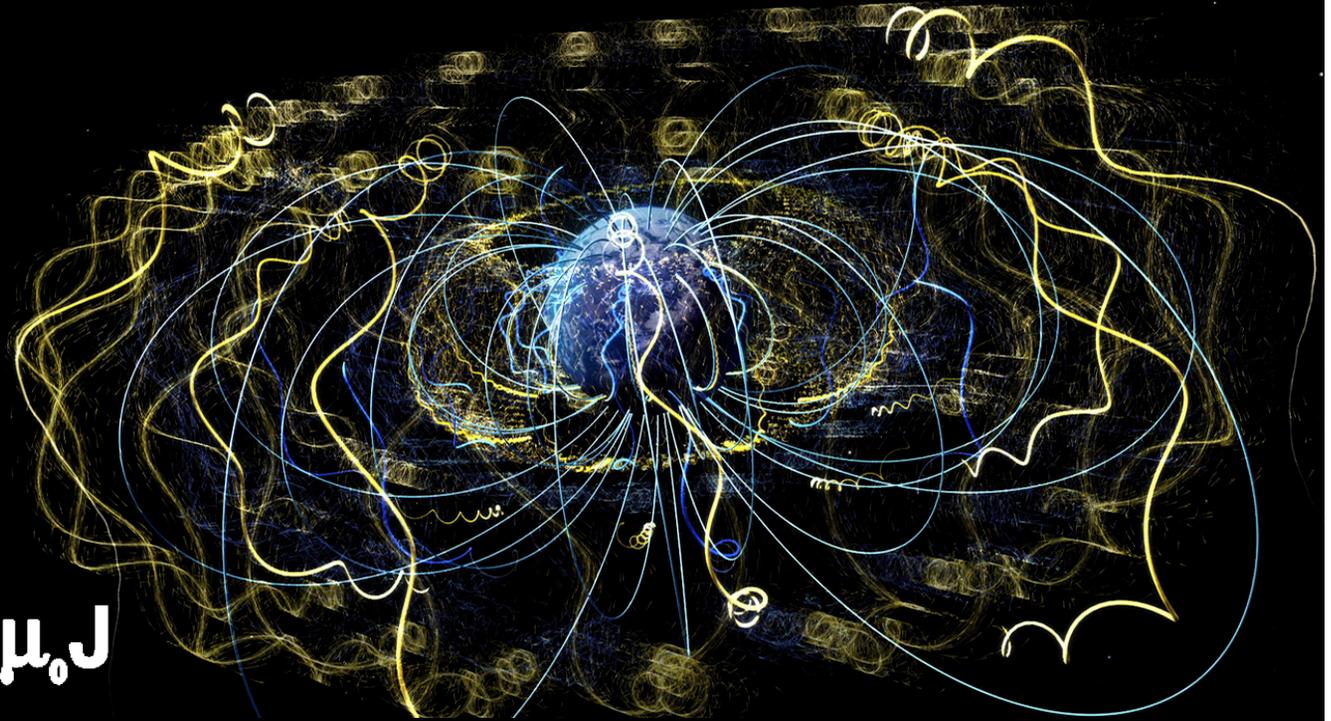


$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} + \mu_0 \mathbf{J}$$



Electricity and Magnetism II: 3812

Professor Jasper Halekas
Van Allen 70
MWF 9:30-10:20 Lecture

Announcements

- We will be in an online format for the rest of the semester
 - We will hold both lectures and office hours virtually using Zoom
 - Homework assignments should be turned in at the usual time/date in electronic form by e-mail
 - The remaining exams will be take-home open-book format
 - The course schedule and syllabus have been modified to reflect these and other changes

Key Syllabus/Schedule Changes

- Lectures and Office Hours
 - Moved to virtual format
- Material Covered
 - Coverage of Ch. 12 reduced to “selected topics” due to extended spring break
- Homework
 - Electronic turn-in rather than physical
 - Number of assignments reduced (10 instead of 11)
 - Delayed by one week due to extended spring break
- Exams
 - Midterm 2 delayed by one week due to extended spring break
 - Midterm 2 and Final both changed to take-home open-book format with two-hour duration

