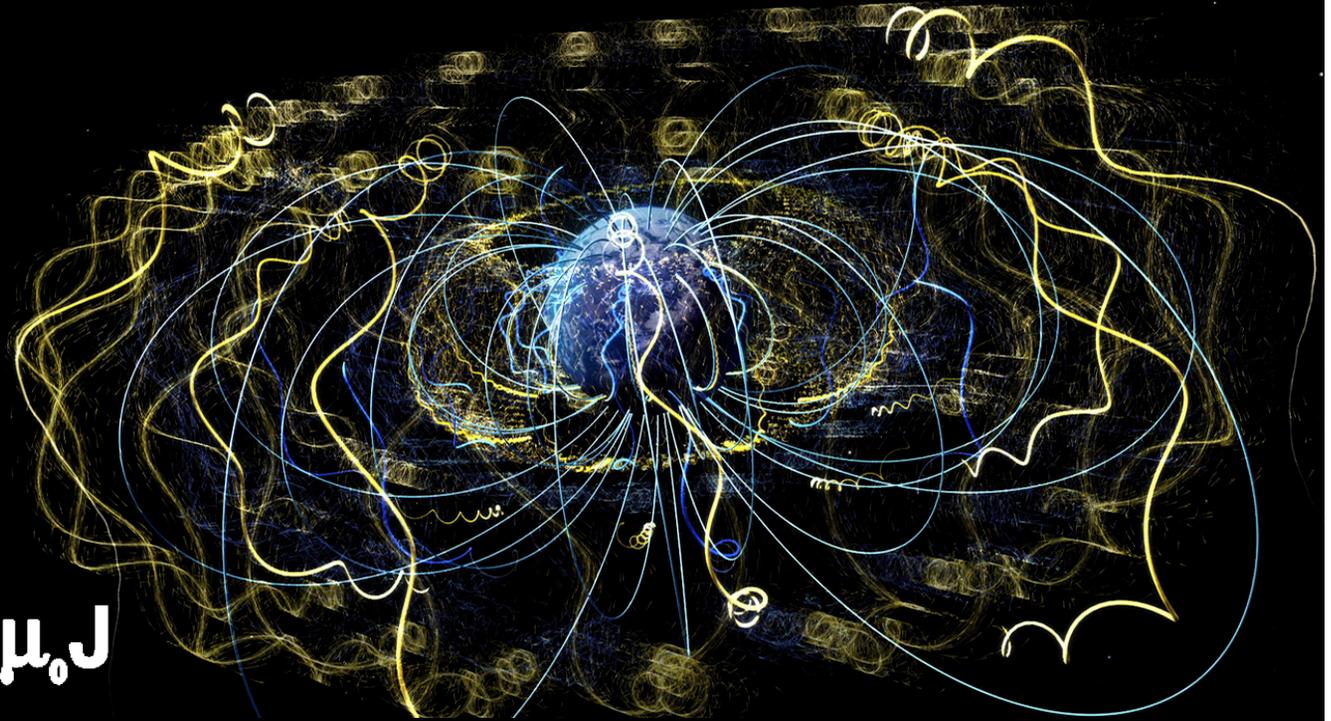


$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} + \mu_0 \mathbf{J}$$



Electricity and Magnetism II: 3812

Professor Jasper Halekas
Van Allen 70
MWF 9:30-10:20 Lecture

Announcements

- Our university is transitioning to online-only instruction, starting March 23 (first Monday following spring break), and continuing until at least April 3 (two weeks)
 - This class will continue in a virtual format
 - Lectures will be webcast via Zoom (at least on a trial basis)
 - Lecture notes will be posted as usual
 - Office hours will be virtual via Zoom
 - Homework #7 and #8 will have to be turned in electronically
 - Due dates/times remain the same
 - There are no exams in this class during the online-only period
 - If the online-only period is eventually extended, we will have take-home exams. If that occurs, more information will be forthcoming.

Zoom Information

- I plan to utilize Zoom for lectures and office hours
 - Zoom is an online videoconferencing tool which allows me to share audio, video, and anything on my computer screen
 - Lectures will be recorded and posted on ICON for those who can't make the lecture, or want to review the lecture
 - Office hours will not be recorded, and will utilize a "waiting room" to allow private discussion if necessary
- Zoom meeting # for lecture: 372-343-761
 - Join URL: <https://uiowa.zoom.us/j/372343761>
- Zoom meeting # for office hours: 146-384-269
 - Join URL: <https://uiowa.zoom.us/j/146384269>
- You will be able to access these meetings directly through these join URLs, or through the Zoom links on the course ICON page

Ch. 10 | Potentials & Fields

10.1 | Potential Formulation

Maxwell's Eqs.

$$\begin{aligned}\nabla \cdot \vec{E} &= \rho/\epsilon_0 & \nabla \times \vec{E} &= -\partial \vec{B} / \partial t \\ \nabla \cdot \vec{B} &= 0 & \nabla \times \vec{B} &= \mu_0 \epsilon_0 \partial \vec{E} / \partial t + \mu_0 \vec{J}\end{aligned}$$

Magnetostatics: $\nabla \cdot \vec{B} = 0$
 $\Rightarrow \boxed{\vec{B} = \nabla \times \vec{A}}$
still valid!

