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LORENTZ BOOSTED NUCLEON-NUCLEON T-MATRIX AND THE TRITON BINDING ENERGY

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The phase equivalent relativistic NN potential, which is related by a nonlinear equation to the original nonrelativistic potential, is used to construct the mass operator (rest Hamiltonian) of the 3- nucleon system. Employing the CD Bonn NN potential, the solution of the relativistic 3N Faddeev equation for ^{3}H shows slightly less binding energy than the corresponding nonrelativistic result. The effect of the Wigner spin rotation on the binding is very small.

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1. Introduction

Considerable experimental effort has been made in measuring proton-deuteron (pd) scattering [1,2,3,4,5] cross sections at intermediate energies. For up to 300 MeV proton energy those data have been analyzed with rigorous three-nucleon (3N) Faddeev calculations [6] based on the CD-Bonn potential [7] and the Tucson-Melbourne 3N force (3NF) [8]. Theoretical predictions based on 2N forces alone are not sufficient to describe the data above about 100 MeV. Some of those defects are known as the Sagara discrepancy [9,10,11]. Though 3NF effects are already seen below 100 MeV, they increase significantly above that energy. However, presently available 3NF's only partially improve the description of cross section data and spin observables. Since most of the cited calculations are based on the non-relativistic formulation of the Faddeev equations [12], one needs to question if in the intermediate energy regime a Poincaré invariant formulation is more adequate.

There are two basic approaches to a relativistic formulation of the 3N problem. One is a manifestly covariant scheme linked to a field theoretical approach [13], the other is based on an exact realization of the symmetry of the Poincaré group in three nucleon quantum mechanics [14]. We employ the second approach, where the mass operator (rest energy operator) consists of relativistic kinetic energies together with two- and many-body interactions including their boost corrections [15].

The first attempt in solving the relativistic Faddeev equation for the 3N bound state based on second approach has been carried out in [16], resulting in a decrease of the binding energy compared to the nonrelativistic result. On the other hand, similar calculations based on the field theory approach [13] increase it. These contradictory results require more investigation. In the following we summarize the results of our calculations based on the second approach: in Section 2 we introduce the relativistic 2N potential, in Section 3 we present the 2N t-matrix, which fulfills the relativistic boosted Lippmann-Schwinger (LS), and in Section 4 we give numerical results for the triton binding energy based on the Poincaré invariant Faddeev equation.

2. The Relativistic Potential

Modern meson theoretical NN potentials, e.g. charge dependent Bonn Potential (CD-Bonn) [7], are derived from a relativistic Lagrangian, then cast into a three-dimensional form using the Blankenbeclar-Sugar equation, which by kinematical redefinitions can be written in the form of a standard nonrelativistic LS equation, which in partial wave decomposed form reads

$$t(p,p';\frac{p'^2}{m}) = v(p,p') + \int_0^\infty \frac{v(p,p')t(p'',p';\frac{p'^2}{m})}{\frac{p'^2}{m} - \frac{p''^2}{m} + i\epsilon} p''^2 dp''. \tag{1}$$

The corresponding relativistic LS equation is given as

$$t^r(p,p';E) = v^r(p,p') + \int_0^\infty \frac{v^r(p,p'')t^r(p'',p';E)}{E - 2\sqrt{p''^2+m^2} + i\epsilon} p''^2 dp'' \tag{2}$$

where

$$E = 2\sqrt{p'^2+m^2} \tag{3}$$

In the relativistic Faddeev equation one needs t^r off-the-energy-shell. According to [17] there is a direct operator relation between the nonrelativistic v and the relativistic v^r:

$$4m\ \hat{v} = 2\sqrt{\hat{p}^2+m^2}\ \hat{v}^r + 2\hat{v}^r\sqrt{\hat{p}^2+m^2} + (\hat{v}^r)^2 \tag{4}$$

In a momentum representation this leads to

$$4m\ v(p,p') = \\ v^r(p,p')\left(2\sqrt{p^2+m^2}+2\sqrt{p'^2+m^2}\right)+\int_0^\infty dp''\ p''^2\ v^r(p,p'')\ v^r(p'',p'). \quad (5)$$

This is the nonlinear relation between the relativistic potential v^r and the nonrelativistic potential v from Eq. (1), which has recently been solved [18]. The resulting on-shell-t-matrix t^r is on-shell identical to the t-matrix t from Eq. (1).

