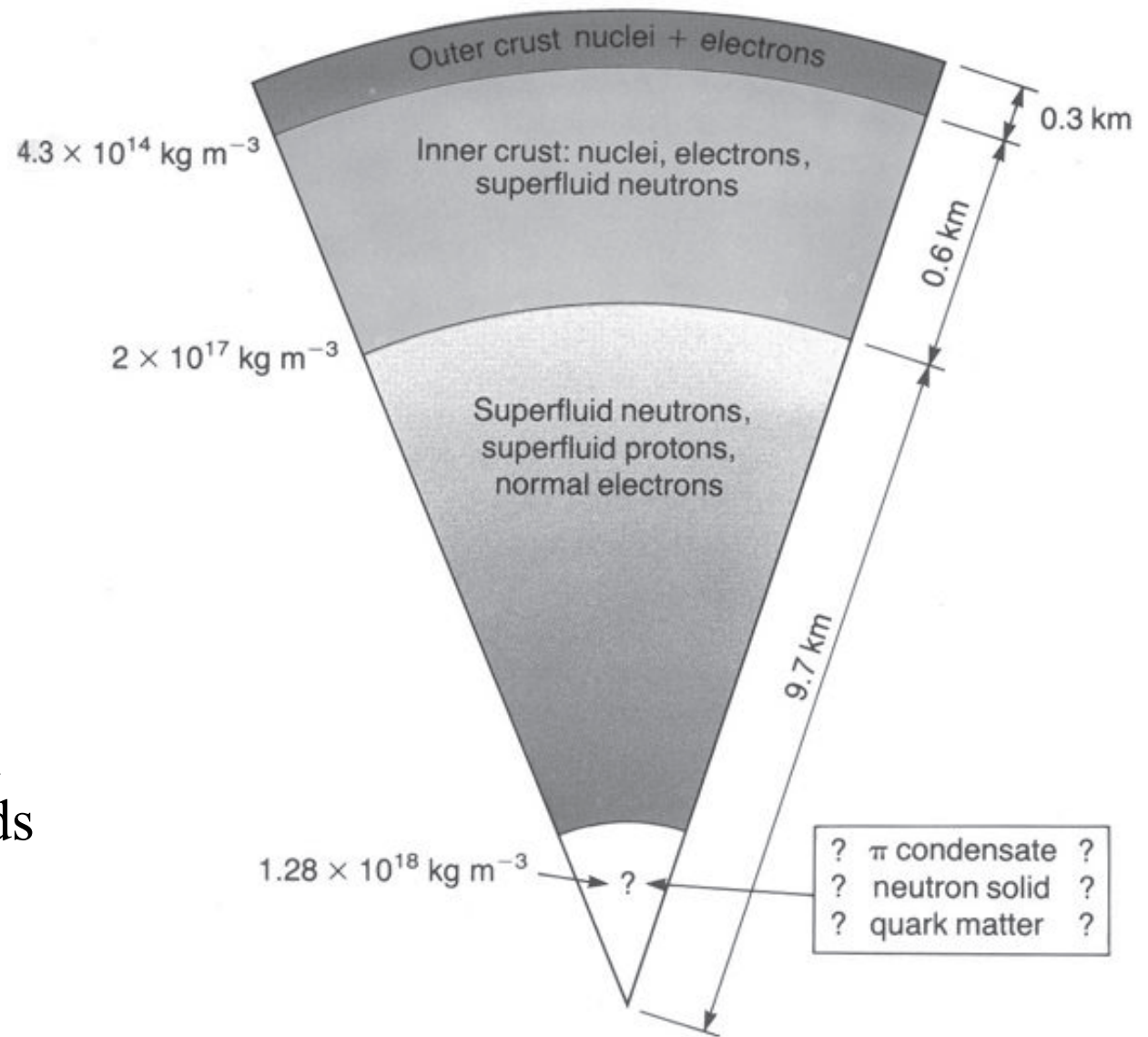


# Relativistic degenerate gas

- Chandrasekhar mass is 1.46 solar masses.
- Can do same calculation for neutron stars.
  - Need to use general relativity because object is so dense.
  - Need to extrapolate nuclear equation of state beyond what has been measured in the laboratory – makes estimates uncertain.
- Maximum neutron star mass is around 2-3 solar masses.

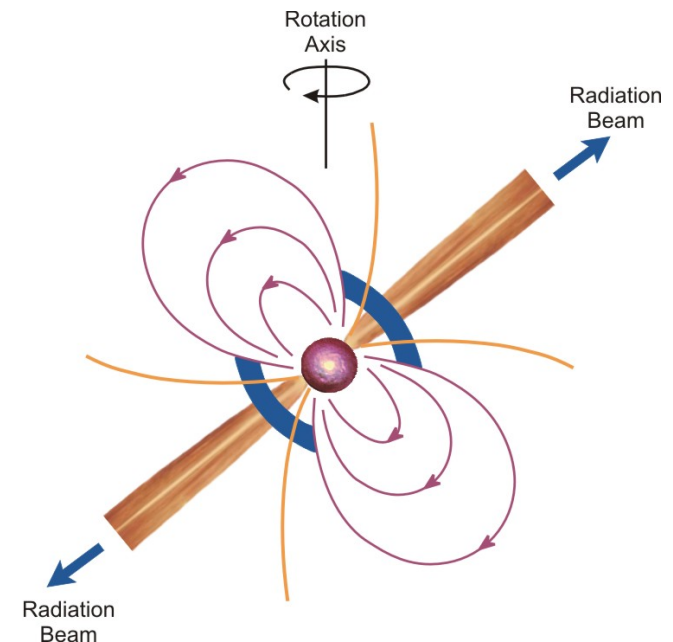
# Neutron Stars

- Radius  $\sim$  10-15 km
- Contains superfluids



# Little Green Men

- First pulsar discovered in 1967 with pulse period of 1.33 s. Found by Jocelyn Bell, her advisor, Anthony Hewish, won the Nobel Prize in Physics for the discovery in 1974.
- What can make extremely regular pulsations?
- Discoverers named object “LGM-1” for “Little Green Men”.
- Additional pulsars discovered within a year, periods 0.03 s to several seconds.



