

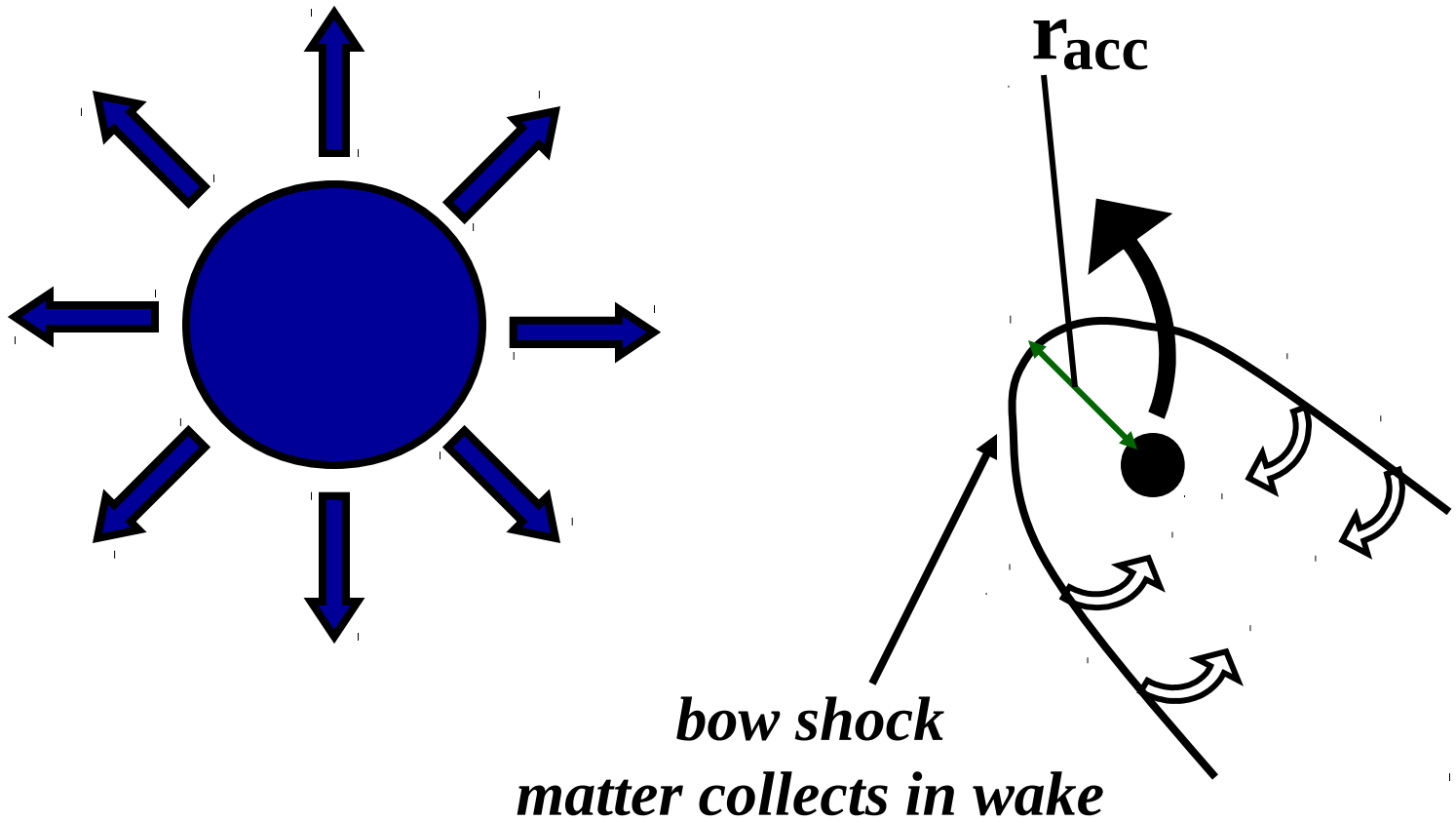
# Accretion in Binaries II

- Classes of X-ray binaries
  - Low-mass (BH and NS)
  - High-mass (BH and NS)
  - X-ray pulsars (NS)
  - Be/X-ray binaries (NS)
- Distinguishing NS versus BH binaries

# X-Ray Binaries

- Low mass: companion star mass less than one solar mass
- High mass: companion star mass greater than one solar mass.

# Wind Fed Accretion



# Wind Fed Accretion

Matter in wind will accrete if its speed is less than the escape speed from the compact object at the radius of closest approach