Revised Schedule

REVISED VERSION 3/18/2020

Physics 3812 Electricity and Magnetism II 2020 Schedule

Dates	Week	Reading (Due Monday unless noted)	HW Due Friday	Notes
Jan. 20-24	Week 1	Ch. 7.1	No HW	<i>Holiday Monday 1/20</i>
Jan. 27-31	Week 2	Ch. 7.2	HW 1	
Feb. 3-7	Week 3	Ch. 7.3	HW 2	
Feb. 10-14	Week 4	Ch. 8	HW 3	
Feb. 17-21	Week 5	Ch. 9.1-9.2	HW 4	
Feb. 24-28	Week 6	No Reading	No HW	Midterm 1 Ch. 7-9.2 Wednesday Feb. 26
Mar. 2-6	Week 7	Ch. 9.3-9.4	HW 5	
Mar. 9-13	Week 8	Ch. 9.5-10.1	HW 6	
Mar. 16-20	Spring Break	No Reading	No HW	Spring Break, Woohoo!
Mar. 23-27	Extended Spring Break	No Reading	No HW	COVID-19, isn't it fun?
Mar. 30-Apr. 3	Week 9	Ch. 10.2-10.3	HW 7	
Apr. 6-10	Week 10	Ch. 11.1	HW 8	
Apr. 13-17	Week 11	Ch. 11.2	HW 9	
Apr. 20-24	Week 12	No Reading	No HW	Midterm 2 Ch. 9.3-11 Wednesday Apr 22
Apr. 27-May 1	Week 13	Ch. 12.1-12.3 (Selected Topics)	HW 10	
May 4-8	Week 14	No Reading	No HW	
May 11-15	Finals Week	No Reading	No HW	Final Exam Ch. 7-12 Monday May 11

Discussion of Values

- The next few weeks (or months) are going to be a learning experience for us all
- Communication will be even more important during this time period
 - Please continue to feel free to ask questions during lecture
 - You can virtually “raise your hand” if you have a question
 - You can use the “chat” feature to ask a question by text
 - You can also just un-mute and speak as needed
 - Please raise any other issues or concerns you have, either during virtual office hours or by e-mail

10.2 | Continuous Distributions

Retarded Potentials

Recall Eqs. for V, \vec{A}
in Lorenz Gauge

$$\square^2 \vec{A} = \nabla^2 \vec{A} - \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2} = -\mu_0 \vec{J}$$

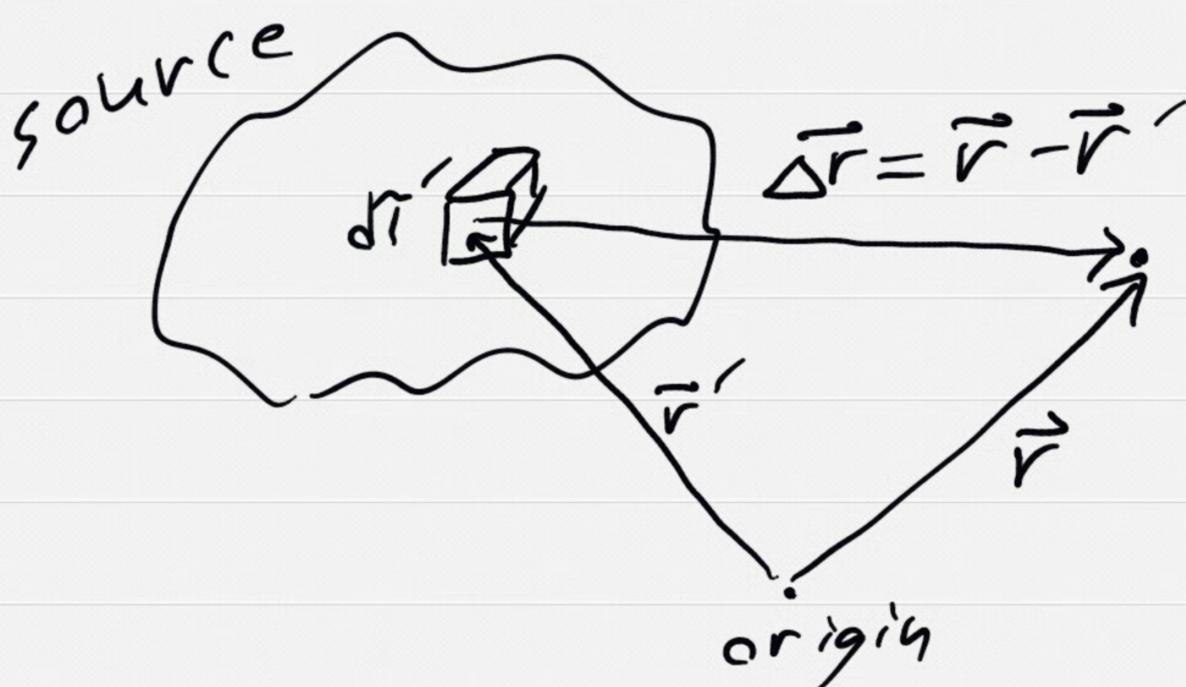
$$\square^2 V = \nabla^2 V - \frac{1}{c^2} \frac{\partial^2 V}{\partial t^2} = -\rho/\epsilon_0$$

- Static Case

$$\nabla^2 \vec{A} = -\mu_0 \vec{J}, \quad \nabla^2 V = -\rho/\epsilon_0$$

$$\Rightarrow \vec{A}(\vec{r}) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}')}{\Delta r} d\tau'$$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{V(\vec{r}')}{\Delta r} d\tau'$$



But what if \vec{r}' changes??