# Crab Pulsar



# Spinning Star?

For a rotating object to remain bound, the gravitational force at the surface must exceed the centripetal acceleration:

$$\frac{GMm}{r^2} > m\omega^2 r \Rightarrow \frac{GM}{r^3} > \frac{4\pi^2}{P^2} \Rightarrow \rho > \frac{3\pi}{P^2 G}$$

For the Crab pulsar,  $P = 33$  ms so the density must be greater than  $1.3 \times 10^{11}$  g cm<sup>-3</sup>.

This exceeds the maximum possible density for a white dwarf, requires a neutron star.

# Magnetic Field

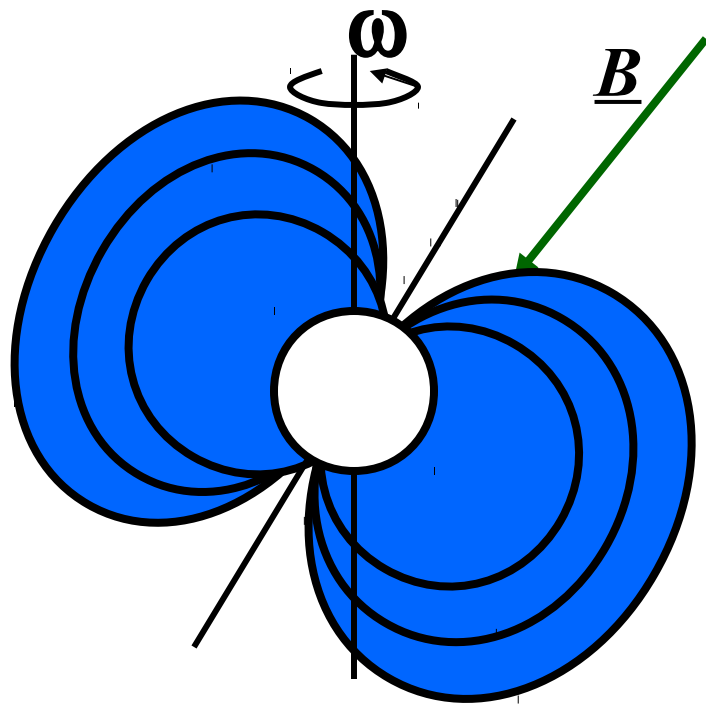
If a solar type star collapses to form a neutron star, while conserving magnetic flux, we would naively expect

$$R_{sun}^2 B_{sun} = R_{ns}^2 B_{ns} \Rightarrow \frac{B_{ns}}{B_{sun}} \approx 5 \times 10^9$$

For the sun,  $B \sim 100$  G, so the neutron star would have a field of magnitude  $\sim 10^{12}$  G.

# Magnetosphere

Neutron star rotating in vacuum:



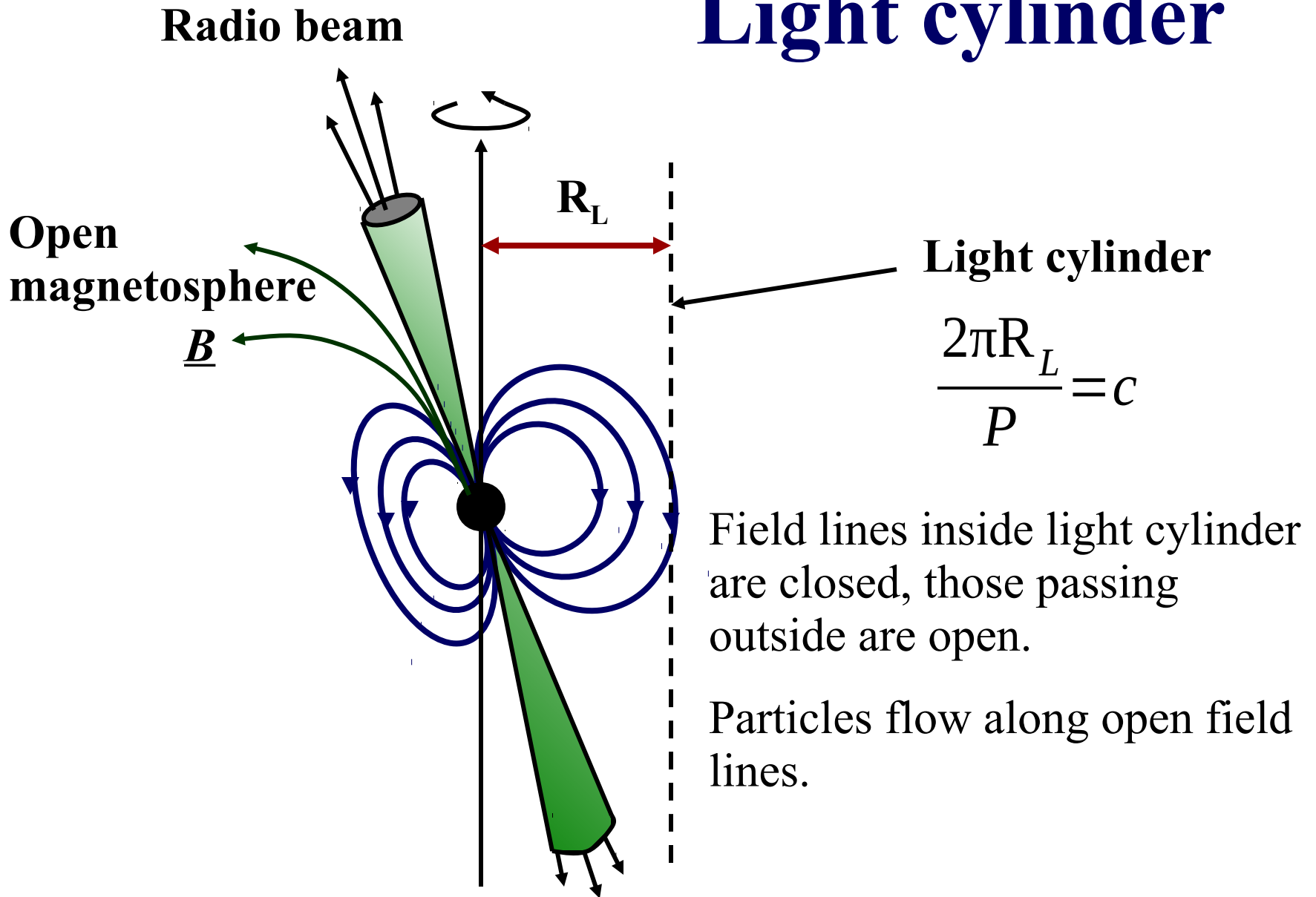
Electric field induced immediately outside NS surface.