$V_w$  = wind velocity

$V_x$  = velocity of compact object

$R_c$  = capture radius

$$R_c = \frac{2GM_x}{V_x^2 + V_w^2}$$

# Luminosity

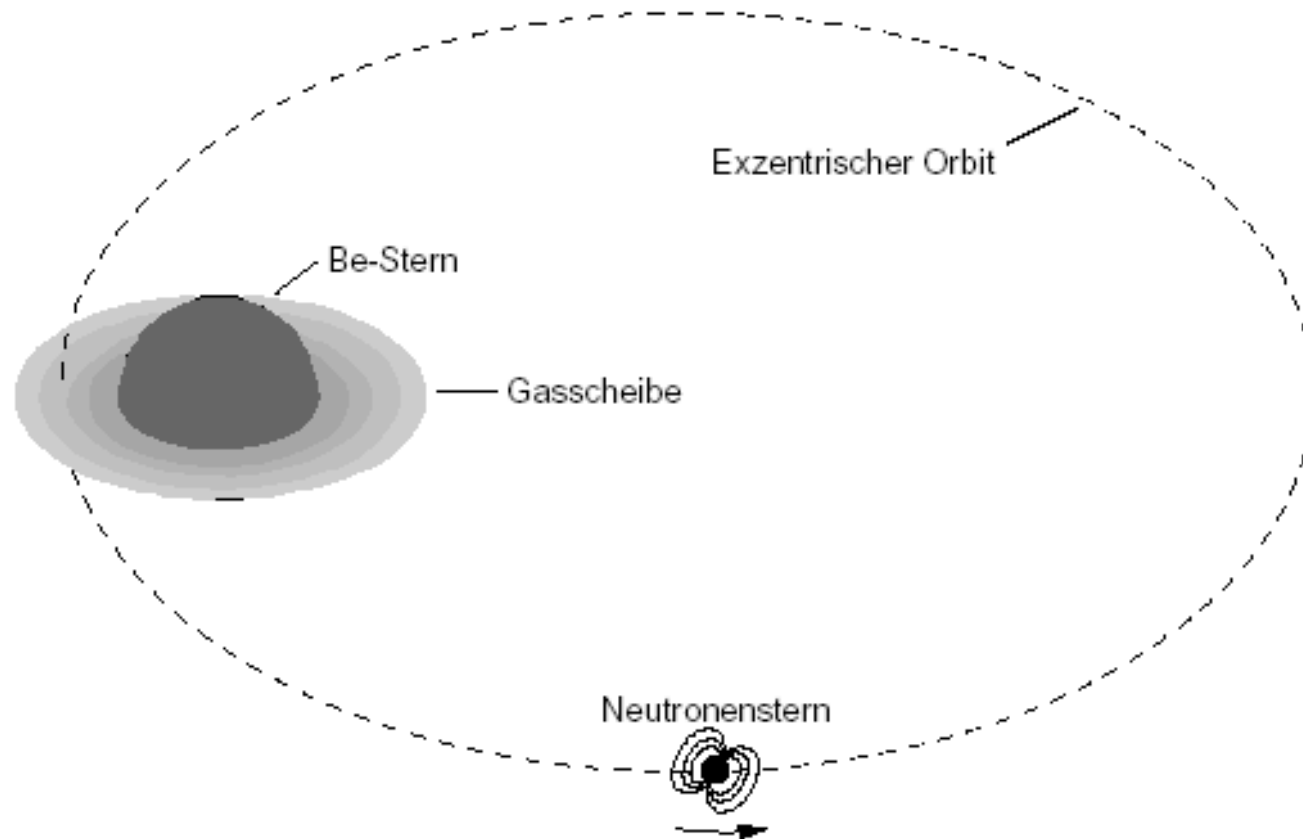
$$\dot{m} = \frac{\dot{m}_w}{4} \left( \frac{r_{acc}}{D_{orb}} \right)^2 \Rightarrow$$

$$L = \xi \dot{m} c^2 \approx \frac{\xi \dot{m}_w c^2}{4} \left( \frac{2GM}{D_{orb}} \right)^2 v_w^{-4}, \text{ if } v_w \gg v_X$$

If a neutron star binary has the following parameters,  $D_{orb} = 10^{12}$  cm,  $v_w = 1000$  km/s,