Retarded Potentials

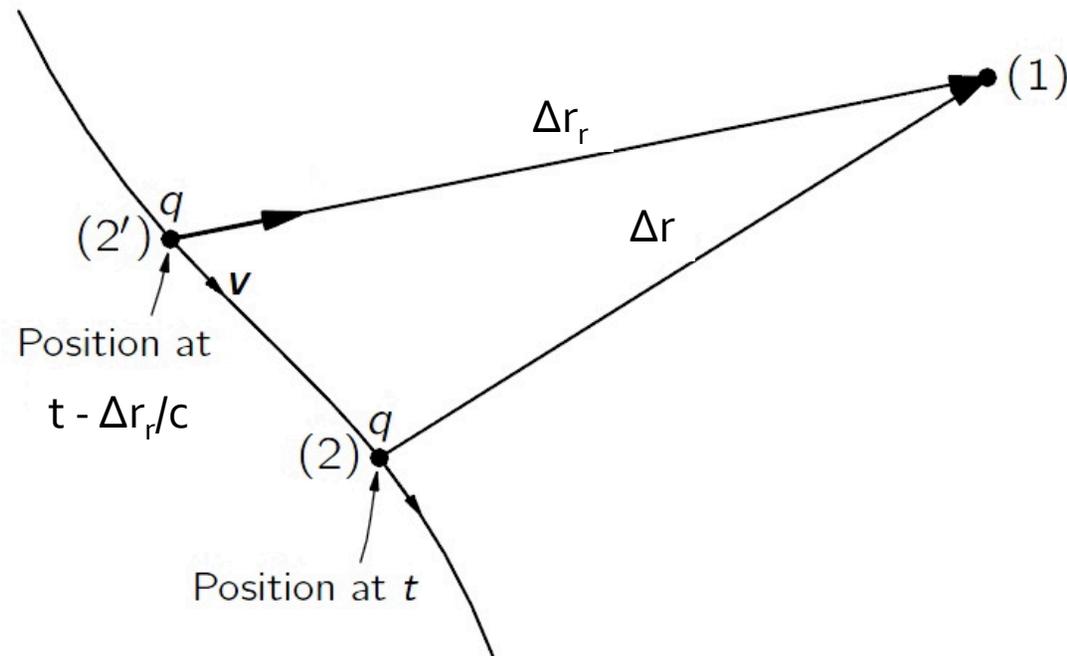


Fig. 21-1. The fields at (1) at the time t depend on the position (2') occupied by the charge q at the time $t - \Delta r_r / c$

- Information about source takes time to propagate at the speed of light

$$\Delta t = \Delta r / c$$

- Effect seen at time t at \vec{r} happened at $t_r = t - \Delta r / c$ at \vec{r}'

- t_r = "retarded time"

$$\text{Guess: } V(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}', t_r)}{\Delta r} d\tau'$$
$$\vec{A}(\vec{r}, t) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}', t_r)}{\Delta r} d\tau'$$

In other words, potentials depend on "old" distribution at sources

We'll prove it for V

$$\nabla V = \nabla \left(\frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}', t_r)}{\Delta r} d\tau' \right)$$
$$= \frac{1}{4\pi\epsilon_0} \int \nabla \left(\frac{\rho(\vec{r}', t_r)}{\Delta r} \right) d\tau'$$

$$= \frac{1}{4\pi\epsilon_0} \int \left[\frac{\nabla \rho(\vec{r}', t_r)}{\Delta r} + \rho(\vec{r}', t_r) \nabla \frac{1}{\Delta r} \right] d\tau'$$

