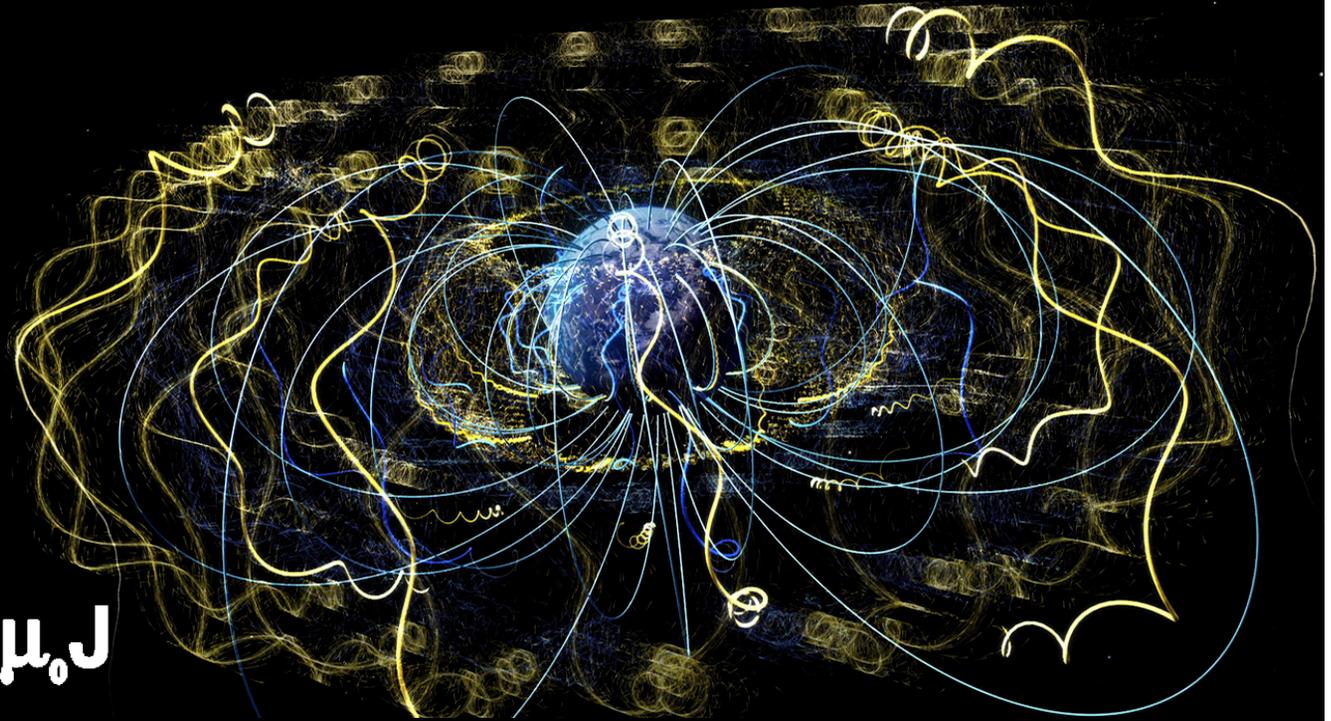


$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} + \mu_0 \mathbf{J}$$



# Electricity and Magnetism II: 3812

Professor Jasper Halekas  
Van Allen 70  
MWF 9:30-10:20 Lecture

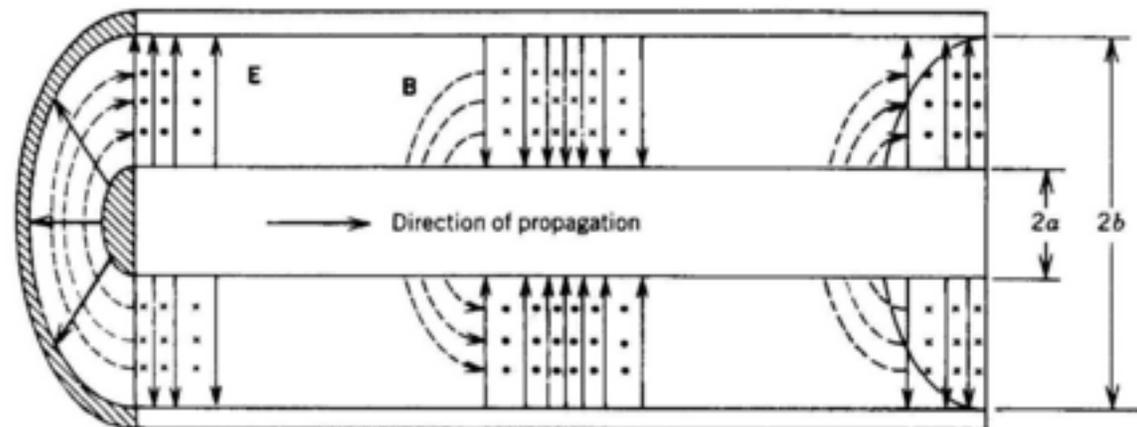
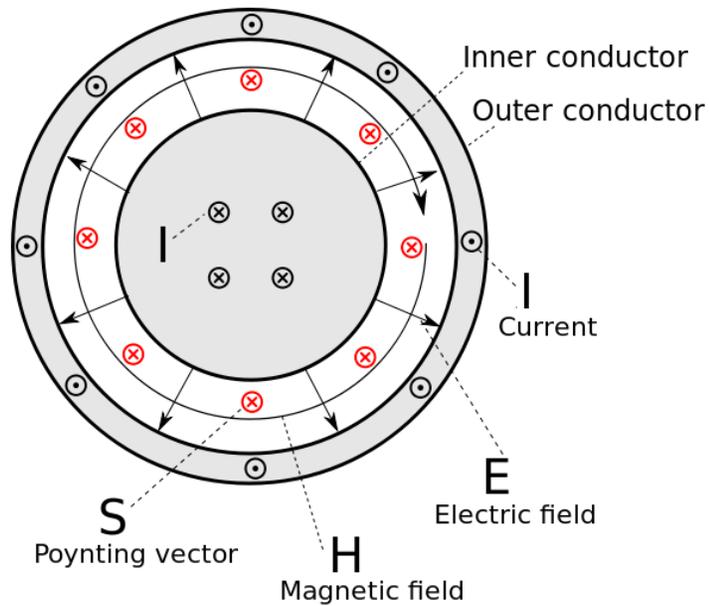
# XKCD with the Best Take as Usual



# Announcements

- It appears very likely that we will be going to online-only classes, most likely after spring break
  - This class will continue, in a virtual format if need be
    - Lecture notes etc. will be posted as usual
    - Lectures will be either recorded or webcast (TBD)
      - Preferences?
  - Homework will have to be turned in electronically
    - Is this an issue for anyone?
  - Exams will likely have to be take-home, open book, with a constrained duration
    - Is this an issue for anyone?
- Suggestions to ease the transition are welcome!

# Coaxial TEM Waveguide



# Coaxial Wave Guide



- Coaxial cable supports TEM modes

$$\vec{E}(\vec{r}, t) = \vec{E}_0(x, y) e^{i(\kappa z - \omega t)}$$
$$\vec{B}(\vec{r}, t) = \vec{B}_0(x, y) e^{i(\kappa z - \omega t)}$$

