

Electricity and Magnetism II: 3812

Professor Jasper Halekas Van Allen 70 MWF 9:30-10:20 Lecture

Xample i Capaci /L Ē = -1/2, È between plates bettem plate Foree $F = \oint F \cdot da$ I Et Area F2 = g(T2xdax + Tzyda, + Tzzdaz) = § Tzzlaz (dax = day = 0 for surface between plates) $\begin{aligned} 1 \neq 2 &= \sum \left(E_{2}E_{2} - E_{2} \int 2 d z z E^{2} \right) \\ &= \sum \left(E_{2}^{2} - E_{2}^{2} - E_{2}^{2} \right) \end{aligned}$ = 5/2 (%)2 0/22. 6122 292 Qut. Etas

Example; Moving Capacitor $d_{2} = V \hat{\mathcal{G}}_{L}$ $\vec{E} = -\sigma_{\vec{k},\sigma} \hat{\vec{x}}$ $\vec{B} = -\eta_{\vec{k},\sigma} \sigma_{\vec{k}} \hat{\vec{x}}$ Between plates;

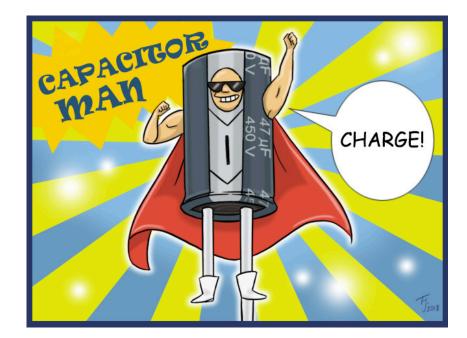
 $\Rightarrow \vec{g}_{EM} = p_0 \varepsilon, \vec{s}$ $= \varepsilon \cdot (\vec{E} \times \vec{s})$ $= p_0 \varepsilon^2 v \hat{\varphi}$ $\tilde{P}_{EM} = \tilde{q}_{EM} \cdot v \cdot lume$ = $\mu \cdot \sigma^2 v d L^2 \hat{q}$ Now move top plate down $w/ u_2 = -u_2^2$

 $F_{y6} = Q_{top}(\hat{u_t} \times \bar{b_x}) \quad \text{on top plate}$ = $\sigma L^2 \cdot (-u\hat{t} \times -\frac{1}{2}p_0 \sigma \sqrt{x})$ (average)= $\frac{1}{2}p_0 \sigma^2 L^2 u \sqrt{y}$ $G \quad top plate$ $\Delta P_{yb} = S F_{yb} \Delta t = I_2 g_0 \sigma^2 L^2 v S u d t$ $= I_2 g_0 \sigma^2 L^2 v d$