ρ does not depend on \vec{r}
 but it does depend on $t_r = t - \frac{\Delta r}{c}$

$$\nabla \rho = \frac{\partial \rho}{\partial t_r} \nabla t_r$$

$$\frac{\partial \rho}{\partial t_r} = \frac{\partial \rho}{\partial t} = \dot{\rho}$$

$$\nabla t_r = -\frac{1}{c} \nabla(\Delta r)$$

$$\Rightarrow \nabla \rho = -\frac{1}{c} \dot{\rho} \nabla(\Delta r)$$

Identities: $\nabla(\Delta r) = \hat{\Delta r}$
 $\nabla\left(\frac{1}{\Delta r}\right) = -\hat{\Delta r}/\Delta r^2$

$$\Rightarrow \nabla V = \frac{1}{4\pi\epsilon_0} \int \left[-\frac{\dot{\rho}}{c} \frac{\hat{\Delta r}}{\Delta r} - \rho \frac{\hat{\Delta r}}{\Delta r^2} \right] d\tau'$$

Next $\nabla^2 V = \frac{1}{4\pi\epsilon_0} \int \left[-\frac{1}{c} \nabla \dot{\rho} \cdot \frac{\hat{\Delta r}}{\Delta r} - \frac{\dot{\rho}}{c} \nabla \cdot \left(\frac{\hat{\Delta r}}{\Delta r} \right) \right. \\ \left. - \nabla \rho \cdot \frac{\hat{\Delta r}}{\Delta r^2} - \rho \nabla \cdot \left(\frac{\hat{\Delta r}}{\Delta r^2} \right) \right] d\tau'$

$$\nabla \dot{\rho} = -\frac{1}{c} \dot{\rho} \nabla(\Delta r) = -\frac{\ddot{\rho}}{c} \hat{\Delta r}$$

Identities: $\nabla \cdot \left(\frac{\hat{\Delta r}}{\Delta r} \right) = \frac{1}{\Delta r^2}$
 $\nabla \cdot \left(\frac{\hat{\Delta r}}{\Delta r^2} \right) = 4\pi \delta^3(\vec{\Delta r})$

$$\Rightarrow \nabla^2 V = \frac{1}{4\pi\epsilon_0} \int \left[\frac{\ddot{\rho}}{c^2} \hat{\Delta r} \cdot \frac{\hat{\Delta r}}{\Delta r} - \frac{\dot{\rho}}{c} \cdot \frac{1}{\Delta r^2} \right. \\ \left. + \frac{\dot{\rho}}{c} \hat{\Delta r} \cdot \frac{\hat{\Delta r}}{\Delta r^2} - \rho \cdot 4\pi \delta^3(\vec{\Delta r}) \right] d\tau'$$

$$= \frac{1}{4\pi\epsilon_0} \int \left[\frac{1}{c^2} \frac{\ddot{\rho}}{\Delta r} - \rho \cdot 4\pi \delta^3(\vec{\Delta r}) \right] d\tau'$$

$$\Rightarrow \nabla^2 V = \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \left(\frac{1}{4\pi\epsilon_0} \int \frac{\rho}{r} d\tau \right) - \rho/\epsilon_0$$

$$\Rightarrow \nabla^2 V = \frac{1}{c^2} \frac{\partial^2 V}{\partial t^2} - \rho/\epsilon_0$$

$$\Rightarrow \square^2 V = -\rho/\epsilon_0 \quad //$$

- So V & \vec{A} just depend on retarded distribution of charge & current

- Beware!!!

The same is not true for \vec{E} and \vec{B}

You can't just put the retarded charge & current distribution into Coulomb's Law, Biot-Savart Law

Jefimenko's Equations

$$\vec{E}(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \int \left[\frac{\rho(\vec{r}', t_r)}{\Delta r^2} \hat{\Delta r} + \frac{\dot{\rho}(\vec{r}', t_r)}{c \Delta r} \hat{\Delta r} - \frac{\ddot{\vec{J}}(\vec{r}', t_r)}{c^2 \Delta r} \right] d\tau'$$

$$\vec{B}(\vec{r}, t) = \frac{\mu_0}{4\pi} \int \left(\left[\frac{\vec{J}(\vec{r}', t_r)}{\Delta r^2} + \frac{\dot{\vec{J}}(\vec{r}', t_r)}{c \Delta r} \right] \times \hat{\Delta r} \right) d\tau'$$