Electrostatics: $\nabla \times \vec{E} = 0$
 $\Rightarrow \vec{E} = -\nabla V$
no longer true!

$$\begin{aligned}\text{Fix: } \nabla \times \vec{E} &= -\partial \vec{B} / \partial t \\ &= -\partial / \partial t (\nabla \times \vec{A}) \\ &= \nabla \times (-\partial \vec{A} / \partial t)\end{aligned}$$

$$\Rightarrow \nabla \times (\vec{E} + \partial \vec{A} / \partial t) = 0$$

$$\Rightarrow \vec{E} + \partial \vec{A} / \partial t = -\nabla V$$

$$\Rightarrow \boxed{\vec{E} = -\nabla V - \partial \vec{A} / \partial t}$$

What about Gauss's Law & Ampere's Law?

$$\text{Gauss: } \nabla \cdot \vec{E} = \nabla \cdot (-\nabla V) + \nabla \cdot \left(-\frac{\partial \vec{A}}{\partial t}\right) = \rho/\epsilon_0$$

$$\Rightarrow \boxed{\nabla^2 V + \frac{\partial}{\partial t}(\nabla \cdot \vec{A}) = -\rho/\epsilon_0}$$

Poisson's Eq. w/ extra term

$$\begin{aligned} \text{Ampere: } \nabla \times \vec{B} &= \nabla \times (\nabla \times \vec{A}) \\ &= \nabla(\nabla \cdot \vec{A}) - \nabla^2 \vec{A} \\ &= \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial}{\partial t}(-\nabla V - \frac{\partial \vec{A}}{\partial t}) \end{aligned}$$

$$\Rightarrow \boxed{\nabla^2 \vec{A} - \mu_0 \epsilon_0 \frac{\partial^2 \vec{A}}{\partial t^2} - \nabla(\nabla \cdot \vec{A} + \mu_0 \epsilon_0 \frac{\partial V}{\partial t}) = -\mu_0 \vec{J}}$$

Ugh!

Gauge Transformations

- would like to simplify - - -

- Anything that leaves \vec{E}, \vec{B} unchanged is okay

- E.g. adding a constant to V

Try $\vec{A}' = \vec{A} + \vec{\alpha}$ w/ $\vec{\alpha}, \beta$ functions
 $V' = V + \beta$

Must have $\nabla \times \vec{A}' = \nabla \times \vec{A}$

so $\nabla \times \vec{\alpha} = 0$

$\Rightarrow \vec{\alpha} = \nabla \lambda$ w/ λ a scalar function

Must have $-\nabla V' - \frac{\partial \vec{A}'}{\partial t} = -\nabla V - \frac{\partial \vec{A}}{\partial t}$

so $\nabla \beta + \frac{\partial \vec{\alpha}}{\partial t} = 0$

$\Rightarrow \nabla \beta + \frac{\partial}{\partial t} (\nabla \lambda) = 0$

$\Rightarrow \nabla (\beta + \frac{\partial \lambda}{\partial t}) = 0$

$\Rightarrow \beta + \frac{\partial \lambda}{\partial t} = f(t)$ ind. of position

We'll pick $f(t) = 0$

$\Rightarrow \beta = -\frac{\partial \lambda}{\partial t}$

So any

$$\boxed{\begin{aligned}\vec{A}' &= \vec{A} + \nabla \lambda \\ V' &= V - \frac{\partial \lambda}{\partial t}\end{aligned}}$$

Gives same \vec{E} and \vec{B}

Coulomb Gauge

$$\text{Pick } \nabla \cdot \vec{A} = 0$$

$$\Rightarrow \nabla^2 V = -\rho/\epsilon_0 \quad \text{Poisson's Eq.}$$

$$\text{But } \nabla^2 \vec{A} - \mu_0 \epsilon_0 \frac{\partial^2 \vec{A}}{\partial t^2} - \nabla \mu_0 \epsilon_0 \frac{\partial V}{\partial t} = -\mu_0 \vec{J}$$

yuck!

