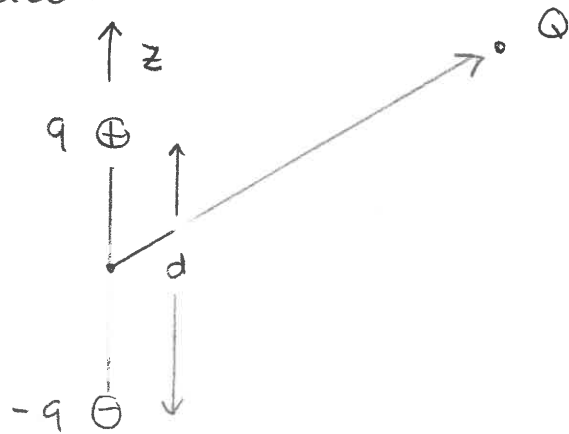


Lecture 3

- ① Force on charge due to electric dipole
- ② Continuous charge distribution
- ③ spherically symmetric charge distribution,
- ④ electric fields



superposition principle +
coulomb's law

$$\begin{aligned}\vec{E}_Q(\vec{r}) &= \vec{E}_{Q,q_1}(\vec{r}) + \vec{E}_{Q,q_2}(\vec{r}) \\ &= kQq \frac{(x,y,z) - (0,0,\frac{d}{2})}{(x^2+y^2+(z-\frac{d}{2})^2)^{3/2}} \\ &\quad + kQ(-q) \frac{(x,y,z) - (0,0,-\frac{d}{2})}{(x^2+y^2+(z+\frac{d}{2})^2)^{3/2}}\end{aligned}$$

Last time we used a trick to write this in a different way

$$\frac{x}{(x^2+y^2+(z \pm \frac{d}{2})^2)^{3/2}} = -\frac{\partial}{\partial x} \frac{1}{(x^2+y^2+(z \pm \frac{d}{2})^2)^{1/2}}$$

$$\frac{y}{(x^2+y^2+(z \pm \frac{d}{2})^2)^{3/2}} = -\frac{\partial}{\partial y} \frac{1}{(x^2+y^2+(z \pm \frac{d}{2})^2)^{1/2}}$$

$$\frac{z \pm d/2}{(x^2+y^2+(z \pm \frac{d}{2})^2)^{3/2}} = -\frac{\partial}{\partial z} \frac{1}{(x^2+y^2+(z \pm \frac{d}{2})^2)^{1/2}}$$

these are the x, y and z components of

$$\frac{(x, y, z) \pm (0, 0, \frac{d}{2})}{(x^2+y^2+(z \pm \frac{d}{2})^2)^{3/2}}$$

this means that we can express

$F_Q(\vec{r})$ as

$$\vec{F} = Qkq \underbrace{\left(-\hat{x} \frac{\partial}{\partial x} - \hat{y} \frac{\partial}{\partial y} - \hat{z} \frac{\partial}{\partial z} \right)}_{-\vec{\nabla}} \left(\frac{1}{\sqrt{x^2+y^2+(z-\frac{d}{2})^2}} - \frac{1}{\sqrt{x^2+y^2+(z+\frac{d}{2})^2}} \right)$$

$\vec{\nabla} f(x,y,z) \equiv \hat{x} \frac{\partial f}{\partial x} + \hat{y} \frac{\partial f}{\partial y} + \hat{z} \frac{\partial f}{\partial z}$ is called

the gradient of $f(\vec{r})$

what can we learn about the dipole if we are very far away? ($|r| \gg d$)

$$\frac{1}{\sqrt{x^2 + y^2 + (z \pm \frac{d}{2})^2}} = \frac{1}{\sqrt{x^2 + y^2 + z^2 \pm zd + \frac{d^2}{4}}} =$$

let $r^2 = x^2 + y^2 + z^2$ $\cos\theta = \frac{z}{r}$

$$\frac{1}{\sqrt{r^2 \pm r^2 \frac{zd}{r^2} + r^2 \frac{d^2}{4r^2}}} = \left(\frac{1}{r} \right) \frac{1}{\sqrt{1 \pm \frac{d}{r} \cos\theta + \frac{d^2}{4r^2}}}$$