$$E \simeq \frac{v}{c} B$$

The potential difference on the scale of the neutron star radius:

$$\Phi = ER \sim 10^{18} V$$

# Light cylinder





# Spin down of a pulsar

$$\text{Energy } E = \frac{1}{2} I (2\pi\nu)^2$$

$$\text{Power } P = -\frac{dE}{dt} = 4\pi^2 I\nu \frac{d\nu}{dt}$$

For Crab pulsar:  $\nu = 30/\text{s}$ ,  $M = 1.4$  solar masses,  $R = 12$  km,  
and  $d\nu / dt = -3.9 \times 10^{-10} \text{ s}^{-2}$ .

Therefore,  $P = 5 \times 10^{38} \text{ erg/s}$ .

Over a year, the spin rate changes by 0.04%.

# Dipole Radiation

Magnetic dipole,  $p_m$ , will radiate if rotated around misaligned axis.

$$\frac{dE}{dt} \propto \Omega^4 p_m^2$$

Equate this to loss of rotational energy:

$$\frac{d}{dt} \left( \frac{1}{2} I \Omega^2 \right) = I \Omega^4 \frac{d\Omega}{dt} \propto \Omega^4 \rightarrow \frac{d\Omega}{dt} \propto \Omega^3$$

# Braking Index

In general can write  $\frac{d\Omega}{dt} \propto \Omega^n$

where  $n$  is the ‘braking index’.

Can estimate  $n$  from spin period  $P$  and its time derivatives:

$$n = 2 - \frac{P \ddot{P}}{\dot{P}^2}$$

“Characteristic” age of the pulsar is then  $\tau = \frac{P}{(n-1)\dot{P}}$

Measurements of  $n$  range 1.8 to 3.0.