Note: V in Coulomb gauge only depends on instantaneous charge distribution. But \vec{A} (& \vec{B} and \vec{E}) have propagation delays since causal influences propagate at speed of light

Lorentz Gauge

$$\text{Pick } \nabla \cdot \vec{A} = -\mu_0 \epsilon_0 \frac{\partial V}{\partial t}$$

$$\Rightarrow \nabla^2 V - \mu_0 \epsilon_0 \frac{\partial^2 V}{\partial t^2} = -\rho/\epsilon_0$$

$$\nabla^2 \vec{A} - \mu_0 \epsilon_0 \frac{\partial^2 \vec{A}}{\partial t^2} = -\mu_0 \vec{J}$$

Symmetric in V, \vec{A} and ρ, \vec{J}

Define d'Alembertian operator

$$\square^2 = \nabla^2 - \mu_0 \epsilon_0 \frac{\partial^2}{\partial t^2}$$

$$\square^2 V = -\rho/\epsilon_0$$

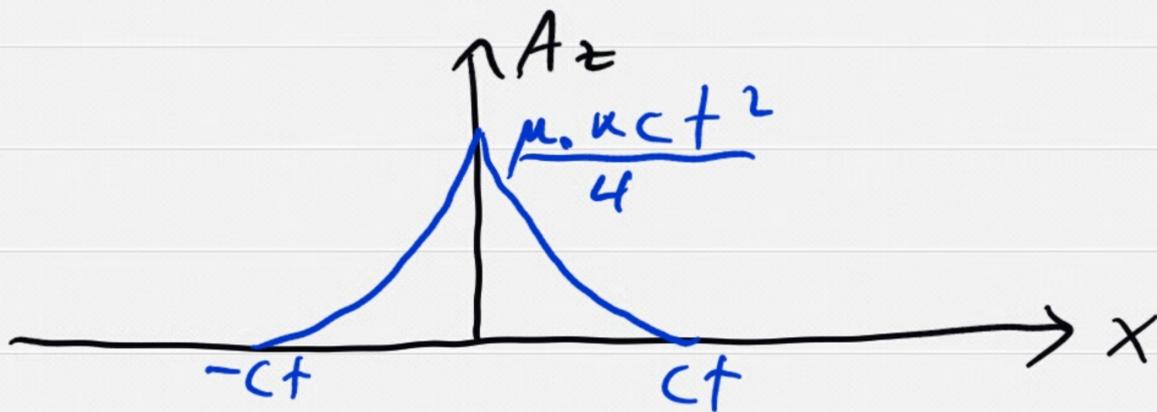
$$\square^2 \vec{A} = -\mu_0 \vec{J}$$

w/ $\rho, \vec{J} = 0 \Rightarrow$ Wave Equation

w/ $\rho, \vec{J} \neq 0 \Rightarrow$ Inhomogeneous
Wave Equation

Example:

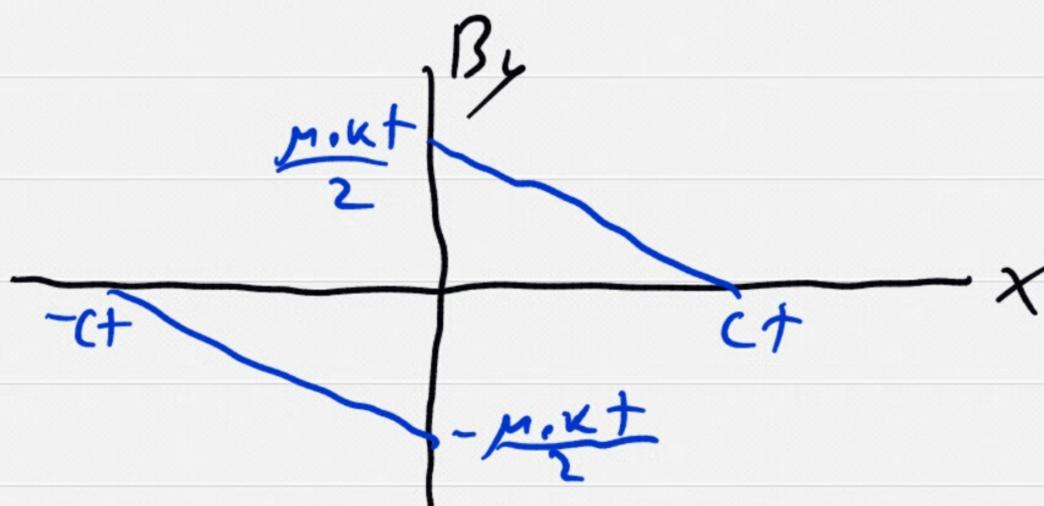
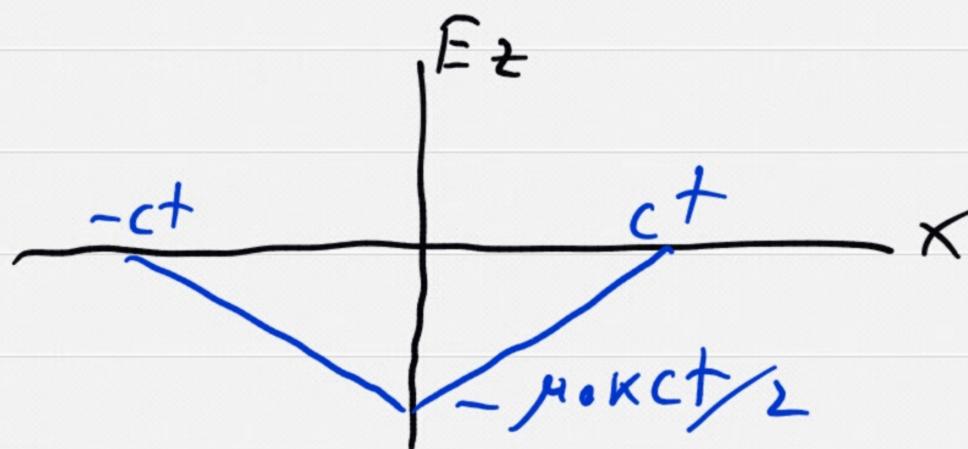
$$V = 0, \quad \vec{A} = \begin{cases} \frac{\mu_0 k}{4c} (ct - |x|)^2 \hat{z} & |x| < ct \\ 0 & |x| > ct \end{cases}$$