3. The Lorentz Boosted T-matrix

Cluster properties require that the energy is additive. Because of the non-linear relations between the mass and energy in special relativity, the additivity of energies in the rest frame implies a non-linear relation between the two-body interactions in the two and three-body mass operators [14]. We call the two-body interaction in the three-body mass operator the "boosted potential".

$$\hat{v_q^r} \equiv \sqrt{\left[2\sqrt{\hat{p}^2+m^2}+\hat{v^r}\right]^2+q^2}-\sqrt{\left[2\sqrt{\hat{p}^2+m^2}\right]^2+q^2}, \quad (6)$$

where the spectator momentum q in the 3-body center of mass is simultaneously the negative total momentum of the pair. Using Eq. (4) this can be rewritten as [18]

$$4m\ v(p,p') = v_q^r(p,p')\left(\sqrt{4(p^2+m^2)+q^2}+\sqrt{4(p'^2+m^2)+q^2}\right) \\ +\int_0^\infty dp''\ p''^2\ v_q^r(p,p'')v_q^r(p'',p'). \quad (7)$$

Thus one can obtain v_q^r by the same technique [18] as v^r. The boosted off-shell t-matrix is the solution of the LS equation

$$t_q^r(p,p';E_q) = v_q^r(p,p') \\ +\int_0^\infty \frac{v_q^r(p,p')\ t_q^r(p'',p';E_q)}{\sqrt{4(k^2+m^2)+q^2}-\sqrt{4(p''^2+m^2)+q^2}+i\epsilon}p''^2 dp''. \quad (8)$$

with $E_q = \sqrt{4(k^2+m^2)+q^2}$.

In Fig. 1 we display the boosted half-shell ($p'=k$) t-matrix of the CD-Bonn [7] potential at E_{lab}=350 MeV for three different spectator momenta q. The magnitude of the t-matrix gradually decreases with increasing the boost momentum q. It can be shown [19,21] that the half-shell t-matrices $t_q^r(p,p'=k;E_k)$ and $t^r(p,p'=k;E_k)$ are related by simple factors

$$t_q^r(p,k;E_k) = \frac{2\sqrt{p^2+m^2}+2\sqrt{k^2+m^2}}{\sqrt{4(p^2+m^2)+q^2}+\sqrt{4(k^2+m^2)+q^2}}t^r(p,k;E_k). \quad (9)$$

Solving Eqs. (2) and (8) to obtain t^r and t_q^r independently we numerically confirmed the relation (9) with high precision. It is the factor on front of t^r in the right hand side of Eq. (9) which attenuates the amplitude of the t-matrix with increasing q.

It can also be shown [19] that the relativistic half-shell t-matrix t^r is related to the corresponding nonrelativistic one via

$$t^r(p,k;E=2\sqrt{k^2+m^2}) = \frac{4m}{2\sqrt{k^2+m^2}+2\sqrt{p^2+m^2}} t(p,k;k^2/m). \tag{10}$$

The explicit construction of first v^r and then t^r is equivalent to obtaining t^r via resolvent equations as suggested in [19] and carried out in [20,21].