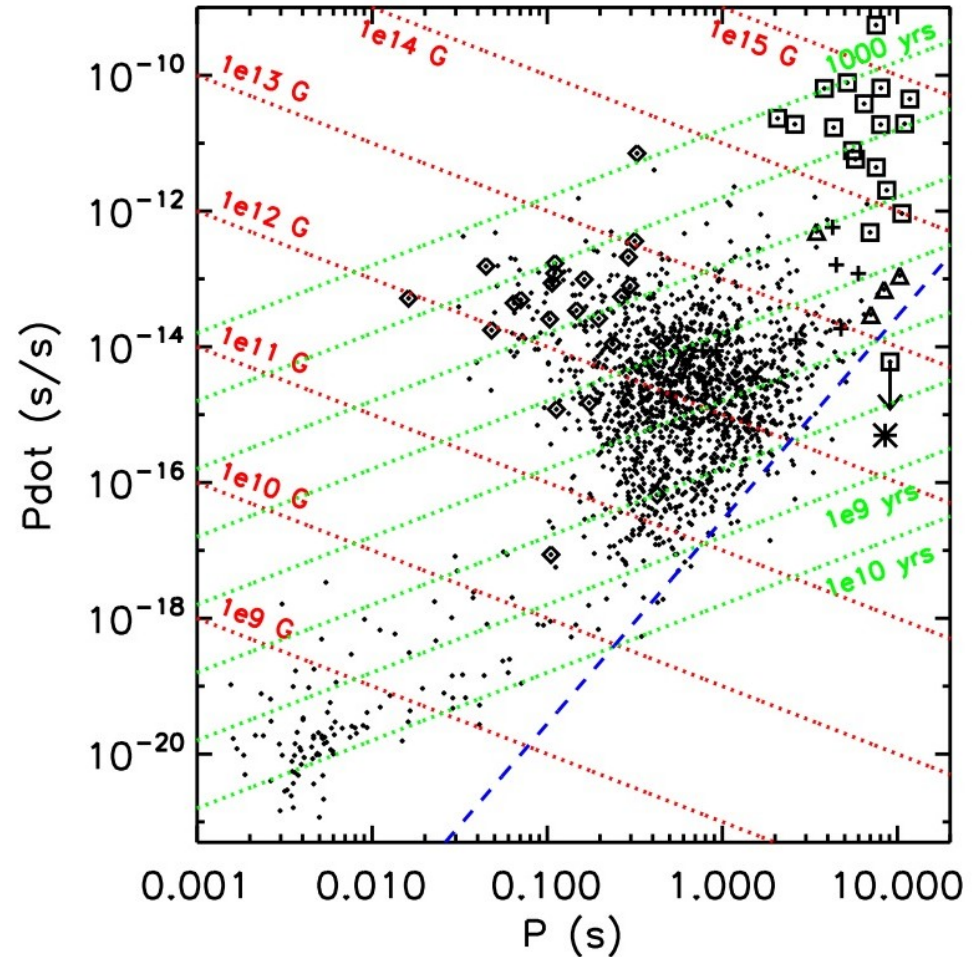
# P P-dot Diagram

If neutron stars have the same radius, then  $B \sim p_m$ ,

Using the dipole radiation formula, we can find

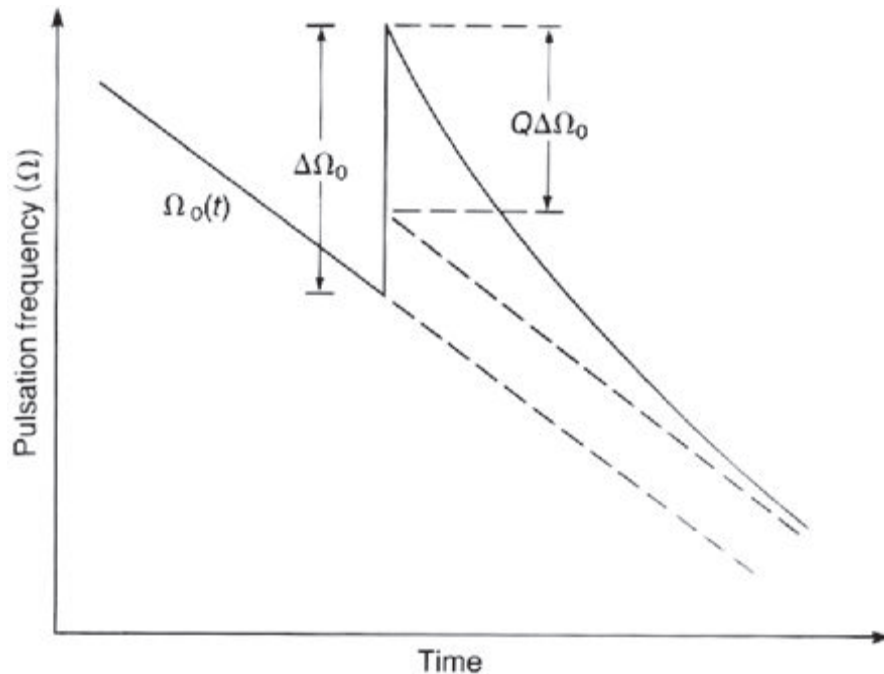
$$B_s \propto (P \dot{P})^{1/2}$$

$$\tau = \frac{P}{2\dot{P}}$$



# Pulsar Glitches

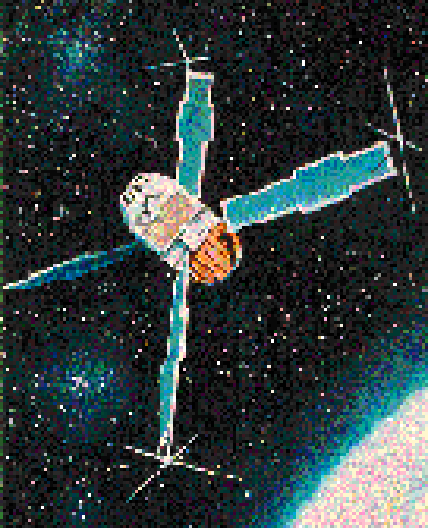
Short timescales - pulsar slow-down rate is **remarkably uniform**  
Longer timescales - irregularities apparent, in particular, 'glitches'