↓ Forget

here we assume $\frac{d}{r} \ll 1$ this means $(\frac{d}{r})^2 \ll (\frac{d}{r})$

$$\begin{aligned} \frac{1}{\sqrt{1+u}} &= 1 + \frac{d}{du} \left(\frac{1}{\sqrt{1+u}} \right) \Big|_{u=0} u + o(u^2) \\ &= 1 - \frac{1}{2} \left(\frac{1}{\sqrt{1+u}} \right)' \Big|_{u=0} u + o(u^2) \\ &= 1 - \frac{1}{2} u + o(u^2) \end{aligned}$$

let $u = \pm \frac{d}{r} \cos\theta + \frac{d^2}{4r^2}$

$$\frac{1}{\sqrt{1 \pm \frac{d}{r} \cos\theta + \frac{d^2}{4r^2}}} = 1 - \frac{1}{2} \left(\pm \frac{d}{r} \cos\theta \right) - \frac{1}{2} \left(\frac{d^2}{4r^2} \right) + o \left(\left(\frac{d}{r} \right)^2 \left(\cos\theta + \frac{d}{4r} \right)^2 \right)$$

all much smaller than d/r

$$\frac{1}{\sqrt{1 \pm \frac{d}{r} \cos \theta + \frac{d^2}{4r^2}}} \approx 1 \mp \frac{1}{2} \frac{d}{r} \cos \theta + o\left(\frac{d^2}{r^2}\right)$$

with this approximation $\vec{F}_0(\vec{r})$ becomes

$$\begin{aligned} \vec{F}_0(\vec{r}) &= Qkq(-\vec{\nabla}) \frac{1}{r} \left(1 + \frac{1}{2} \frac{d}{r} \cos \theta - 1 - \left(-\frac{1}{2} \frac{d}{r} \cos \theta\right) + o\left(\frac{d^2}{r^2}\right) \right) \\ &= Qkqd(-\vec{\nabla}) \frac{1}{r} \left(\frac{\cos \theta}{r} \right) \\ &= Qkqd(-\vec{\nabla}) \frac{1}{r} \left(\frac{z}{r^3} \right) \end{aligned}$$

This can be expressed with $\vec{p} = qd\hat{z}$

$$\boxed{\vec{F}_0(\vec{r}) = Q(-\vec{\nabla}) \left(k \frac{\vec{p} \cdot \vec{r}}{r^3} \right) =}$$

we can see that this force depend on the vector \vec{p} , (called the dipole moment)

$$\begin{aligned} &= QkP(-\vec{\nabla}) \left(\frac{z}{(x^2+y^2+z^2)^{3/2}} \right) \\ &= QkP \left(\frac{3xz}{(x^2+y^2+z^2)^{5/2}} \hat{x} + \frac{3yz}{(x^2+y^2+z^2)^{5/2}} \hat{y} + \frac{2z^2-x^2-y^2}{(x^2+y^2+z^2)^{3/2}} \hat{z} \right) \\ &= QkP \frac{3 \vec{r}(\vec{r} \cdot \hat{z}) - r^2 \hat{z}}{r^5} \end{aligned}$$

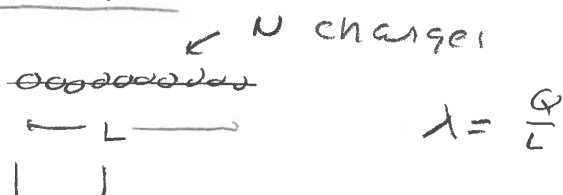
$$\vec{F}_e(\vec{r}) = Qk \frac{3\vec{r}(\vec{r} \cdot \vec{p}) - r^2\vec{p}}{r^5}$$

continuous charge distributions

while electric charge comes in multiples of the charge of an electron, for every day phenomena we are dealing with 10^{18} charges.