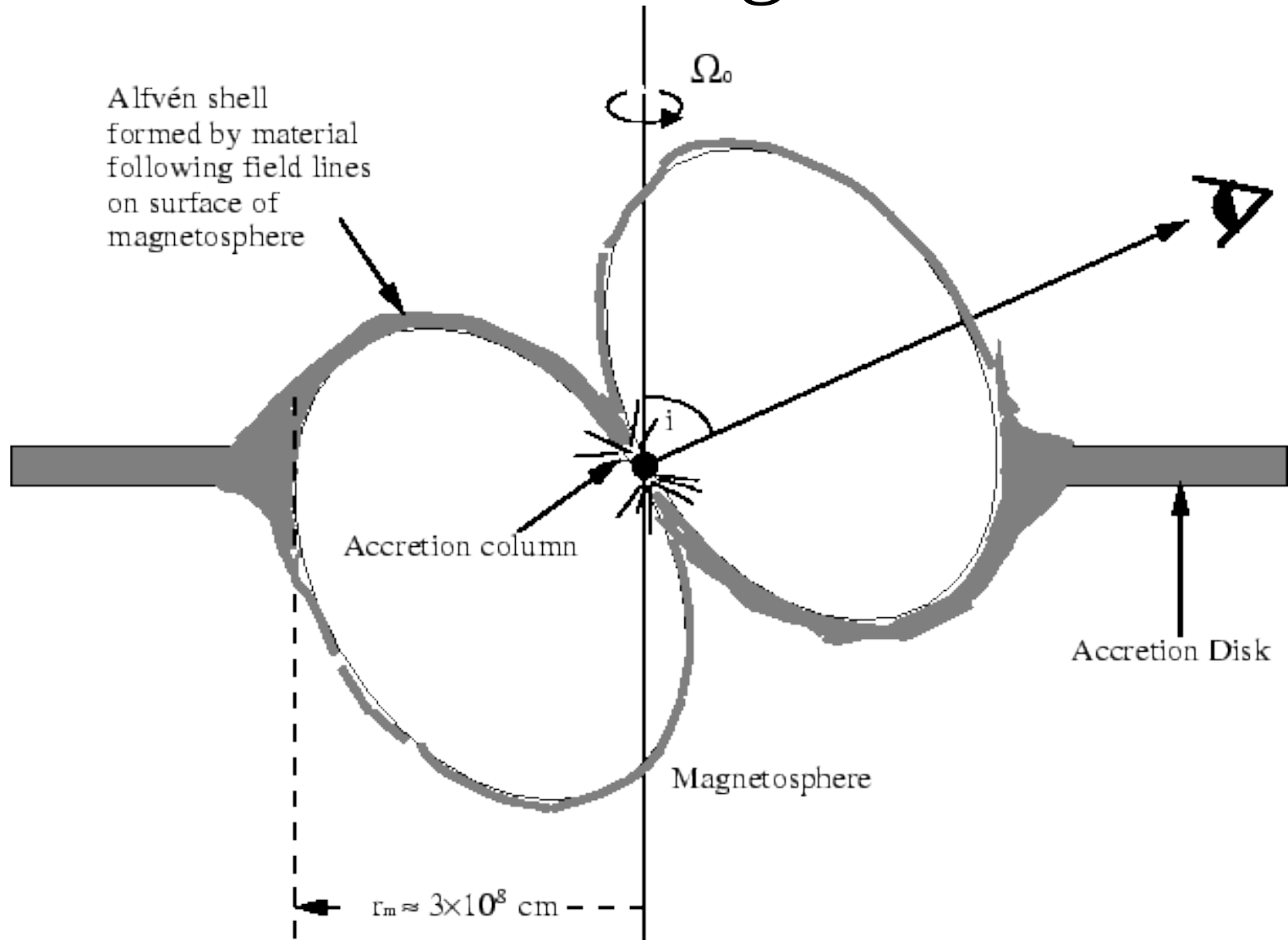
$$L \approx \left[ \frac{1}{4} \left( \frac{R_s}{D_{orb}} \right)^2 \left( \frac{c}{v_w} \right)^4 \right] \xi \dot{m}_w c^2 = 4 \times 10^{-4} (\xi \dot{m}_w c^2)$$

# Accreting Pulsars



High mass X-ray binary

# Accretion in magnetic field



# Magnetosphere boundary

Magnetospheric boundary,  $r_M$ , is where the magnetic pressure balances the ram pressure of accreted matter

$$p_{mag} = \frac{B^2}{8\pi} = p_{ram} = \rho v^2$$

Assuming spherical accretion and a dipole field with the dipole moment  $\mu = B_S R_S^3$ , where  $R_S$  and  $B_S$  are the radius of the star and the field strength at the surface.

$$r_M = 5.1 \times 10^8 \text{ cm} \left( \frac{\dot{m}}{10^{16} \text{ g/s}} \right)^{-2/7} \left( \frac{M}{M_\odot} \right)^{-1/7} \left( \frac{\mu}{10^{30} \text{ G} \cdot \text{cm}^3} \right)^{4/7}$$



# Accretion Torque

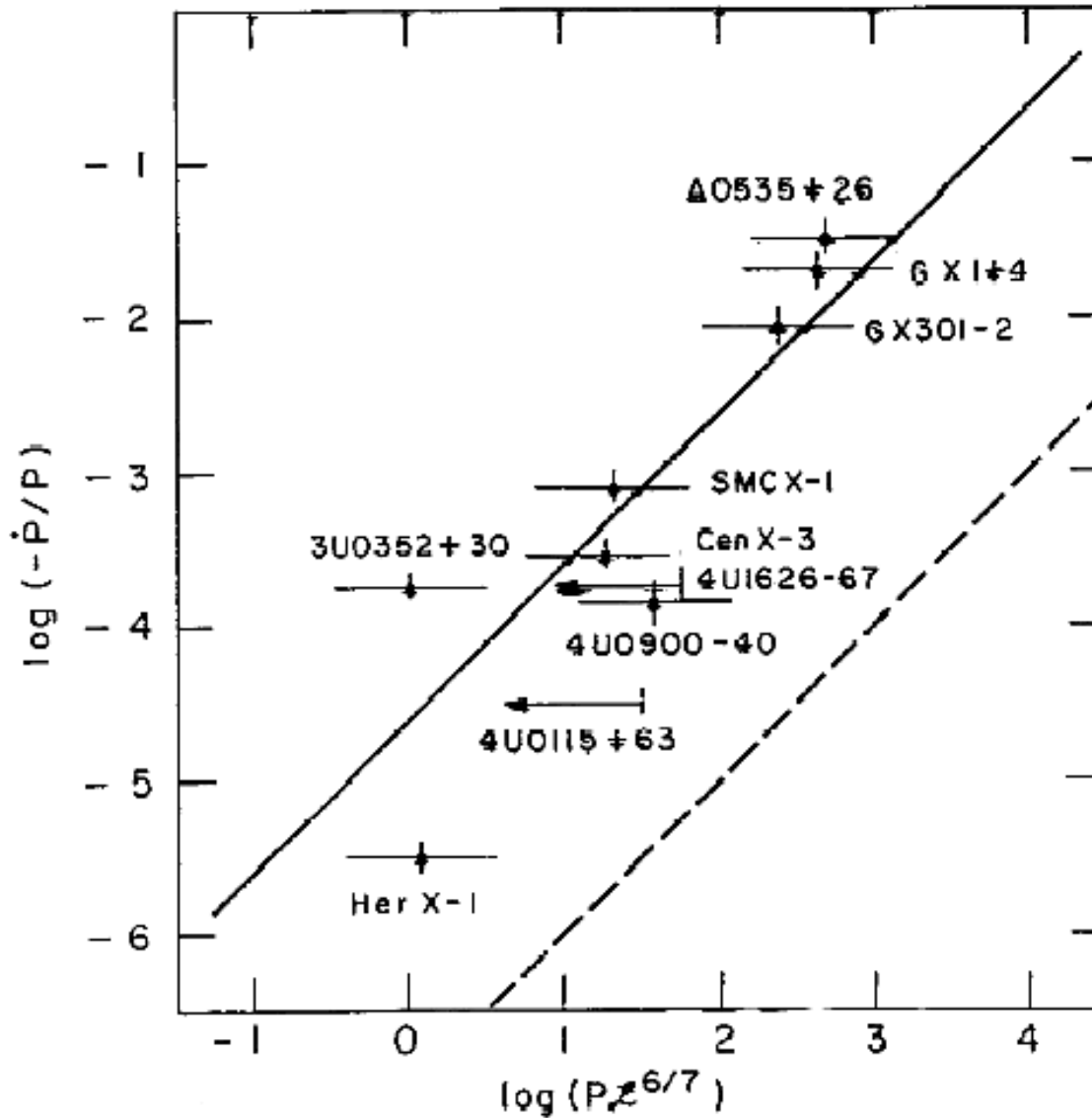
Angular momentum transport =  $\dot{m} v_{\phi} r$