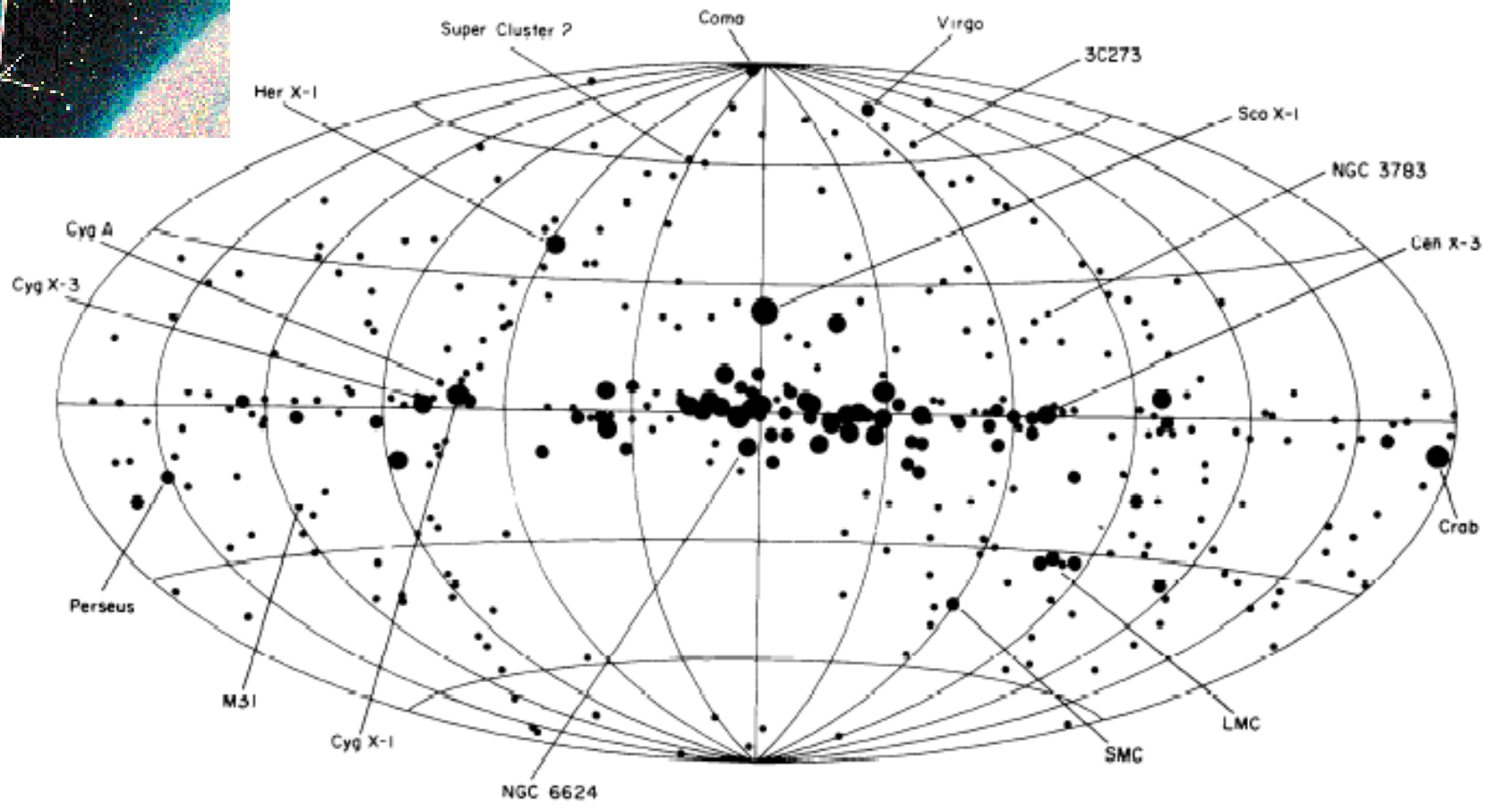
$$\frac{\Delta P}{P} \sim 10^{-10}$$

for Crab pulsar

Due to stresses and fractures in the crust?

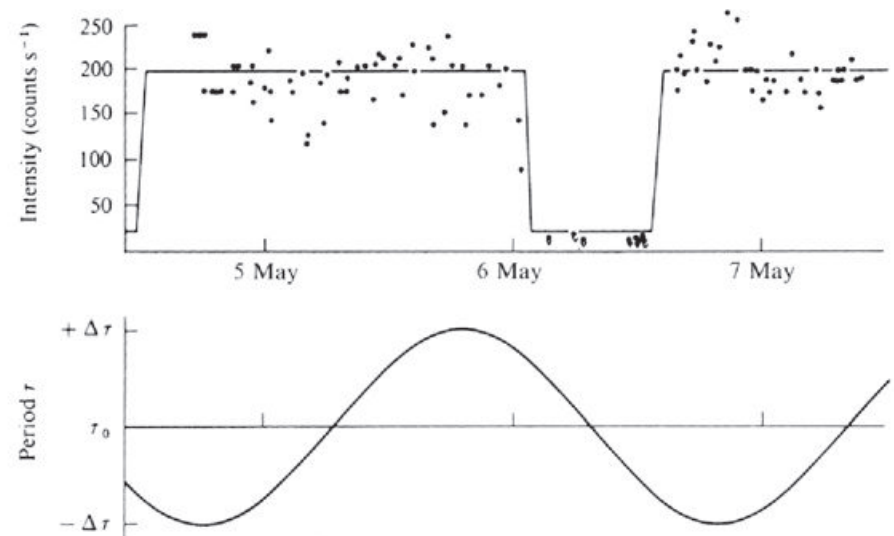
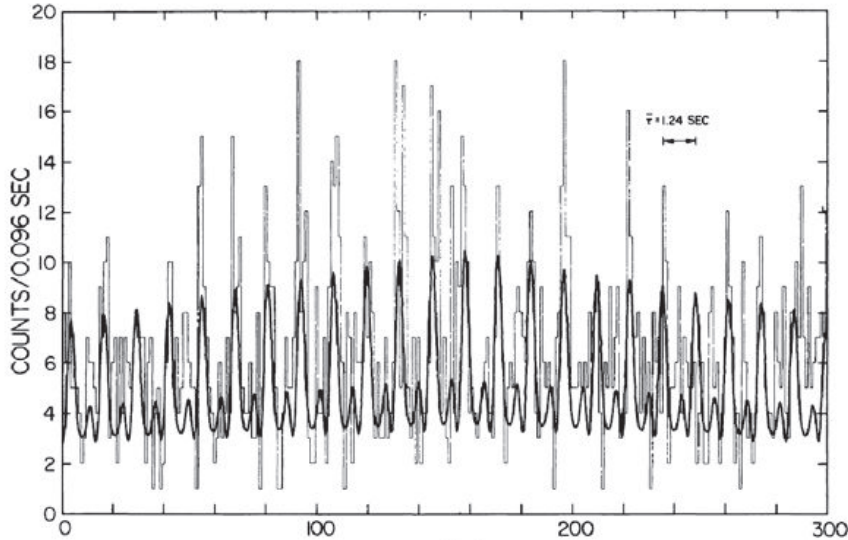


# Uhuru (1970-1973)



**339 sources**

# X-Ray Pulsar Cen X-3



- Pulses occur at intervals of 4.84 seconds and are modulated at orbital period of 2.09 days.
- There is an eclipse when companion star blocks X-ray source.