It is useful for computations to replace them by continuous charge distributions.

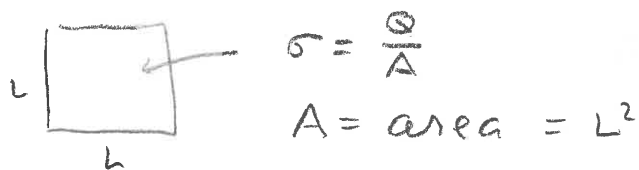
line charges



$\lambda dx = \# \text{ charges between } ||$

$$Q = \int_0^L \lambda dx = Q \frac{L}{L} =$$

$\lambda = \text{linear charge density}$



$$Q = \int_{\square} \sigma dA = \int_0^L dx \int_0^L dy \sigma = L \cdot L \cdot \sigma = Q$$

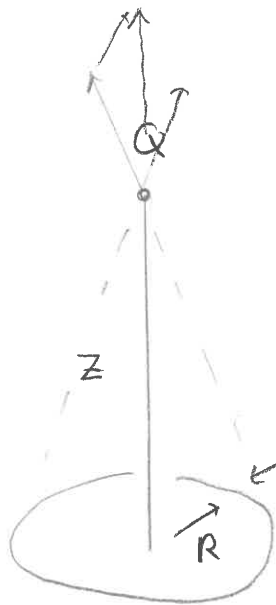
$\sigma =$ surface charge density



$$\rho = \frac{Q}{V} = \frac{Q}{L^3} = \text{volume charge density}$$

$$Q = \int_V \rho dV = \int_0^L dx \int_0^L dy \int_0^L dz \rho = L^3 \frac{Q}{L^3}$$

example



ring of positive charge Q_R

$$\lambda = \frac{Q_R}{2\pi R}$$

choose the origin to be at the center of the circle

$$\vec{r} = (0, 0, z) \quad s = R\theta \quad ds = R d\theta$$

$$\vec{r}'(\theta) = (R \cos \theta, R \sin \theta, 0)$$

$$\vec{F}_0(z) = kQ \int_0^{2\pi} (\lambda R d\theta) \frac{\vec{r} - \vec{r}'(\theta)}{|\vec{r} - \vec{r}'(\theta)|^3}$$

↑
change in $ds = R d\theta$

$$= kQ \int_0^{2\pi} \lambda R d\theta \frac{(-R \cos \theta, -R \sin \theta, z)}{(R^2 \cos^2 \theta + R^2 \sin^2 \theta + z^2)^{3/2}}$$

$$= kQ \lambda R \frac{R}{(R^2 + z^2)^{3/2}} \left(\underbrace{-\int_0^{2\pi} \cos \theta d\theta}_0, \underbrace{-\int_0^{2\pi} \sin \theta d\theta}_0, \frac{z}{R} \int_0^{2\pi} d\theta \right)$$

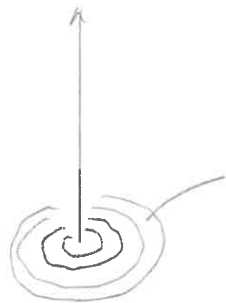
$$= \frac{kQ \lambda R^2 \cdot 2\pi z}{(R^2 + z^2)^{3/2}} (0, 0, 1)$$

using $\lambda = \frac{Q_R}{2\pi R}$

$$\boxed{\vec{F}_0(z) = \frac{kQ Q_R R z}{(R^2 + z^2)^{3/2}} \hat{z}}$$

The net force due to all of the charges on the ring is up as expected.