$$\text{TEM} \Rightarrow \begin{aligned} \vec{E}_0 &= E_{0x} \hat{x} + E_{0y} \hat{y} \\ \vec{B}_0 &= B_{0x} \hat{x} + B_{0y} \hat{y} \end{aligned}$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$
$$\Rightarrow \begin{aligned} \frac{\partial E_{0y}}{\partial x} - \frac{\partial E_{0x}}{\partial y} &= 0 \\ -i\kappa E_{0y} &= i\omega B_{0x} \\ i\kappa E_{0x} &= i\omega B_{0y} \end{aligned}$$

$$\nabla \times \vec{B} = \frac{1}{c^2} \frac{\partial \vec{E}}{\partial t}$$
$$\Rightarrow \begin{aligned} \frac{\partial B_{0y}}{\partial x} - \frac{\partial B_{0x}}{\partial y} &= 0 \\ -i\kappa B_{0y} &= -\frac{i\omega}{c^2} E_{0x} \\ i\kappa B_{0x} &= -\frac{i\omega}{c^2} E_{0y} \end{aligned}$$

Combine to find:

$$E_{0x} = c B_{0y}, \quad E_{0y} = -c B_{0x}$$

$$\& \quad k = \frac{\omega}{c}$$

$$\frac{\partial E_{0y}}{\partial x} - \frac{\partial E_{0x}}{\partial y} = 0$$

$$\Rightarrow \frac{\partial B_{0x}}{\partial x} + \frac{\partial B_{0y}}{\partial y} = 0$$

$$\frac{\partial B_{0y}}{\partial x} - \frac{\partial B_{0x}}{\partial y} = 0$$

$$\Rightarrow \frac{\partial E_{0x}}{\partial x} + \frac{\partial E_{0y}}{\partial y} = 0$$