# Mass Function

- If one can detect spectral lines from the companion, then one can measure the binary orbit and constrain the mass of the neutron star.
- $K$  = peak velocity of companion,  $i$  = orbit inclination
- ‘Mass function’ is

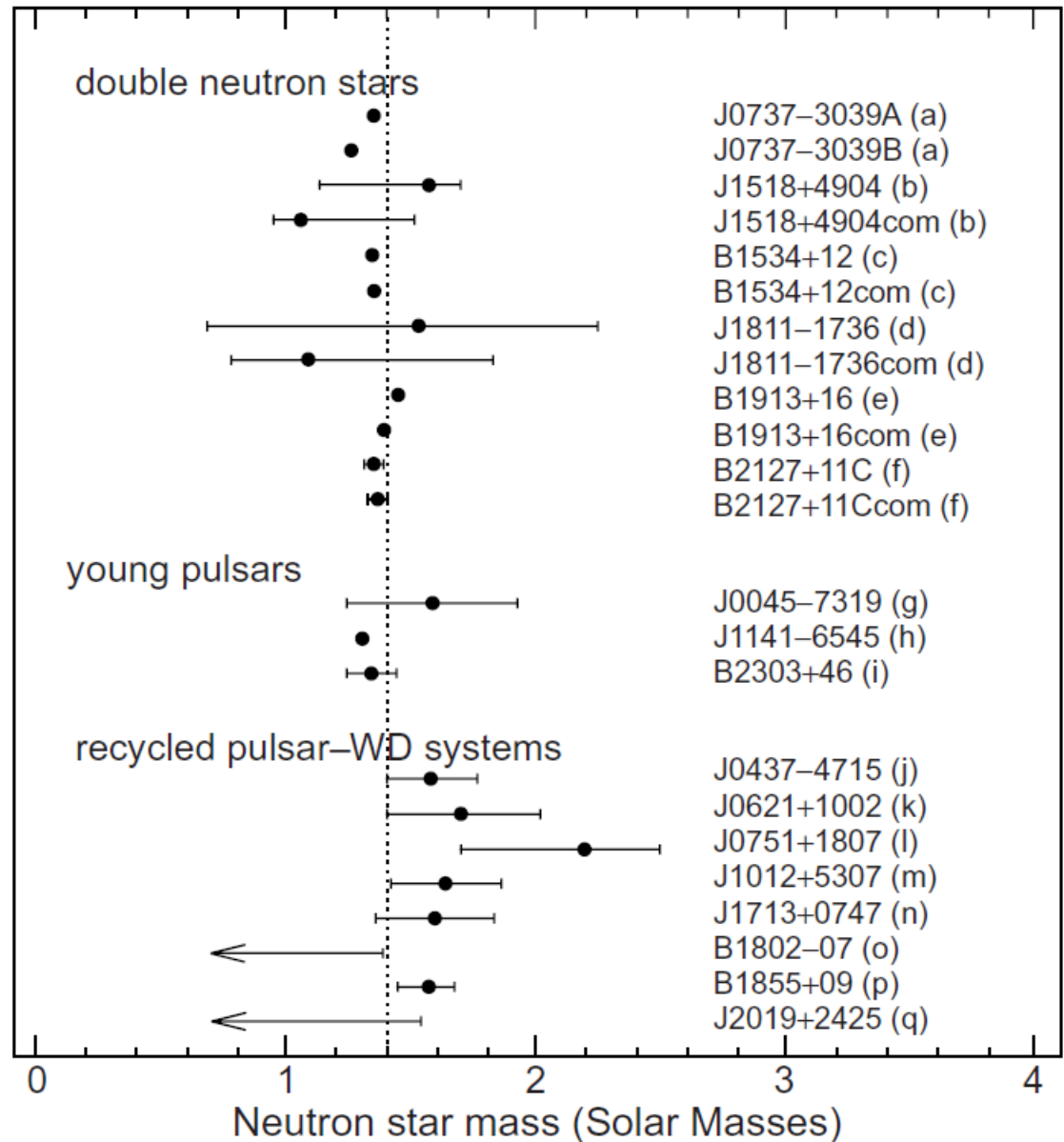
$$f = \frac{M_X^3 \sin^3 i}{(M_0 + M_X)^2} = \frac{P_{orb} K^3}{2\pi G} < M_X$$

- Note that book has mass function measured using velocity shifts of X-ray source, which gives constraint on mass of companion star.

