

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

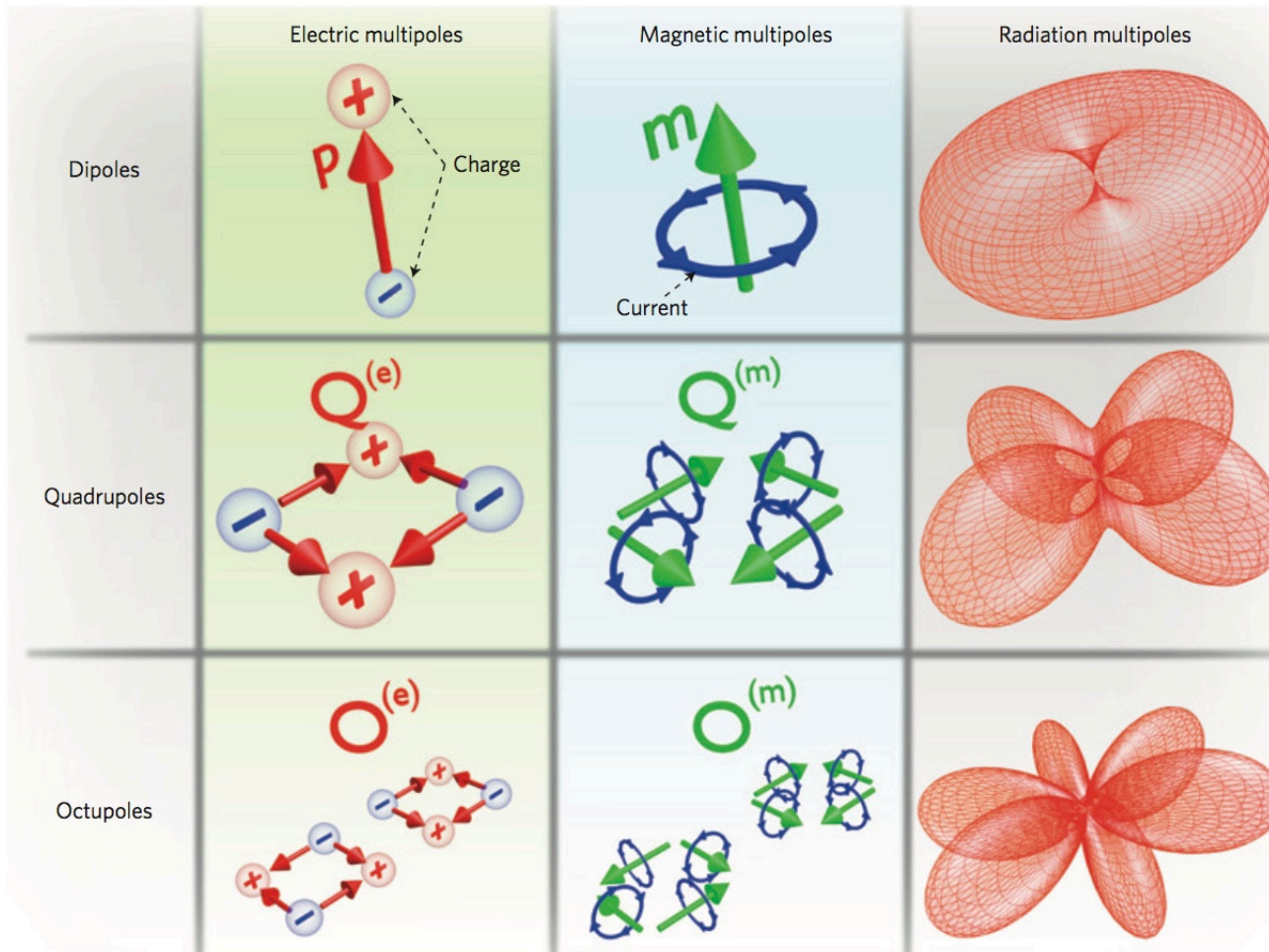
$$\nabla \times \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} + \mu_0 \mathbf{J}$$



