

Lecture #22 Green's Functions

I. Green's Functions

A. Basic Concept

1. Used to solve inhomogeneous differential equations, where the inhomogeneous term acts as a source, and you integrate over the source to obtain solution.

2. Ex: Poisson Equation: $-\nabla^2 \psi(r) = \frac{1}{\epsilon_0} \rho(r)$

a. Solution: $\psi(r) = \frac{1}{4\pi\epsilon_0} \int d^3r' \frac{\rho(r')}{|r-r'|}$

b. Integral contributes over entire region where $\rho(r') \neq 0$.

c. Right-hand side is an integral operator

\Rightarrow Converts charge density $\rho(r)$ to potential $\psi(r)$.

d. Green's Function: $G(r, r') = \frac{1}{4\pi\epsilon_0} \frac{1}{|r-r'|}$

i) This is the kernel that multiplies the inhomogeneous source term.

ii) Function of two variables (r, r'): r is constant, r' is integrated.

e. Therefore $\psi(r) = \int d^3r' G(r, r') \rho(r')$

3. For more general problems, Green's function may depend on

a. Boundary Conditions

b. Form of differential operator.

4. KEY CONCEPT: a. $G(r, r')$ gives contribution to ψ at point r due to a unit-magnitude source (delta function) at r'

b. We can determine ψ at r by summing (integrating) over all points of the source because the differential operator is linear

I. A.4. (Continued)

Hence ②

c. Ex : i. Potential due to charge q at r' $\phi(x) = \frac{1}{4\pi\epsilon_0} \frac{q}{|x-r'|}$

ii. If I have a collection of charges q_j , then

$$\phi(x) = \frac{1}{4\pi\epsilon_0} \sum_j \frac{q_j}{|x-r'_j|}$$

iii. The continuum limit yields $\phi(x) = \frac{1}{4\pi\epsilon_0} \int d^3r' \frac{\rho(r')}{|x-r'|}$
where charge density is $\int d^3r' \rho(r) = Q = \sum_j q_j$
total charge

II. Green's Functions in 1D

A. Fundamental Definition

1. Consider a self-adjoint, inhomogeneous ODE $\mathcal{L}y = \frac{d}{dx} \left(p(x) \frac{dy}{dx} \right) + q(x)y = f(x)$

a. Interval $a \leq x \leq b$

b. Boundary Conditions: Homogeneous at $x=a$ & $x=b$

(Meaning if y_1 satisfies BC's, so does Cy_1)

2. Defining Equation for G :

$$\mathcal{L}G(x, x') = \delta(x-x')$$

3. Solving for $y(x)$:

$$y(x) = \int_a^b G(x, x') f(x') dx'$$

a. Check: $\mathcal{L}y(x) = \int_a^b \mathcal{L}G(x, x') f(x') dx' = \int_a^b \delta(x-x') f(x') dx' = f(x)$
 $= \delta(x-x')$

B. General Properties

Note:

→ 1. Take $\mathcal{L}(x') G(x, x') = \delta(x-x')$ and integrate $\int_{x-\epsilon}^{x+\epsilon} dx'$
 x is constant, but x' is variable that is differentiated.

a. $\int_{x-\epsilon}^{x+\epsilon} f(x') G(x, x') dx' = \int_{x-\epsilon}^{x+\epsilon} \delta(x-x') dx' = 1$

b. $\int_{x-\epsilon}^{x+\epsilon} \frac{d}{dx'} \left[p(x') \frac{dG(x, x')}{dx'} \right] dx' + \int_{x-\epsilon}^{x+\epsilon} q(x') G(x, x') dx' = 1$

II. B. (Continued)

2. First integral is exact $\Rightarrow p(x') \frac{dG(x, x')}{dx'} \Big|_{x=\epsilon}^{x+\epsilon}$ Hence ③

3. Necessary Discontinuity:

$$p(x') \frac{dG(x, x')}{dx'} \Big|_{x=\epsilon}^{x+\epsilon} + \int_{x-\epsilon}^{x+\epsilon} q(x) G(x, x') dx = 1$$

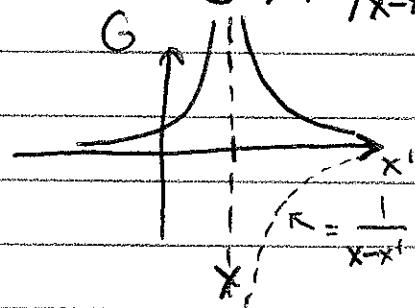
a. In $\lim_{\epsilon \rightarrow 0^+}$, this cannot be satisfied if both $\frac{dG(x, x')}{dx'}$ and $G(x, x')$ are continuous.