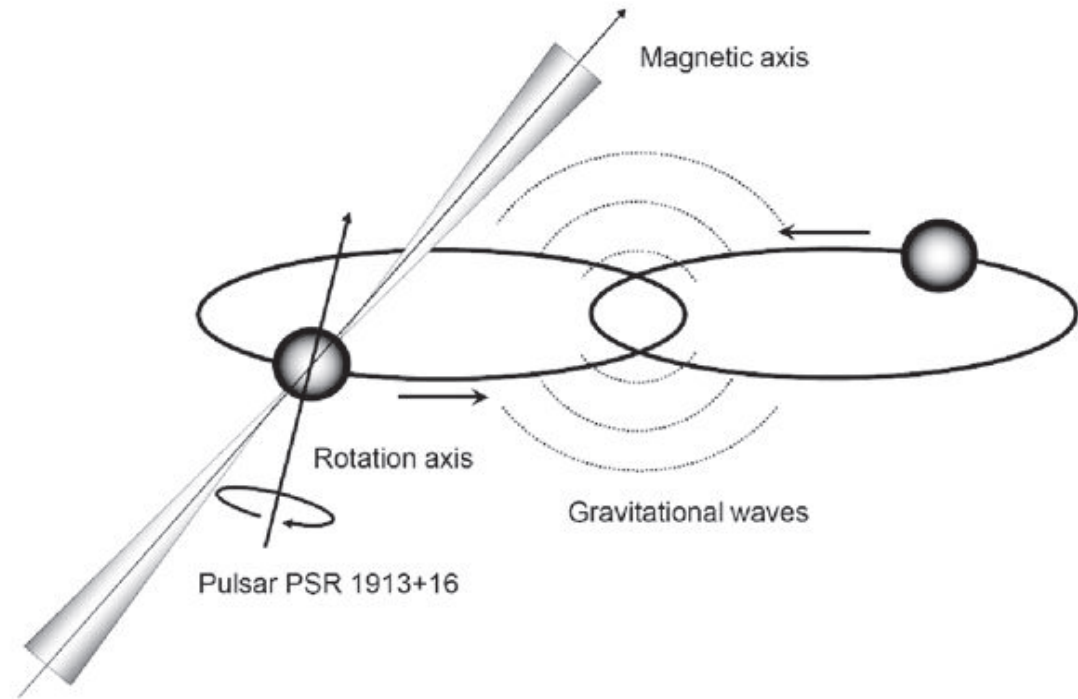


# Binary Pulsars

- Systems with two neutron stars permit accurate mass measurements and tests of general relativity.
- First such binary was found by Hulse and Taylor in 1975; they won the Nobel prize. Double pulsar system found in 2006.



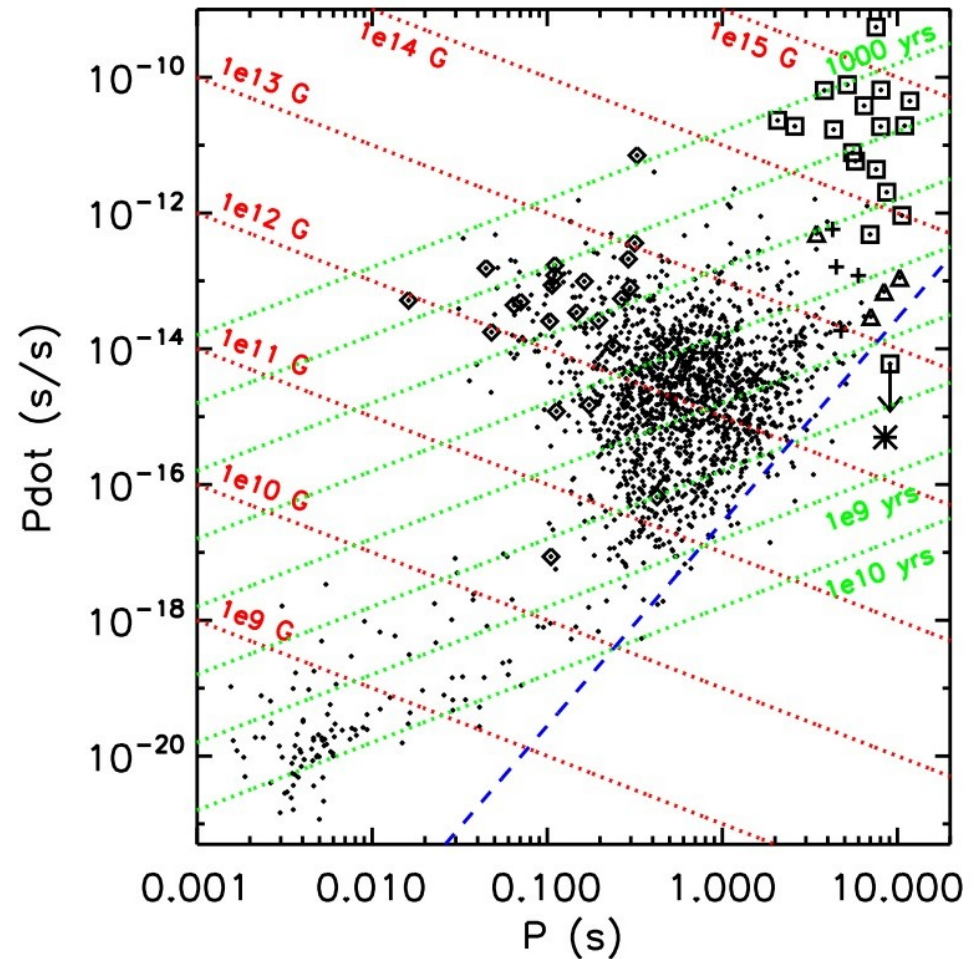
# Binary Pulsars



- Hulse and Taylor found that the binary orbit was gradually shrinking in a manner consistent with the binary losing energy to gravitational waves.
- Gravitational waves finally detected last year, but in a BH-BH binary.