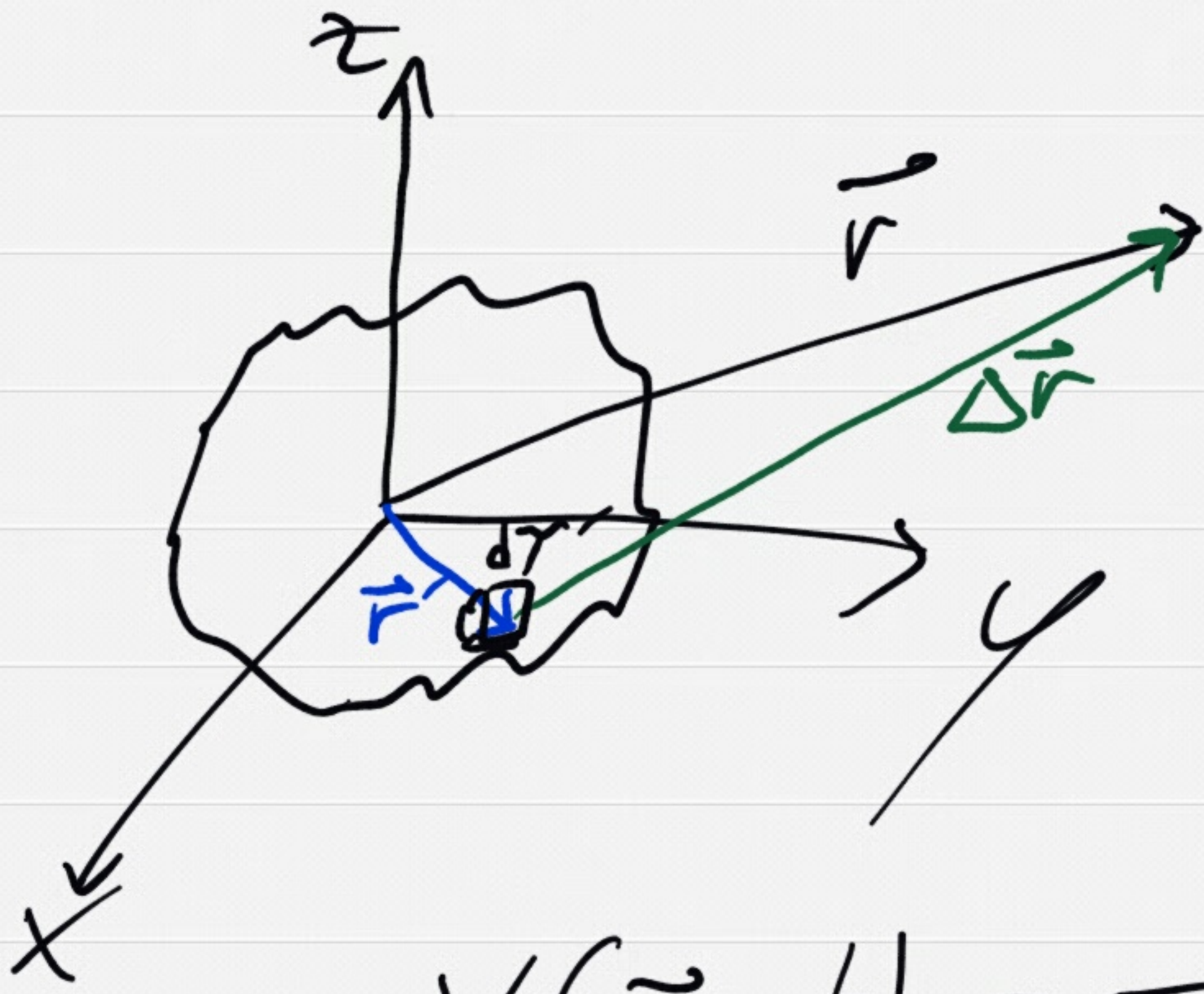
Electricity and Magnetism II: 3812

Professor Jasper Halekas
Virtual by Zoom!
MWF 9:30-10:20 Lecture

Multipole Radiation



Multipole Radiation



$$V(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}', t_r)}{\Delta r} d\gamma'$$
$$= \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}', t - \Delta r/c)}{\Delta r} d\gamma'$$

$$\Delta r = \sqrt{r^2 + r'^2 - 2\vec{r} \cdot \vec{r}'}$$

Assume $r' \ll r$ (far from source)

$$\Rightarrow \Delta r \sim r \left(1 - \frac{\vec{r} \cdot \vec{r}'}{r^2}\right) = r \left(1 - \frac{\hat{r} \cdot \vec{r}'}{r}\right)$$

$$1/\Delta r \sim \frac{1}{r} \left(1 + \frac{\vec{r}' \cdot \vec{r}}{r^2}\right) = \frac{1}{r} \left(1 + \frac{\hat{r} \cdot \vec{r}'}{r}\right)$$

$$\rho(\vec{r}', t - \Delta r/c) \sim \rho(\vec{r}', t - \frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c})$$

Expand ρ about $t_0 = t - r/c$
(retarded time at origin)

$$\Rightarrow \rho(\vec{r}', t - \frac{\Delta r}{c}) \sim \rho(\vec{r}', t_0) + \dot{\rho}(\vec{r}', t_0) \cdot \frac{\hat{r} \cdot \vec{r}'}{c} + \dots$$

$$V(\vec{r}, t) \sim \frac{1}{4\pi\epsilon_0} \int \left\{ \frac{1}{r} \left(1 + \frac{\hat{r} \cdot \vec{r}'}{r} \right) \left[\rho(\vec{r}', t_0) + \dot{j}(\vec{r}', t_0) \frac{\hat{r} \cdot \vec{r}'}{c} \right] \right\} d\tau'$$