Torque on star =  $I \dot{\omega} = \dot{m} v_{\phi} r_m$

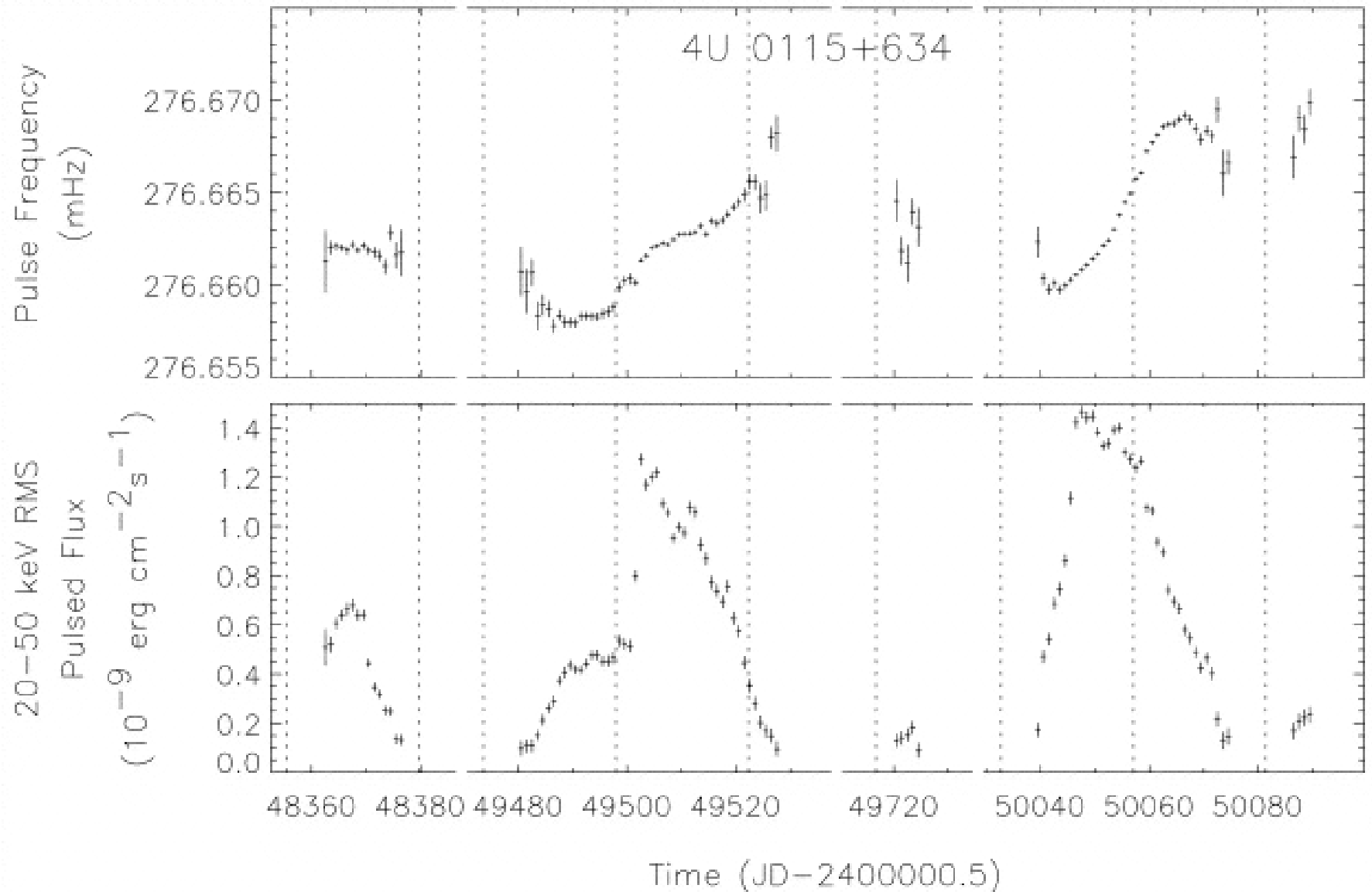
Note  $L = \frac{GM \dot{m}}{r_M}$  and  $\dot{\omega} = \frac{-2\pi}{P^2} \dot{P}$