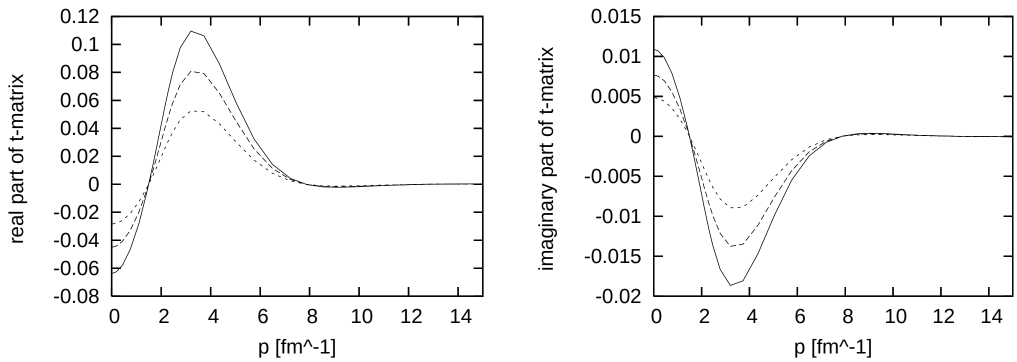

Fig. 1. The boosted half-on-the-mass-shell t-matrix of the CD-Bonn potential at E_{lab}=350 MeV. The left and right plots are real and imaginary parts, respectively. The solid, dashed and dotted lines are related to the boosting momentum q= 0, 10 and 20 fm^{-1}, respectively.

4. The Triton Binding Energy

The relativistic bound state Faddeev equation was solved using the boosted t-matrix t^r_q. In Table 1 the results for the triton binding energy using the CD-Bonn potential as input are shown. The triton binding energy obtained from the relativistic calculation is about 100 keV smaller compared to the one calculated nonrelativistically. This value is significantly smaller than a previously published result [22] in which a reduction of the triton binding energy by about 400 keV was given. The reason for this overestimation of a relativistic effect on the binding energy can be attributed to a different construction of the relativistic off-shell t-matrix t^r. The scaling transformation employed in [22] does not keep the 2N scattering data invariant as function of the 2N c.m. momentum.

We also included the Wigner spin rotation as outlined in [23]. Thereby the the Balian-Brezin method[24] in handling the permutations is quite useful. In Table 2 the triton binding energies are shown allowing charge independence breaking (CIB)[26] and Wigner spin rotations. Wigner spin rotation effects reduce the binding energy by only about 2 keV.

Table 1. The theoretical predictions of the relativistic and nonrelativistic triton binding energies in MeV. The first line indicates how many partial waves we took into account. The second and third lines are the results of the nonrelativistic and relativistic calculations, respectively. The difference between the nonrelativistic and relativistic results is given in the last line. Only the np force of the CD-Bonn potential was used.

	5ch (S-wave)	18ch ($j_{max}=2$)	26ch ($j_{max}=3$)	34ch ($j_{max}=4$)
nonrel.	-8.331	-8.220	-8.241	-8.247
rel.	-8.219	-8.123	-8.143	-8.147
diff.	0.112	0.107	0.098	0.100

Table 2. The theoretical predictions for the relativistic and nonrelativistic triton binding energies in MeV. All numbers are 34 channels results. The second column is the same as the last column in Table 1. The results in the third column take charge dependence[26] into account. In addition the result of the fourth column contains also Wigner spin rotation effects.

	np force only	np+nn forces	Wigner rotation	diff.
nonrel.	-8.247	-8.005	-	-
rel.	-8.147	-7.916	-7.914	-0.002
diff.	0.100	0.089	-	-

5. Summary and Outlook

A phase-shift equivalent 2N potential $\hat{v}^r$ in the relativistic 2N Schrödinger equation is related to the potential v in the nonrelativistic Schrödinger equation by the nonlinear relation given in Eq. (4). The boosted potential $\hat{v}^r_q$ is related to $\hat{v}^r$ by a similar expression, Eq. (6). With these potentials we generate the relativistic fully-off-shell t-matrix t^r_q, which enters into the relativistic Faddeev equation. We solve the relativistic bound state Faddeev equation and compare the binding energy for the triton with the one obtained from a nonrelativistic calculation with the same input interaction. We find that the difference between the two calculations is only about 90 keV including CIB, where the relativistic calculation gives slightly less binding. Taking Wigner spin rotations into account in the relativistic calculation reduces the binding energy by a very small amount, ≈ 2 keV, indicating that Wigner rotations of the spin have essentially no effect on the predicted value of the binding energy.

Applications to the 3-body continuum are in progress. Recently [23] the formulation lined out above has been used to study the low energy A_y puzzle in neutron-deuteron scattering. Details are presented by Witała in this conference. In the intermediate energy regime the formulation has been applied to exclusive proton-deuteron scattering cross sections at 508 MeV [20,21] based on a formulation of of the Faddeev equations which does not employ a partial wave decomposition. The approach can also be extended and applied to electromagnetic processes[27,28].

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