Keep only first order terms in r'

$$\Rightarrow V(\vec{r}, t) \sim \frac{1}{4\pi\epsilon_0 r} \left[\int \rho(\vec{r}', t_0) d\tau' + \frac{\hat{r}}{r} \cdot \int \vec{r}' \rho(\vec{r}', t_0) d\tau' + \frac{\hat{r}}{c} \cdot \int \vec{r}' \dot{j}(\vec{r}', t_0) d\tau' \right]$$

$$V(\vec{r}, t) \approx \frac{1}{4\pi\epsilon_0} \left[\frac{Q}{r} + \frac{\hat{r} \cdot \vec{p}(t_0)}{r^2} + \frac{\hat{r} \cdot \dot{\vec{p}}(t_0)}{rc} \right]$$

\uparrow Monopole \uparrow Dipole \uparrow radiation

$$\vec{A}(\vec{r}, t) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}', t - \frac{\Delta r}{c})}{\Delta r} d\tau'$$

$$\sim \frac{\mu_0}{4\pi r} \int \vec{J}(\vec{r}', t_0) d\tau'$$

Turns out $\int \vec{J} d\tau' = \dot{\vec{p}}$

$$\Rightarrow \vec{A}(\vec{r}, t) \approx \frac{\mu_0}{4\pi r} \dot{\vec{p}}(t_0)$$

Radiation Fields

Dropping all $1/r^2$ terms ---

$$\nabla V \cong \nabla \left(\frac{1}{4\pi\epsilon_0} \frac{\hat{r} \cdot \dot{p}(t_0)}{rc} \right)$$

$$\cong \frac{1}{4\pi\epsilon_0} \frac{\hat{r} \cdot \ddot{p}(t_0)}{rc} \nabla t_0$$

$$\text{but } \nabla t_0 = -\frac{1}{c} \nabla r = -\frac{\hat{r}}{c}$$

$$\Rightarrow \nabla V \cong -\frac{1}{4\pi\epsilon_0 c^2} \frac{\hat{r} \cdot \ddot{p}(t_0)}{r} \hat{r}$$

$$\nabla \times \vec{A} \cong \frac{\mu_0}{4\pi r} \nabla \times \dot{\vec{p}}(t_0)$$

$$\cong \frac{\mu_0}{4\pi r} \nabla t_0 \times \ddot{\vec{p}}(t_0)$$

$$= -\frac{\mu_0}{4\pi r c} \hat{r} \times \ddot{\vec{p}}(t_0)$$

$$\frac{\partial \vec{A}}{\partial t} \cong \frac{\mu_0}{4\pi} \frac{\ddot{\vec{p}}(t_0)}{r}$$

$$\Rightarrow \vec{E}(\vec{r}, t) \cong \frac{\mu_0}{4\pi r} [(\hat{r} \cdot \ddot{\vec{p}}) \hat{r} - \ddot{\vec{p}}]$$

$$\boxed{\vec{E}(\vec{r}, t) \cong \frac{\mu_0}{4\pi r} [\hat{r} \times (\hat{r} \times \ddot{\vec{p}})]}$$

$$\boxed{\vec{B}(\vec{r}, t) \cong \frac{\mu_0}{4\pi r c} \hat{r} \times \ddot{\vec{p}}}$$

w/ $\ddot{\vec{p}}$ evaluated
at $t_0 = t - r/c$

Spherical Coordinates

$$\vec{E}(r, \theta, t) \approx \frac{\mu_0 \ddot{p}_0(t)}{4\pi} \frac{\sin\theta}{r} \hat{\theta}$$

$$\vec{B}(r, \theta, t) \approx \frac{\mu_0 \dot{\ddot{p}}_0(t)}{4\pi c} \frac{\sin\theta}{r} \hat{\phi}$$

w/ z-axis along \ddot{p}
(not along \vec{r} !)

$$\vec{S}(r, \theta, t) = \frac{1}{\mu_0} (\vec{E} \times \vec{B})$$

$$\vec{S}(r, \theta, t) = \frac{\mu_0}{16\pi^2 c} [\ddot{p}(t)]^2 \frac{\sin^2\theta}{r^2} \hat{r}$$

$$P(r, t) = \oint \vec{S} \cdot d\vec{a}$$

$$P(r, t) = \frac{\mu_0}{6\pi c} [\ddot{p}(t - r/c)]^2$$

Example 1: Oscillating Dipole

$$p(t) = p_0 \cos(\omega t)$$

$$\ddot{p}(t) = -p_0 \omega^2 \cos(\omega t)$$

$$\Rightarrow P_{\text{rad}} = \frac{\mu_0 \omega^4 p_0^2}{6\pi c} \cos^2(\omega(t - r/c))$$

$$\text{and } \langle P_{\text{rad}} \rangle = \frac{\mu_0 \omega^4 p_0^2}{12\pi c} \quad \text{same as previous derivation}$$