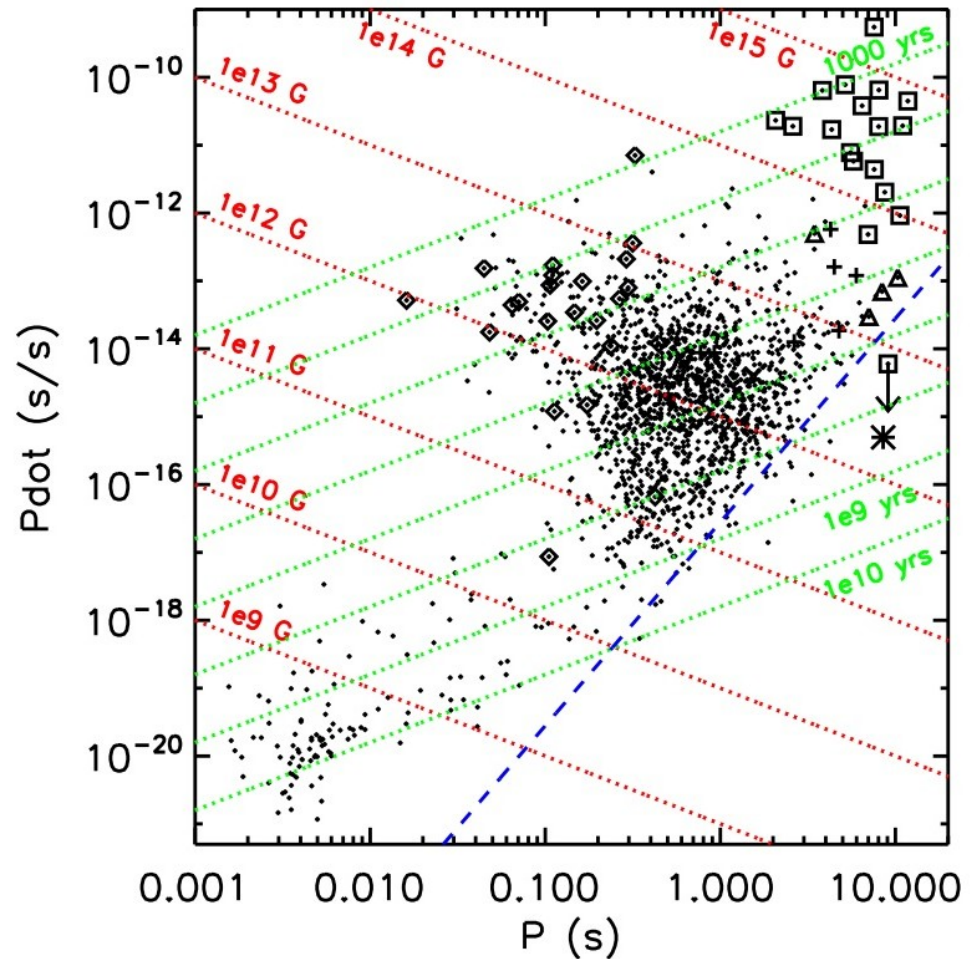
# Millisecond Pulsars

- There are pulsars with very short periods, milliseconds.
- How are they produced?
- Faster than the youngest pulsars. Also, have long characteristic ages and low magnetic fields. Often found in globular clusters.
- Millisecond pulsars are thought to be spun-up via accretion from a companion star.

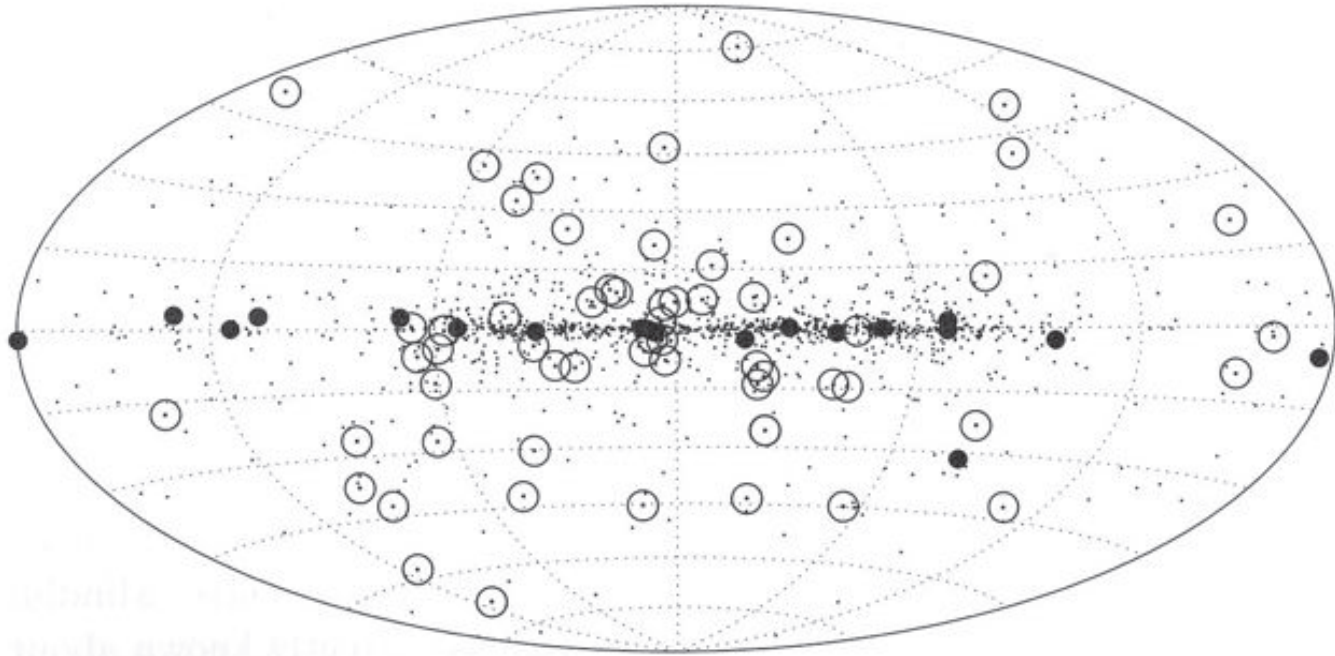


# Magnetars

- There are pulsars with very high magnetic fields (squares) and few second periods.
- They also produce giant outbursts as ‘soft gamma repeaters’ and ‘anomalous X-ray pulsars’.
- Spinning down too slowly to power outbursts via spin.
- Outburst thought to be due to release of magnetic energy via starquakes.



# Galactic Distribution of Pulsars



- Large dots are pulsars with SNR.
  - Why all in plane?
- Small dots are other regular pulsars.
  - Broader distribution due to velocities up to 1600 km/s.
- Dots in circles are millisecond pulsars.
  - Why wide distribution?