Find  $\frac{\dot{P}}{P} \propto P L^{6/7}$

# Accretion Torque



# Pulse Period Variations



What happens if the spin rate of the pulsar is faster than Keplerian rotation rate at the magnetospheric boundary?

Corotation radius lies outside magnetospheric boundary

# Equilibrium Period

The accreted matter ceases to transfer angular momentum to the neutron star when

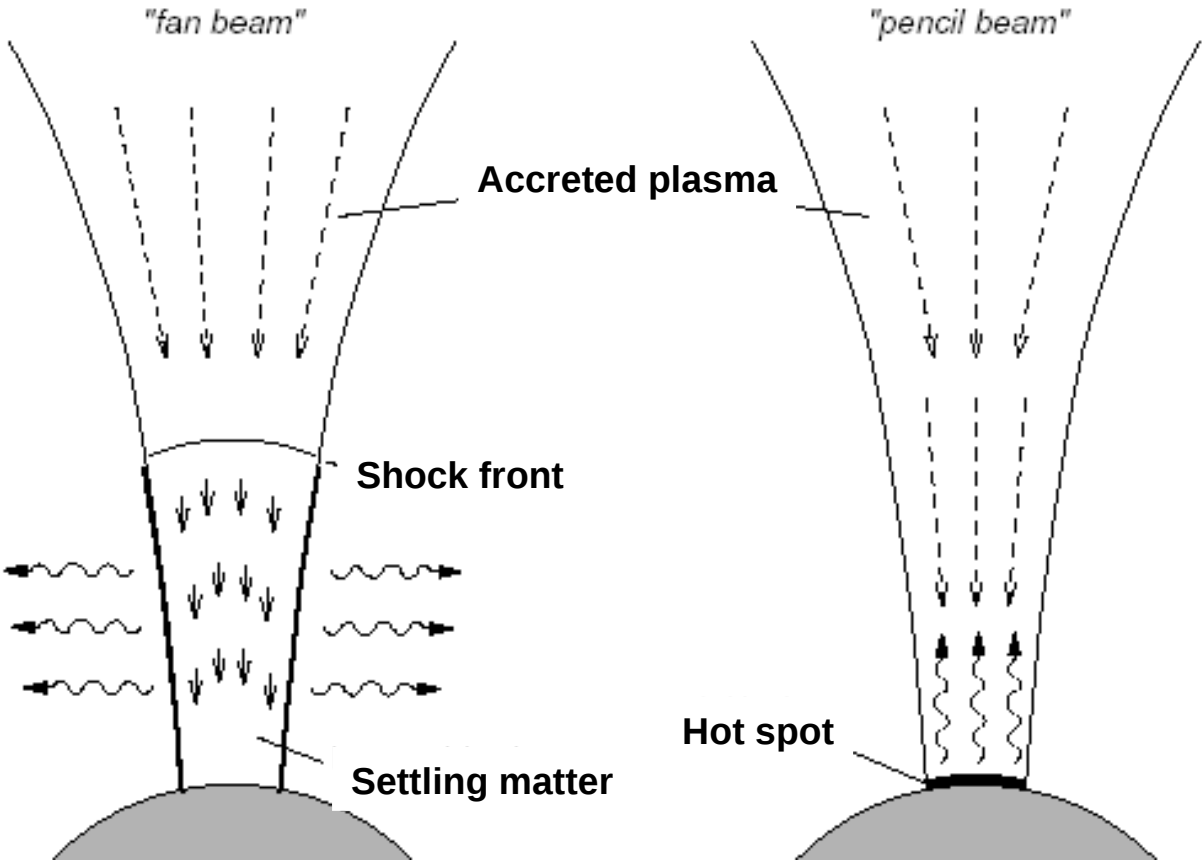
$$\omega = \Omega_K(r_M)$$

$$P_{eq} = 2\pi \left( \frac{r_M^3}{GM} \right)^{1/2} = \frac{2\pi}{(GM)^{1/2}} \left( \frac{R^{12}}{8\pi GM \dot{m}} \right)^{3/14} B_S^{6/7}$$