Summary:  $B_{0z}, E_{0z} = 0$

$$\begin{array}{l|l} \nabla \cdot \vec{B}_0 = 0 & \nabla \cdot \vec{E}_0 = 0 \\ \nabla \times \vec{B}_0 = 0 & \nabla \times \vec{E}_0 = 0 \end{array}$$

Solutions: - Cylindrical symmetry  
- Vacuum

- Boundary conditions  $\vec{E}_\perp = 0$   
&  $B_\perp = 0$  at  $r=a, b$

$$\Rightarrow \vec{E}_0(s, \varphi) = \frac{A}{s} \hat{s}$$

$$B = \frac{A}{c}$$

$$\vec{B}_0(s, \varphi) = \frac{B}{s} \hat{\varphi}$$

Finally

$$\vec{E}(s, \varphi, z, t) = \frac{A}{s} \cos(\kappa z - \omega t + \delta) \hat{s}$$

$$\vec{B}(s, \varphi, z, t) = \frac{A}{cs} \cos(\kappa z - \omega t + \delta) \hat{\varphi}$$

$$\vec{S} = \frac{\vec{E} \times \vec{B}}{\mu_0}$$

$$= \frac{A^2}{cs^2 \mu_0} \cos^2(\kappa z - \omega t + \delta) \hat{z}$$

propagates along coax

# Check Your Understanding I

- Consider a wave with an electric field amplitude  $E_0$  in vacuum, and a wave with the same electric field amplitude  $E_0$  in a dielectric material with permittivity  $\epsilon > \epsilon_0$ 
  - In which case is the electric field energy density greater?
  - Physically, why is the electric field energy density greater in that case?

$$Q/\therefore U_E = \frac{1}{2} \epsilon E_0^2$$

$$U_{E_{vac}} = \frac{1}{2} \epsilon_0 E_0^2$$

$$U_E > U_{E_{vac}}$$

$$\begin{array}{l} \longrightarrow \vec{E}_{imposed} \\ \longrightarrow \vec{P} = \epsilon_0 \chi_e \vec{E}_{imp} \\ \longleftarrow \Delta \vec{E} \\ \longrightarrow \vec{E}_0 = \vec{E}_{imp} - \Delta \vec{E} \end{array}$$

- Field weakened by polarization
- Wave energy goes with imposed field, not weakened field

# Check Your Understanding II

- Consider a wave with a magnetic field amplitude  $B_0$  in vacuum, and a wave with the same magnetic field amplitude  $B_0$  in a dielectric material with permeability  $\mu > \mu_0$ 
  - In which case is the magnetic field energy density greater?
  - Physically, why is the magnetic field energy density greater in that case?

Q2:  $U_B = B_0^2 / 2\mu$   
 $U_{Bvac} = B_0^2 / 2\mu_0$

$U_B < U_{Bvac}$

$\longrightarrow \vec{B}_{imposed}$

$\longrightarrow \vec{M} = \chi_m \vec{B}_{imp} / \mu$

$\longrightarrow \Delta \vec{B}$

$\longrightarrow \vec{B}_0 = \vec{B}_{imp} + \Delta \vec{B}$

- Field amplified by magnetization
- Wave energy goes with imposed field, not amplified field

---

Note: For wave  $B_0/E_0 = \eta/\omega$   
 $= \sqrt{\mu\epsilon} > 1$

- This factor ensures that  $U_B = U_E$  in both vacuum & dielectric

# Check Your Understanding III

- Why are electromagnetic waves damped in conducting materials?
  - A. Free currents are driven by the electric fields
  - B. Faraday's Law
  - C. Ampere's Law
  - D. Ohm's Law
  - E. All of the above

Q3:

Ohm's Law  $\vec{J} = \sigma \vec{E}$

Faraday's Law  $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$

Ampere's Law  $\nabla \times \vec{B} = \mu \epsilon \frac{\partial \vec{E}}{\partial t}$   
Becomes  $\nabla \times \vec{B} = \mu \epsilon \frac{\partial \vec{E}}{\partial t} + \mu \vec{J}$

Wave: Changing  $B \rightarrow E$   
Changing  $E \rightarrow B$

Damped Wave: Changing  $B \rightarrow E$   
 $E \rightarrow \vec{J}$   
Changing  $E, \vec{J} \rightarrow B$

Dissipation  $\vec{J} \cdot \vec{E} = \sigma E^2 = \frac{dW}{dt}$   
energy loss to particles  
from fields