Example 2: Accelerated Particle

Point charge at position $\vec{d}(t)$

$$\vec{p}(t) = q \vec{d}(t)$$

$$\begin{aligned} \ddot{\vec{p}}(t) &= q \frac{d^2 \vec{d}(t)}{dt^2} \\ &= q \vec{a}(t) \end{aligned}$$

$$\Rightarrow \boxed{P_{\text{rad}} = \frac{\mu_0 q^2 a^2}{6\pi c}} \quad \begin{array}{l} \checkmark a \\ \text{at retarded} \\ \text{time} \end{array}$$

Larmor Formula!

Check Your Understanding

- Verify Faraday's law $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$ for electric dipole radiation:

$$\vec{E} = -\hat{\theta} \frac{p_0}{4\pi\epsilon_0 r} \frac{\omega^2}{c^2} \sin\theta \cos\omega \left(t - \frac{r}{c}\right)$$

$$\vec{B} = -\frac{\mu_0 p_0 \omega^2}{4\pi r c} \hat{\phi} \sin\theta \cos\omega \left(t - \frac{r}{c}\right)$$

$$\nabla \times \vec{A} = \frac{1}{r \sin\theta} \left(\frac{\partial(\sin\theta A_\phi)}{\partial\theta} - \frac{\partial A_\theta}{\partial\phi} \right) \hat{r} + \frac{1}{r} \left(\frac{1}{\sin\theta} \frac{\partial A_r}{\partial\phi} - \frac{\partial(r A_\phi)}{\partial r} \right) \hat{\theta} + \frac{1}{r} \left(\frac{\partial(r A_\theta)}{\partial r} - \frac{\partial A_r}{\partial\theta} \right) \hat{\phi}$$

Check Your Understanding

- Verify Ampere's law $\nabla \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$ for electric dipole radiation:

$$\vec{E} = -\hat{\theta} \frac{p_0}{4\pi\epsilon_0 r} \frac{\omega^2}{c^2} \sin\theta \cos\omega \left(t - \frac{r}{c}\right)$$

$$\vec{B} = -\frac{\mu_0 p_0 \omega^2}{4\pi r c} \hat{\phi} \sin\theta \cos\omega \left(t - \frac{r}{c}\right)$$

$$\nabla \times \vec{A} = \frac{1}{r \sin\theta} \left(\frac{\partial(\sin\theta A_\phi)}{\partial\theta} - \frac{\partial A_\theta}{\partial\phi} \right) \hat{r} + \frac{1}{r} \left(\frac{1}{\sin\theta} \frac{\partial A_r}{\partial\phi} - \frac{\partial(r A_\phi)}{\partial r} \right) \hat{\theta} + \frac{1}{r} \left(\frac{\partial(r A_\theta)}{\partial r} - \frac{\partial A_r}{\partial\theta} \right) \hat{\phi}$$

$$\begin{aligned}
 Q1: \nabla \times \vec{E} &= \frac{1}{r} \frac{\partial}{\partial r} (r E_{\theta}) \hat{\phi} \\
 &= \frac{1}{r} \frac{\partial}{\partial r} \left(\frac{-\rho_0 \omega^2}{4\pi \epsilon_0 c^2} \sin \theta \cos(\omega(t - r/c)) \right) \hat{\phi} \\
 &= \frac{-\rho_0 \omega^3}{4\pi \epsilon_0 c^3 r} \sin \theta \sin(\omega(t - r/c)) \hat{\phi}
 \end{aligned}$$

$$\frac{\partial \vec{B}}{\partial t} = \frac{\mu_0 \rho_0 \omega^3}{4\pi r c} \sin \theta \sin(\omega(t - r/c)) \hat{\phi}$$

$$\frac{1}{\epsilon_0 c^2} = \mu_0 //$$

$$\begin{aligned}
 Q2: \nabla \times \vec{B} &= \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta B_{\phi}) \hat{r} \\
 &\quad - \frac{1}{r} \frac{\partial}{\partial r} (r B_{\phi}) \hat{\theta}
 \end{aligned}$$

- The $\hat{\theta}$ term will match up w/ $\mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$

- But, there's an extra radial part, which would only be compensated by one of the terms we dropped in our derivation

- This term didn't affect \vec{S} , as it turns out