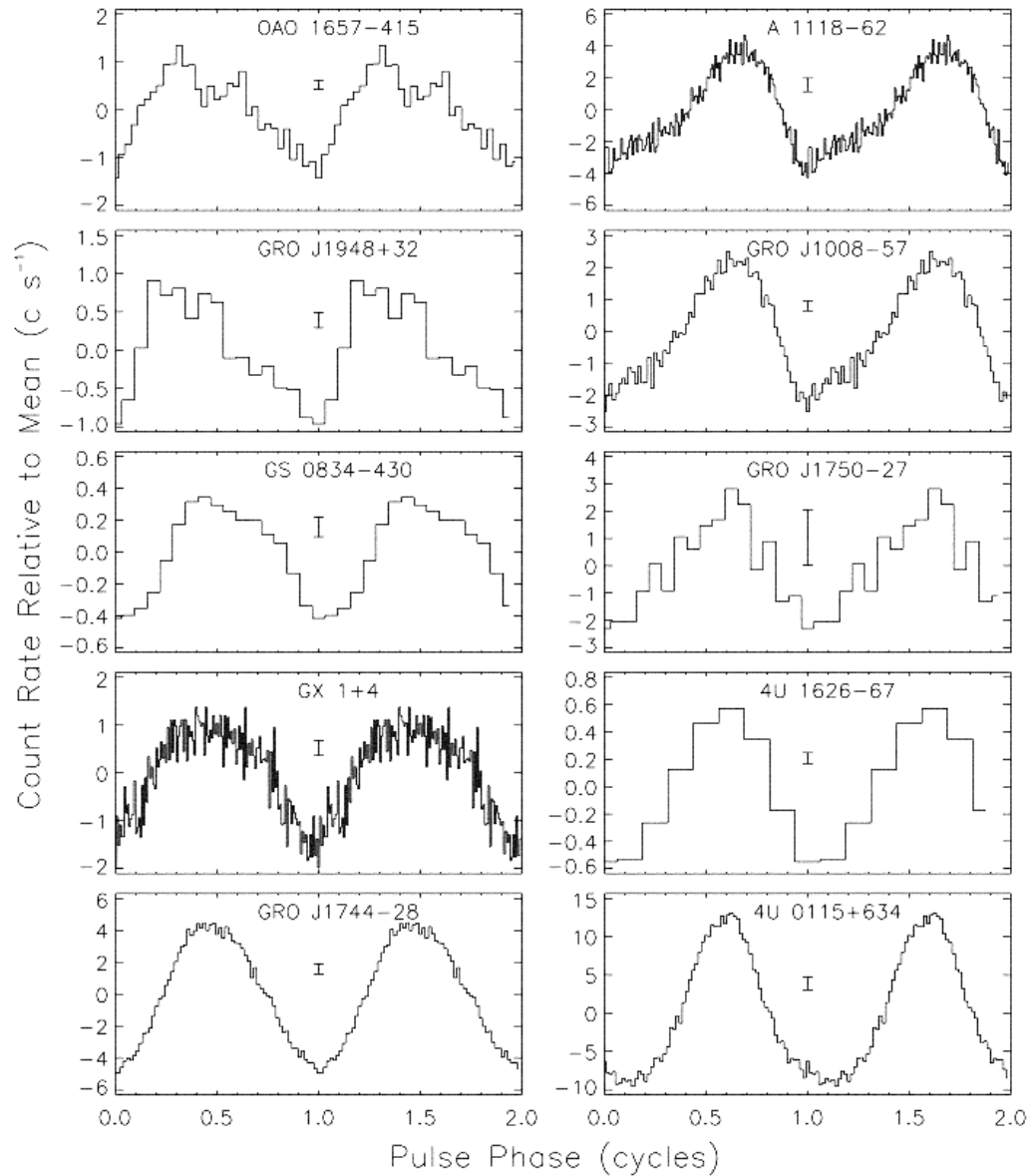
For a maximum (Eddington) luminosity and typical neutron star mass and radius, we have

$$\frac{P_{\min}}{0.002s} \approx \left( \frac{B}{10^9 G} \right)^{6/7}$$

# Emission Geometry



# Pulse Profiles



# Landau Levels

Quantization of energy due to magnetic field:

$$E_n = m_e c^2 \sqrt{1 + \left(\frac{p_{\parallel}}{m_e c}\right)^2 + 2n \frac{B}{B_{\text{crit}}}} \quad \text{Landau levels}$$

where  $p_{\parallel}$  is the momentum of the electron parallel to the field,  $n$  is the quantum number, and  $B_{\text{crit}}$  is the critical field,