$$\begin{aligned} \nabla \cdot \vec{A} &= \frac{\partial A_z}{\partial z} = 0 \\ \frac{\partial V}{\partial t} &= 0 \end{aligned} \quad \Rightarrow \quad \begin{array}{l} \text{satisfies} \\ \text{both Lorenz} \\ \& \text{ Coulomb} \end{array}$$

$$\vec{E} = -\frac{\partial \vec{A}}{\partial t} = -\frac{\mu_0 k}{2} (ct - |x|) \hat{z} \quad |x| < ct$$

$$\begin{aligned} \vec{B} &= \nabla \times \vec{A} = -\frac{\partial A_z}{\partial x} \hat{y} \\ &= \frac{\mu_0 k}{2c} (ct - |x|) \cdot \text{sign}(x) \hat{y} \quad |x| < ct \end{aligned}$$



\vec{B} has a discontinuity
at $x = 0$

$$\Delta \vec{B}_{||} = \mu_0 \vec{k}_f \times \hat{y} = \mu_0 \vec{k}_f \times \hat{x} \\ = \mu_0 \kappa t \hat{y}$$

$$\Rightarrow \boxed{\vec{k}_f = \kappa t \hat{z}}$$

surface current increasing w/ time

$$\nabla \cdot \vec{E} = 0 \Rightarrow \rho_f = 0 \\ \vec{E} \text{ continuous} \Rightarrow \sigma_f = 0$$

\vec{E} completely induced

- Note that current starts
flowing at $t = 0$

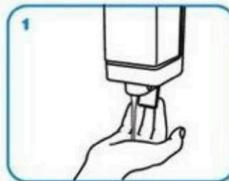
- Fields don't respond until
 $t = |x|/c$

- Information propagates @ c

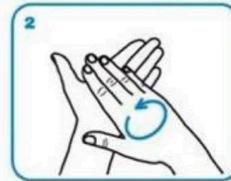
Happy Spring Break – And Wash Your Hands!



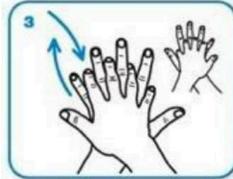
Three Rings for the Elven-kings



under the sky



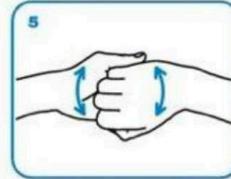
Seven for the dwarf-lords



in their halls of stone



Nine for Mortal Men doomed to die



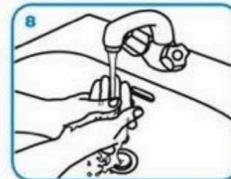
One for the Dark Lord on his dark throne



In the Land of Mordor where the Shadows lie



One Ring to rule them all



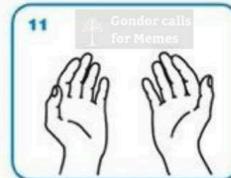
One Ring to find them



One Ring to bring them all



and in the darkness bind them



In the Land of Mordor where the Shadows lie