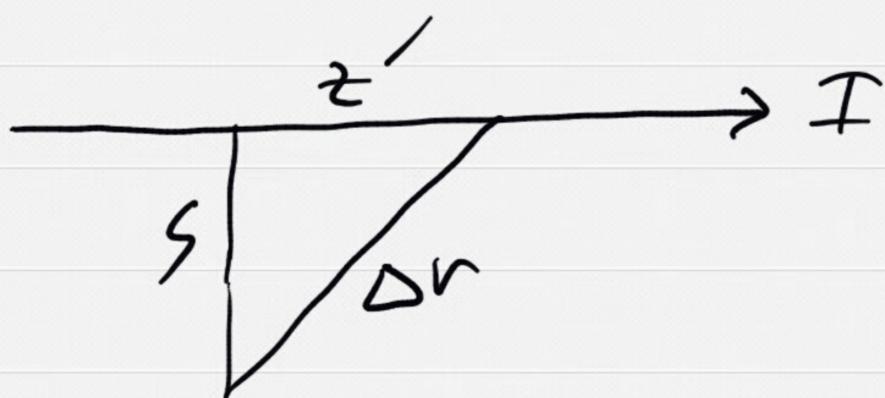
- Leading order terms are Coulomb & Biot-Savart
- These equations are so ugly no one uses them!

Example



Line current
 $I = I_0$, turns
on everywhere
at $t = 0$

$$I(t) = 0 \quad t \leq 0 \\ I_0 \quad t > 0$$



$$\rho = 0 \Rightarrow V = 0$$

$$\vec{A}(\vec{r}, t) = \frac{\mu_0}{4\pi} \int_{-\infty}^{\infty} \frac{I(t_r) \hat{z}}{\Delta r} dz' \\ = \frac{\mu_0}{2\pi} \int_0^{\infty} \frac{I(t_r) \hat{z}}{\Delta r} dz'$$

$$t_r = t - \Delta r/c = t - \frac{1}{c} \sqrt{s^2 + z'^2}$$

$$I(t_r) = 0 \quad t_r \leq 0 \\ I_0 \quad t_r > 0$$

$$t_r > 0 \Rightarrow t > \frac{1}{c} \sqrt{s^2 + z'^2} \\ \Rightarrow z'^2 < (ct)^2 - s^2$$

$$\begin{aligned}
 \text{So } \vec{A}(\vec{r}, t) &= \frac{\mu_0 I_0}{2\pi} \hat{z} \int_0^{\sqrt{(ct)^2 - s^2}} \frac{1}{\sqrt{s^2 + z'^2}} dz' \\
 &= \frac{\mu_0 I_0}{2\pi} \hat{z} \ln(\sqrt{s^2 + z'^2} + z') \Big|_0^{\sqrt{(ct)^2 - s^2}} \\
 &= \frac{\mu_0 I_0}{2\pi} \hat{z} \ln\left(\frac{ct + \sqrt{(ct)^2 - s^2}}{s}\right)
 \end{aligned}$$

$$\begin{aligned}
 \vec{E}(\vec{r}, t) &= -\partial \vec{A} / \partial t \\
 &= -\frac{\mu_0 I_0}{2\pi} \hat{z} \cdot \frac{s \cdot \frac{1}{s} (ct + \frac{1}{s} \sqrt{(ct)^2 - s^2} \cdot 2ct)}{ct + \sqrt{(ct)^2 - s^2}} \\
 &= -\frac{\mu_0 I_0}{2\pi} \hat{z} \frac{c(1 + ct/\sqrt{(ct)^2 - s^2})}{ct + \sqrt{(ct)^2 - s^2}} \\
 &= \boxed{-\frac{\mu_0 I_0 c}{2\pi} \frac{\hat{z}}{\sqrt{(ct)^2 - s^2}}}
 \end{aligned}$$

$$\begin{aligned}
 \vec{B}(\vec{r}, t) &= \nabla \times \vec{A} = -\partial A_z / \partial s \hat{\phi} \\
 &= -\frac{\mu_0 I_0}{2\pi} \hat{\phi} \frac{s}{ct + \sqrt{(ct)^2 - s^2}} \cdot \left(-ct/s^2 + \frac{1}{2\sqrt{(ct/s)^2 - 1}} \cdot \frac{2ct}{s} \cdot -ct/s^2\right) \\
 &= -\frac{\mu_0 I_0}{2\pi} \hat{\phi} \cdot \frac{-ct/s}{ct + \sqrt{(ct)^2 - s^2}} \left(1 + \frac{ct/s}{\sqrt{(ct/s)^2 - 1}}\right) \\
 &= \boxed{\frac{\mu_0 I_0}{2\pi s} \frac{ct \hat{\phi}}{\sqrt{(ct)^2 - s^2}}}
 \end{aligned}$$

$t \rightarrow \infty \Rightarrow \vec{E} \rightarrow 0, \vec{B} \rightarrow \frac{\mu_0 I_0}{2\pi s} \hat{\phi}$
 static solutions