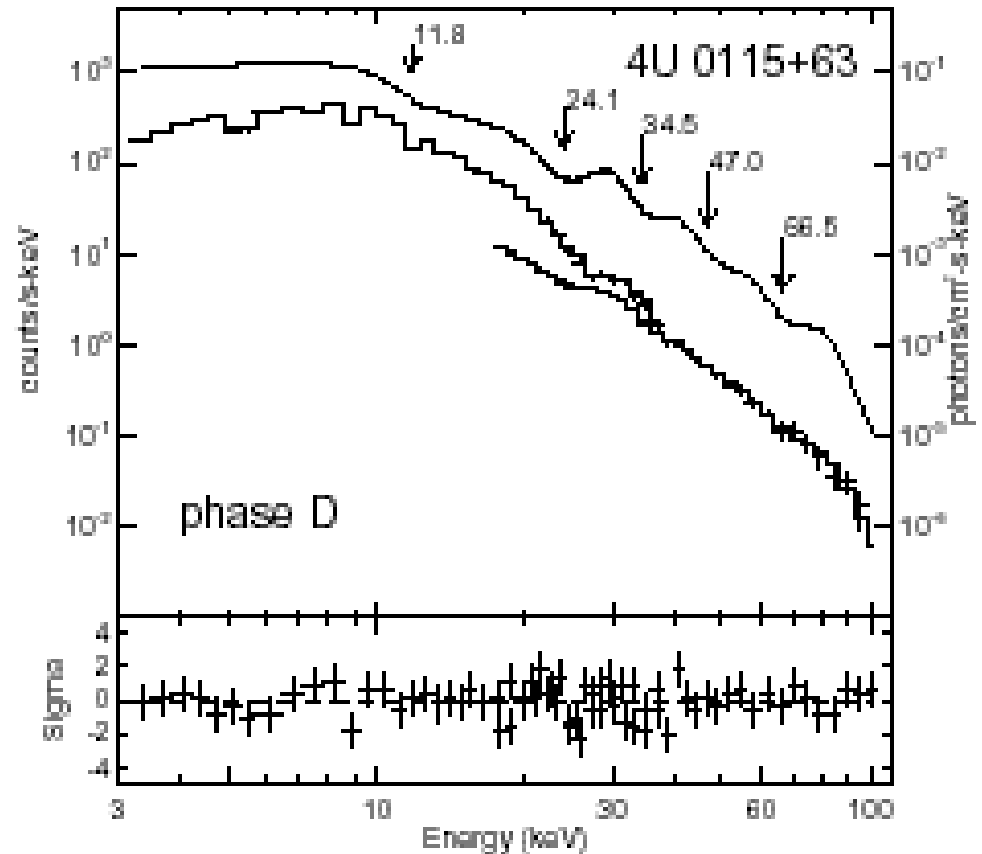
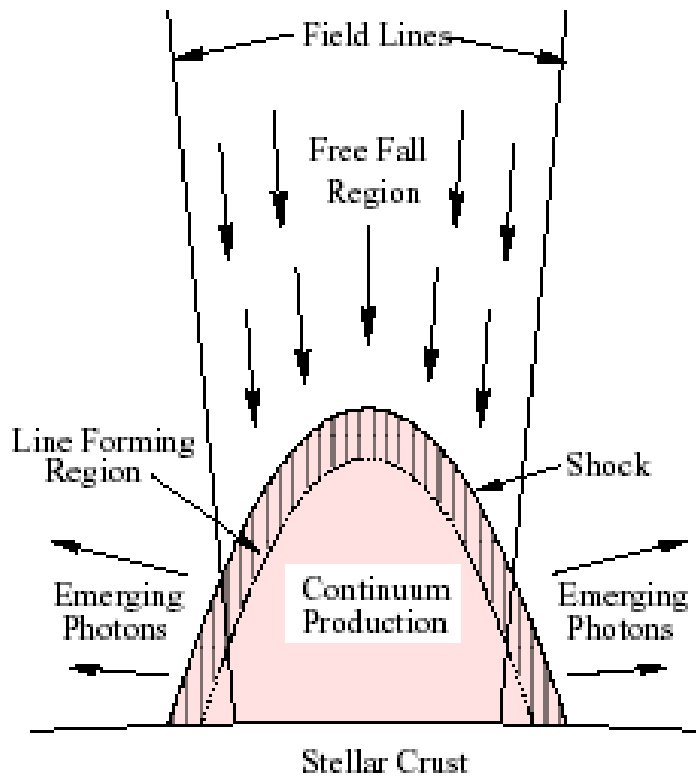
$$B_{\text{crit}} = \frac{m_e^2 c^3}{e \hbar} \approx 4.4 \times 10^{13} \text{ G}$$