What if we replace the ring by a disk. We can think of the force due to the disk as the sum of forces due to each ring



$$\sigma = \frac{Q_D}{\pi R^2}$$

$$dQ_R = \underbrace{\sigma \cdot 2\pi r dr}_{\text{area of ring}}$$

$$\vec{F}_{\text{ring}}(z) = \frac{kQ dQ_R r z}{(r^2 + z^2)^{3/2}} = \frac{kQ\sigma \cdot 2\pi r dr \cdot r z}{(r^2 + z^2)^{3/2}}$$

by the superposition principle, the force is the integral from this quantity from $0 \rightarrow R$

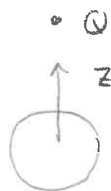
$$\vec{F} = \int_0^R \frac{kQ\sigma \cdot 2\pi r^2 dr}{(r^2 + z^2)^{3/2}} z \cdot \hat{z} =$$

$$= \frac{kQQ_D}{\pi R^2} \cdot 2\pi z \int_0^R \frac{r^2 dr}{(r^2 + z^2)^{3/2}} \hat{z}$$

$$= \frac{2kQQ_D}{R^2} z \int_0^R \frac{r^2 dr}{(r^2 + z^2)^{3/2}} \hat{z}$$

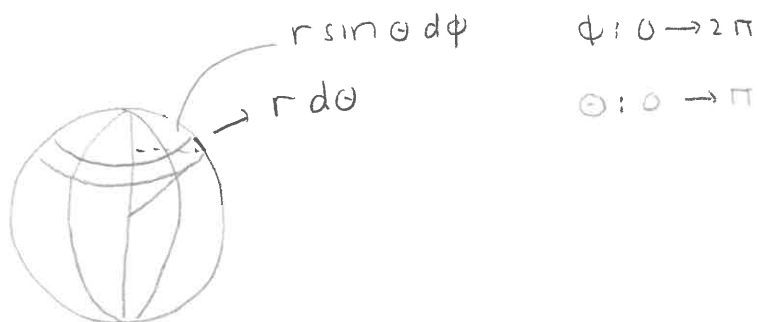
(this integral can be done analytically)

Next we consider a spherical charge distribution $\rho = \rho(r)$



Choose the origin of the coordinate system at the center of the spherical charge. Choose the z axis along the line from the center of the charge to Q.

spherical volume integrals



$$dV = (r d\theta)(r \sin\theta d\phi) dr$$

check $\int_0^R dr \int_0^\pi d\theta \int_0^{2\pi} d\phi \cdot r^2 \sin\theta =$

$$\int_0^R r^2 dr \int_0^\pi \sin\theta d\theta \int_0^{2\pi} d\phi$$

$$\left(\frac{R^3}{3} - \frac{0^3}{3}\right) \underbrace{(-\cos\pi + \cos 0)}_2 (2\pi - 0) = \frac{4}{3} \pi R^3$$

which is the volume of a sphere

$$\vec{r} = (0 \ 0 \ z)$$

$$\vec{r}' = (x \ y \ z) = (r \sin \theta \cos \phi, r \sin \theta \sin \phi, r \cos \theta)$$

$$\begin{aligned} |\vec{r} - \vec{r}'|^2 &= r^2 \sin^2 \theta \cos^2 \phi + r^2 \sin^2 \theta \sin^2 \phi \\ &\quad + z^2 - 2zr \cos \theta + r^2 \cos^2 \theta \\ &= r^2 \sin^2 \theta + r^2 \cos^2 \theta + z^2 - 2rz \cos \theta \\ &= r^2 + z^2 - 2rz \cos \theta \end{aligned}$$

$$d\omega = \rho \, dV = \rho \, r^2 \sin \theta \, d\theta \, d\phi \, dr$$

$$\begin{aligned} F_z(z) &= kq \int_0^R dr \int_0^\pi d\theta \int_0^{2\pi} d\phi \, \rho(r) r^2 \sin \theta \times \\ &\quad \frac{(-r \sin \theta \cos \phi, -r \sin \theta \sin \phi, z - r \cos \theta)}{(r^2 + z^2 - 2rz \cos \theta)^{3/2}} \end{aligned}$$

this is just the superposition principle. Integrating each component separately

$$\int_0^{2\pi} \cos \phi \, d\phi = \int_0^{2\pi} \sin \phi \, d\phi = 0$$

so the x and y components of the flux vanish

what remains is in the z direction -
there is a 2π from the ϕ integral

$$F_{z0}(z) = \rho Q 2\pi \int_0^R \rho(r) r^2 dr \int_0^\pi \frac{z - r \cos \theta}{(r^2 + z^2 - 2rz \cos \theta)^{3/2}} \sin \theta d\theta$$