b. We require $G(x, x')$ continuous, but allow discontinuity in $\frac{dG(x, x')}{dx'}$

c. In $\lim_{\epsilon \rightarrow 0^+}$, integral vanishes, and we obtain

$$\lim_{\epsilon \rightarrow 0^+} \left[\frac{dG(x, x')}{dx'} \Big|_{x'=x+\epsilon} - \frac{dG(x, x')}{dx'} \Big|_{x'=x-\epsilon} \right] = \frac{1}{p(x)}$$

d. Ex: 1-D $G(x, x') = \frac{1}{|x-x'|}$



$$\frac{dG(x, x')}{dx'} \begin{cases} \frac{1}{(x-x')^2} & x' < x \\ -\frac{1}{(x-x')^2} & x' > x \end{cases} \Rightarrow \text{Discontinuous slope.}$$

4. NOTE: a. The key to understanding Green's functions is to be very careful handling its two variables

b. One variable is always a constant, while the other varies (and is the one over which you integrate).

5. Eigenfunction expansions in two variables:

a. We are going to expand $G(x, x')$ in the orthonormal eigenfunctions of the Hermitian operator \mathcal{L} , ϕ_i

b. NOTE: Both x and x' are separately expanded in the ϕ_i

II. B.5. (Continued)

Thus ④

c. To expand a function $f(x, x')$ in ϕ_i

i. First, expand in $\phi_i(x)$ with $x' = \text{constant}$

$$f(x, x') = \sum_n a_n(x') \phi_n(x), \text{ where } a_n(x') = \langle \phi_n(x) | f(x, x') \rangle$$

Here $\langle \phi_n(x) | f(x, x') \rangle = \int_a^b \phi_n^*(x) f(x, x') dx = a_n(x')$ with $x = \text{const}$.

ii. Second, expand $a_n(x')$ in $\phi_m(x')$ with $x = \text{constant}$

$$a_n(x') = \sum_m b_m \phi_m(x') \quad \text{where } b_m = \langle \phi_m(x') | a_n(x') \rangle$$

iii. Thus, $b_m = \langle \phi_m(x) | a_n(x') \rangle = \int_a^b \phi_m^*(x) a_n(x') dx'$

$$= \int_a^b \phi_m^*(x) \left[\int_a^b \phi_n^*(x) f(x, x') dx \right] dx' \equiv C_{nm}$$

iv. Finally

$$f(x, x') = \sum_{n,m} C_{nm} \phi_n(x) \phi_m(x')$$

6. Now, let's expand $G(x, x')$ in terms of eigenfunctions of \mathcal{L}

a. First, expand $\delta(x-x')$ in terms of $\phi_i(x)$

$$\text{i. } \delta(x-x') = \sum_n c_n(x') \phi_n(x)$$

$$\text{ii. } c_n(x') = \langle \phi_n(x) | \delta(x-x') \rangle = \int_a^b \phi_n^*(x) \delta(x-x') dx = \phi_n^*(x')$$

$$\text{iii. Thus } \delta(x-x') = \sum_n \phi_n(x) \phi_n^*(x')$$

Take $\phi_m^*(x')$, instead of

$$\text{b. Now, } G(x, x') = \sum_{n,m} g_{nm} \phi_n(x) \phi_m^*(x') \leftarrow \phi_m(x'), \text{ to match } \delta(x-x').$$

c. Use defining eq. for G : $\mathcal{L}(x)G(x, x') = \delta(x-x')$ and substitute.

i. NOTE: Here $\mathcal{L}(x)$ is operating on x , not x' !

$$\text{ii. } \mathcal{L}(x) \sum_{n,m} g_{nm} \phi_n(x) \phi_m^*(x') = \sum_m \phi_m(x) \phi_m^*(x')$$

$$\text{iii. } \sum_{n,m} g_{nm} \lambda_n \phi_n(x) \phi_m^*(x') = \sum_m \phi_m(x) \phi_m^*(x')$$

II. B. 6. (Continued)

d. Take scalar product with $\phi_m^*(x')$

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$$\sum_{nm} g_{nm} \lambda_n \phi_n(x) \underbrace{\langle \phi_m^*(x') | \phi_m^*(x') \rangle}_{\delta_{mm'}} = \sum_m \phi_m(x) \underbrace{\langle \phi_m^*(x') | \phi_m^*(x') \rangle}_{\delta_{mm'}} = \delta_{mm'}$$

e. Take scalar product with $\phi_{n'}(x)$

$$\sum_n g_{nm'} \lambda_n \underbrace{\langle \phi_{n'}(x) | \phi_{n'}(x) \rangle}_{\delta_{n'n'}} = \underbrace{\langle \phi_{n'}(x) | \phi_m(x) \rangle}_{\delta_{n'm'}}$$

f. Thus, we obtain

$$g_{n'm'} \lambda_{n'} = \delta_{n'm'} \text{ or } g_{nm} = \frac{\delta_{nm}}{\lambda_n}$$

$$g_{nm} = \frac{\delta_{nm}}{\lambda_n}$$

g. Plugging back into expansion for $G(x, x')$:

$$G(x, x') = \sum_n \frac{\phi_n(x) \phi_n^*(x')}{\lambda_n}$$

7. Symmetry of $G(x, x')$:

$$G(x, x') = G^*(x', x)$$

a. This key property will enable us to establish a procedure for computing Green's Functions.