For  $B \ll B_{\text{crit}}$ , the spacing between adjacent levels is

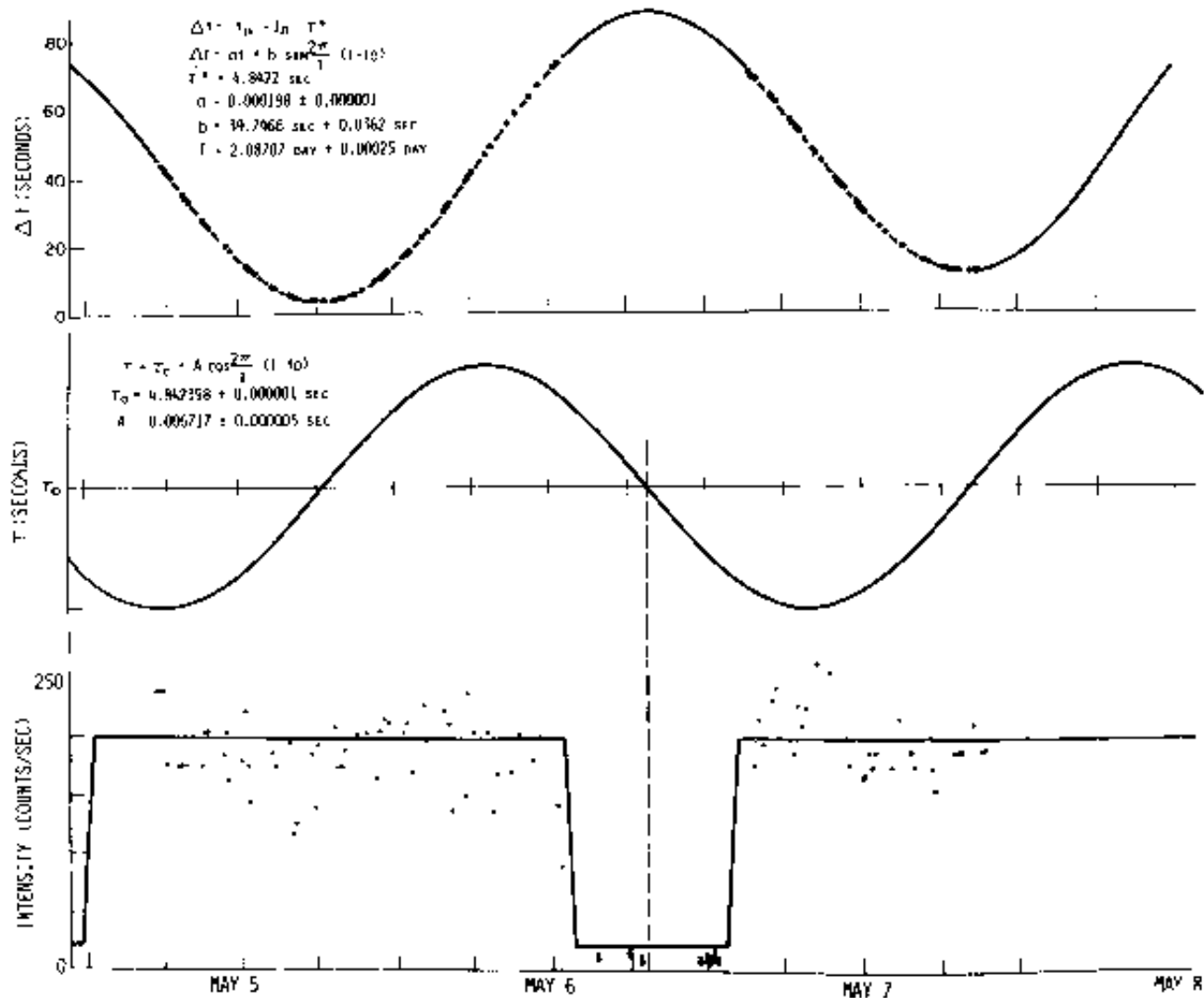
$$E_{\text{cyc}} = E_{n+1} - E_n = \frac{\hbar e}{m_e c} B = 11.6 \text{ keV} \left( \frac{B}{10^{12} \text{ G}} \right)$$



# Cyclotron Lines



# X-Ray Pulsar Cen X-3

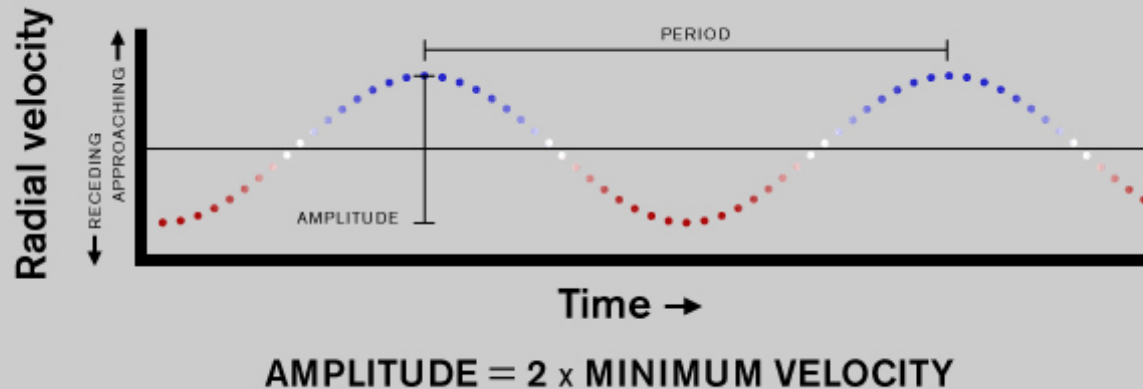
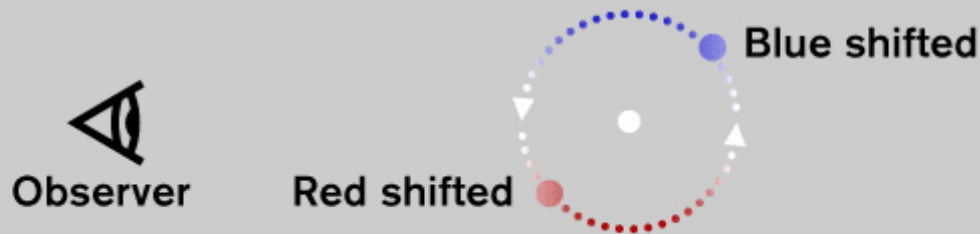


Pulses are modulated at orbital period of 2.09 days

Figure 6. Period variations and occultations of Cen X-3. From Giacconi, 1974.

# Distinguishing BH vs NS

**Determining mass of compact object  
in x-ray binary**



Mass function

$$f_0 \equiv \frac{M_X^3 \sin^3 i}{(M_X + M_0)^2} = \frac{P_0 K_0^3}{2\pi G} \quad f_0 \leq M_X$$