To do the integral use

$$\begin{aligned} - \frac{d}{dz} \frac{1}{(r^2 + z^2 - 2rz \cos \theta)^{1/2}} &= -(-\frac{1}{2}) \frac{2z - 2r \cos \theta}{(r^2 + z^2 - 2rz \cos \theta)^{3/2}} \\ &= \frac{z - r \cos \theta}{(r^2 + z^2 - 2rz \cos \theta)^{3/2}} \end{aligned}$$

using this in the integral

$$F_{z0}(z) = 2\pi R Q \int_0^R \rho(r) r^2 \left(-\frac{d}{dz}\right) \int_0^\pi \frac{\sin \theta d\theta}{(r^2 + z^2 - 2rz \cos \theta)^{1/2}}$$

$$\text{let } v = r^2 + z^2 - 2rz \cos \theta \quad (\text{replace } \theta)$$

$$dv = 2rz \sin \theta d\theta \quad v: |r-z|^2 \rightarrow |r+z|^2$$

$$\sin \theta d\theta = \frac{dv}{2rz}$$

$$\bar{F}_{z0}(z) = 2\pi R Q \int_0^R \rho(r) r^2 \left(-\frac{d}{dz}\right) \int_{|z-r|}^{|z+r|} \left(\frac{1}{2rz}\right) \frac{dv}{v^{1/2}}$$

$$= 2\pi R Q \int_0^R \rho(r) r \left(-\frac{d}{dz}\right) \left(\frac{1}{z} (|z+r| - |z-r|)\right)$$

$$z > r \rightarrow 2\frac{r}{z^2}$$

$$z < r \quad \frac{d}{dz}(z) = 0$$

case 1 $z > R$

$$F = \frac{4\pi RQ}{z^2} \int \rho(r) r^2 dr = \frac{kQ_{\text{sphere}}}{z^2}$$

case 2 $z = r < R$

$$F = \frac{4\pi RQ}{z^2} \int_0^z \rho(r) r^2 dr = \frac{kQ}{z^2} \text{ charge inside sphere of radius } z$$

conclusion:

* spherical charge distribution gives the same force as a point charge at the origin with the same charge.

(note z is just the distance between the particle and the center of the sphere.)

* there are no forces on a charge inside of a hollow sphere.

electric fields

In coulombs law

$$\vec{F}_Q(\vec{r}) = \sum k Q q_i \frac{\vec{r} - \vec{r}_i}{|\vec{r} - \vec{r}_i|^3} = Q \left(k \sum q_i \frac{\vec{r} - \vec{r}_i}{|\vec{r} - \vec{r}_i|^3} \right)$$

The quantity in () only depends on \vec{r} and the other charges. It is a vector called the electric field at \vec{r}

The above equation can be expressed as

$$\vec{F}_Q(\vec{r}) = Q \vec{E}(\vec{r})$$

$$\vec{E}(\vec{r}) = \frac{\vec{F}_Q(\vec{r})}{Q}$$

is called the electric field at \vec{r}

* It exists in the absence of the charge Q

* To measure the electric field at \vec{r}

$$\vec{E}(\vec{r}) = \lim_{Q \rightarrow 0} \frac{\vec{F}_Q(\vec{r})}{Q}$$

This prevents Q from distributing
the position of the other charges

$\vec{E}(\vec{r})$ assigns a vector at each point
in space. The direction is the
direction that a positive charge
would move if placed at that
point



hollow sphere - field outside
is radial, no field inside