C. General Solution for $G(x, x')$

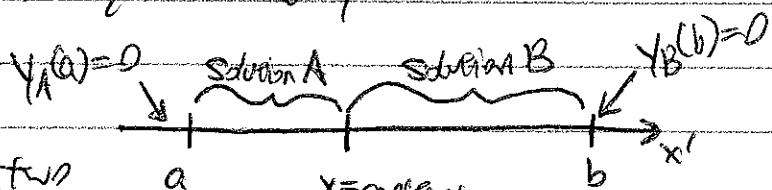
1. Consider a Hermitian operator \mathcal{L} with homogeneous BCs on (a, b) .

$$\mathcal{L}y = \frac{d}{dx} \left[p(x) \frac{dy}{dx} \right] + q(x)y = f(x) \quad \text{Source in inhomogeneous ODE}$$

2. We will determine $G(x, x')$ in terms of solutions $y_1(x), y_2(x)$ of the homogeneous equation $\mathcal{L}y = 0$.

3. Strategy:

- Split range $a \leq x' \leq b$ into two sections, $a \leq x' < x$ and $x < x' \leq b$, where x is constant.
- In each half, satisfy BC's: $y_A(a) = 0, y_B(b) = 0$!



II.C. (Continued)

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4. The general homogeneous solution $y(x) = C_1 y_1(x) + C_2 y_2(x)$

5. Interval $a \leq x' < x$

a. Choose C_1 & C_2 such that $y_A(a) = C_1 y_1(a) + C_2 y_2(a) = 0$

b. The most general form may have a constant $C_A(x) \leftarrow$ can be function of x .

c. Thus $G(x, x') = C_A(x) y_A(x')$, $a \leq x' < x$

6. Following same procedure, $G(x, x') = C_B(x) y_B(x')$, $x < x' \leq b$

7. Thus

$$G(x, x') = \begin{cases} C_A(x) y_A(x') & a \leq x' < x \\ C_B(x) y_B(x') & x < x' \leq b \end{cases}$$

8. Apply Symmetry Condition $G(x, x') = G^*(x', x)$

a. Solution is consistent only if $C_A^*(x) = A y_B(x)$ and $C_B^*(x) = A y_A(x)$
where A is a constant

b. If we assume y_A and y_B can be chosen to be real, we obtain

$$G(x, x') = \begin{cases} A y_A(x') y_B(x) & a \leq x' < x \\ A y_B(x') y_A(x) & x < x' \leq b \end{cases}$$

Solution for $G(x, x')$

c. But, we still need to determine A .

9. Apply $\lim_{\epsilon \rightarrow 0^+} \left[\frac{dG(x, x')}{dx'} \Big|_{\substack{x \leq x+\epsilon \\ x' > x}} - \frac{dG(x, x')}{dx'} \Big|_{\substack{x'+\epsilon = x \\ x' < x}} \right] = \frac{1}{p(x)}$

a. $A [y_B'(x) y_A(x) - y_A'(x) y_B(x)] = \frac{1}{p(x)}$

b. Thus $A = \left\{ p(x) [y_B'(x) y_A(x) - y_A'(x) y_B(x)] \right\}^{-1}$

Solution for A

c. Note! A is actually a constant \rightarrow does not depend on x !

ii. Term in brackets is Wronskian, $[\dots] = W = \frac{F}{p(x)} \leftarrow$ can be shown.

II. (Continued)

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D. Example 1:

1. Consider $-y'' = f(x)$, where $y(0) = y(1) = 0$,

2. General Solution to Homogeneous equation: $y'' = 0$

$$y = C_0 + C_1 x$$

3. Determine Solutions y_0 and y_1 that satisfy B.C.'s:

$$a. y(0) = 0 = C_0 + C_1(0) \Rightarrow C_0 = 0 \Rightarrow y_0(x) = x$$

$$b. y(1) = 0 = C_0 + C_1(1) \Rightarrow C_1 = -C_0 \Rightarrow y_1(x) = 1-x$$

4. Compute coefficients A: $p(x) = -1$, $y_0'(x) = 1$, $y_1'(x) = -1$

$$A = \left\{ p(x) [y_1'(x)y_0(x) - y_0'(x)y_1(x)] \right\}^{-1} = \left\{ (-1) [(-1)(x) - (1)(1-x)] \right\}^{-1} = \frac{1}{x-1+x}$$

$$5. \text{Thus } G(x, x') = \begin{cases} x'(1-x), & 0 \leq x' \leq x \\ (1-x')x, & x \leq x' \leq 1 \end{cases}$$