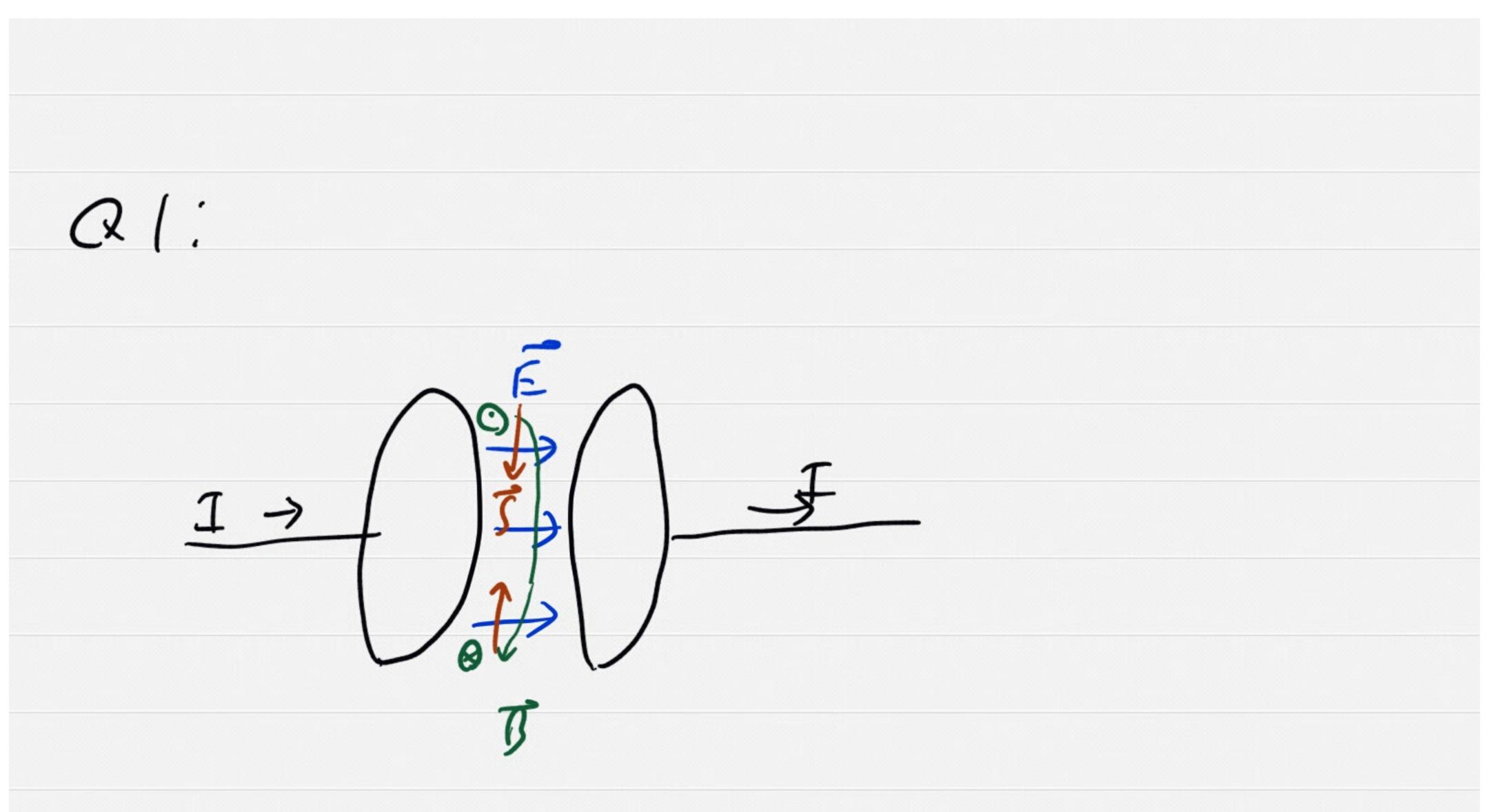
É frim 25/37 as to plate -Also maves Babove = 0 Bbelow = -m. ovx SE-le = 1/dt SB-da (Eabove - Ebelow). L = 1/dt (Bbelow L.d) = Brelow · L · - U => Eabove - Ebelow $= \beta_{below} - \mu$ $= p_{0} \sigma V \mu$ => Eabore = 12 por vug Ebelow = - 12 por vug Fye = abottom - Etelow on bottom plate = - ~ L2. - カア. ~ Vuý = 12 poor L2 Vuý => DPye = tr partivelý => Dpy = mor2 L2 Vdý = PyEm (0) // Capacitor gains momentum from EM fields!

Check Your Understanding

- Which direction do the electric and magnetic fields point between the plates of a charging capacitor?
- Which direction is the Poynting vector?



 $\vec{S} = \frac{1}{\mu_0} \left(\vec{E} \times \vec{B} \right)$



 $\overline{E} = 9/\epsilon. \hat{z}$ $\overline{B} = \frac{p_0 Fr}{2TR^2} \hat{q}$ $\vec{s} = \frac{1}{r} (\vec{E} \times \vec{B})$ $= \int \frac{\sigma I r}{2\pi \epsilon R^2} \hat{r} \Big|$

Check Your Understanding

 Use the Maxwell stress tensor to show that there is no net force on a finite portion of an infinite sheet of charge with area A.

$$E_{\perp} = E$$

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$$Gaussian$$

$$Surface$$

$$F = \frac{d\vec{p}_{mech}}{dt} = \oint \vec{T} \cdot \vec{da} - \frac{d\vec{p}_{EM}}{dt}$$

 T_{ij}

 $Q_2: \vec{E} = \mathscr{D}_{\mathcal{L}}. \hat{\eta} \quad \mathcal{U} = \pm \frac{1}{2}$ $T = \varepsilon \begin{pmatrix} -\lambda E^2 & 0 & 0 \\ 0 & -\lambda E^2 & 0 \\ 0 & 0 & \lambda E^2 \end{pmatrix}$ $= \frac{0}{8\epsilon_{0}} \begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}$

 $\begin{cases} \stackrel{d}{=} \cdot d \vec{a} = \int T_{22} d a_2 = \frac{\sigma^2}{\delta \epsilon}.$ Sittin 7.1a = STat. - daz = - 5/82. Ssides F.da = STxx.dax x + STyp.day g = 0

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 $\Rightarrow \int f da = 0$