6. Thus, Solution $y(x)$ is $y(x) = \int_0^1 G(x, x') f(x') dx'$ for any $f(x')$.

7. a. For example, $f(x) = \sin(\pi x)$ yields

$$b. y(x) = \int_0^1 G(x, x') \sin(\pi x') dx' = \int_0^x x'(1-x) \sin(\pi x') dx' + \int_x^1 (1-x) x \sin(\pi x') dx'$$

$$c. \text{Using } \int_a^b x \sin(\pi x') dx' = \frac{\sin(\pi x)}{\pi^2} - \frac{x \cos(\pi x)}{\pi} \Big|_a^b, \text{ we obtain}$$

$$(1-x) \left[\frac{\sin(\pi x)}{\pi^2} - \frac{x \cos(\pi x)}{\pi} \right]_0^x + x \left[\frac{-\cos(\pi x)}{\pi} \right]_x^1 - x \left[\frac{\sin(\pi x)}{\pi^2} - \frac{x \cos(\pi x)}{\pi} \right]_x^1$$

$$= (1-x) \left[\frac{\sin(\pi x)}{\pi^2} - \frac{x \cos(\pi x)}{\pi} \right] + x \left[\frac{-\cos(\pi x)}{\pi} \right] - x \left[\frac{-\sin(\pi x)}{\pi^2} - \frac{x \cos(\pi x)}{\pi} \right]$$

$$= \frac{\sin(\pi x)}{\pi^2} - \frac{x \cos(\pi x)}{\pi} - x \left(\frac{\sin(\pi x)}{\pi^2} - \frac{x \cos(\pi x)}{\pi} \right) + \frac{x \sin(\pi x)}{\pi} + x \left(\frac{\sin(\pi x)}{\pi^2} - \frac{x \cos(\pi x)}{\pi} \right)$$

$$y(x) = \frac{\sin(\pi x)}{\pi^2}$$

8. NOTE! With Green's functions, if you change $f(x)$, you need only repeat the integration of the Green's function kernel with $f(x)$.

II. (Continued)

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E. Other Forms of Boundary Conditions

1. Non-homogeneous boundary conditions: $y(a) = c_1$, $y(b) = c_2$

a. Simply change y to $U = y - \frac{c_1(b-x) + c_2(x-a)}{b-a}$

b. In terms of U , $U(a) = 0$, $U(b) = 0$.

2. Initial Value Problems: $L[y(t)] = f(t)$, $f(0) = 0$, $f'(0) = 0$.

a. Construct $G(t, t')$ by invoking continuity of G at $t=t'$.

b. BUT, $G(t, t')$ is no longer symmetric.

3. Ext Initial Value Problem: $L[y] = \frac{d^2y}{dt^2} + y = f(t)$, $y(0) = 0$, $y'(0) = 0$

a. Homogeneous Solution: $L[y] = 0 \Rightarrow y_1 = \sin t$, $y_2 = \cos t$

b. Key point! No linear combination $y_0 = C_1 y_1 + C_2 y_2$ satisfies BC's at $t=0$, except for $C_1 = C_2 = 0$.

i. Thus, for $t' < t$, we have $G(t, t') = 0$.

c. Second key point: At $t' > t$, we have no constraining BC's, so

i. For $t' > t$, $G(t, t') = C_1(t) \sin t' + C_2(t) \cos t'$

d. Thus $G(t, t') = \begin{cases} 0 & t' < t \\ C_1(t) \sin t' + C_2(t) \cos t' & t' > t \end{cases}$

e. Enforce continuity requirements:

$$\lim_{\epsilon \rightarrow 0^+} [G(t, t+\epsilon) - G(t, t-\epsilon)] = 0 \Rightarrow [C_1(t) \sin t + C_2(t) \cos t - 0] = 0$$

$$\text{ii. } \lim_{\epsilon \rightarrow 0^+} \left[\left| \frac{dG(t, t)}{dt} \right|_{t=\pm\epsilon} - \left| \frac{dG(t, t)}{dt} \right|_{t=t-\epsilon} \right] = \frac{1}{P(t)} \quad \text{where } P(t) = 1 \text{ for } L.$$

$$C_1(t) \cos t + C_2(t)(-\sin t) - 0 = 1$$

IV. Solving these equations yields $C_1(t) = \cos t$, $C_2(t) = \sin t$.

$$\text{v. Thus } \cos t \sin t' - \sin t \cos t' = \sin(t-t') \Rightarrow G(t, t') = \begin{cases} 0 & t' < t \\ \sin(t-t') & t' > t \end{cases}$$