

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

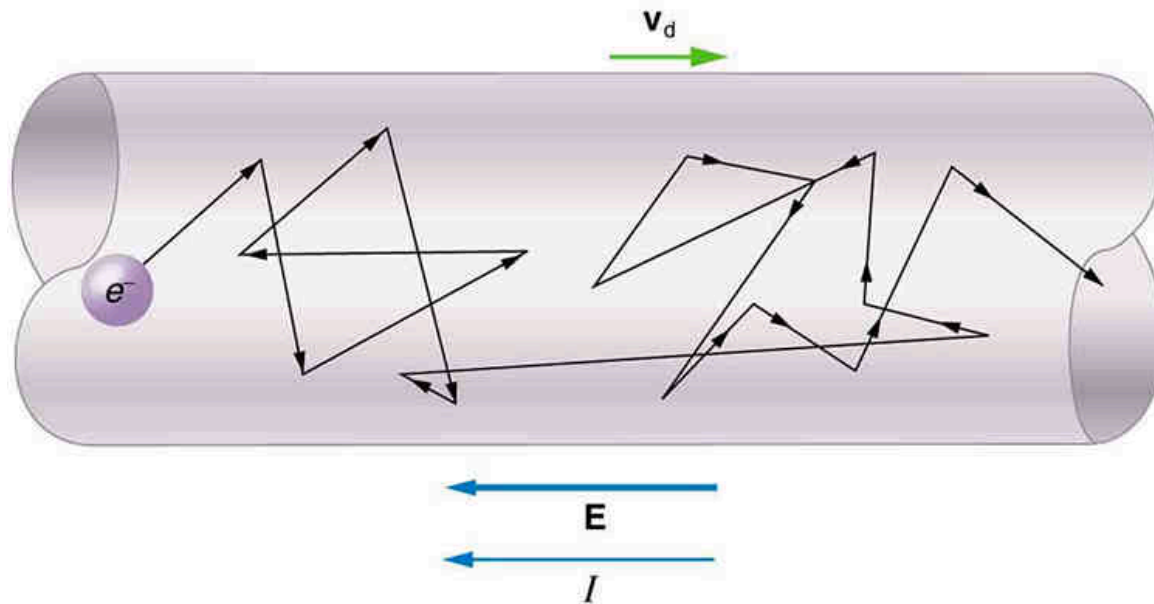
$$\nabla \times \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} + \mu_0 \mathbf{J}$$



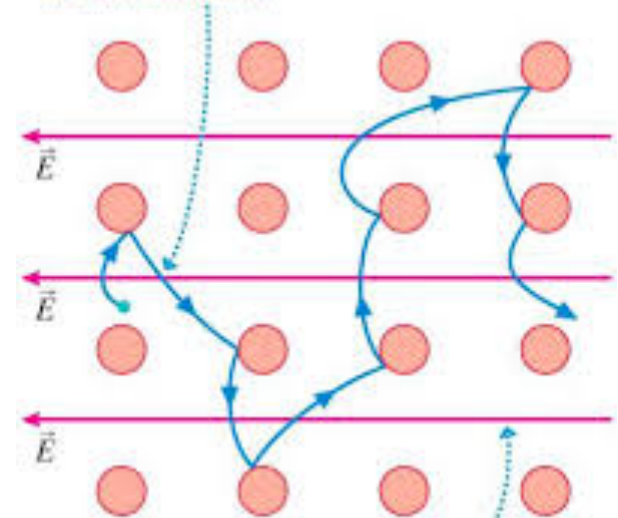
# Electricity and Magnetism II: 3812

Professor Jasper Halekas  
Van Allen 70  
MWF 9:30-10:20 Lecture

# Ohm's Law



The collisions "reset" the motion of the electron. It then accelerates until the next collision.



The electron has a net displacement opposite the electric field.

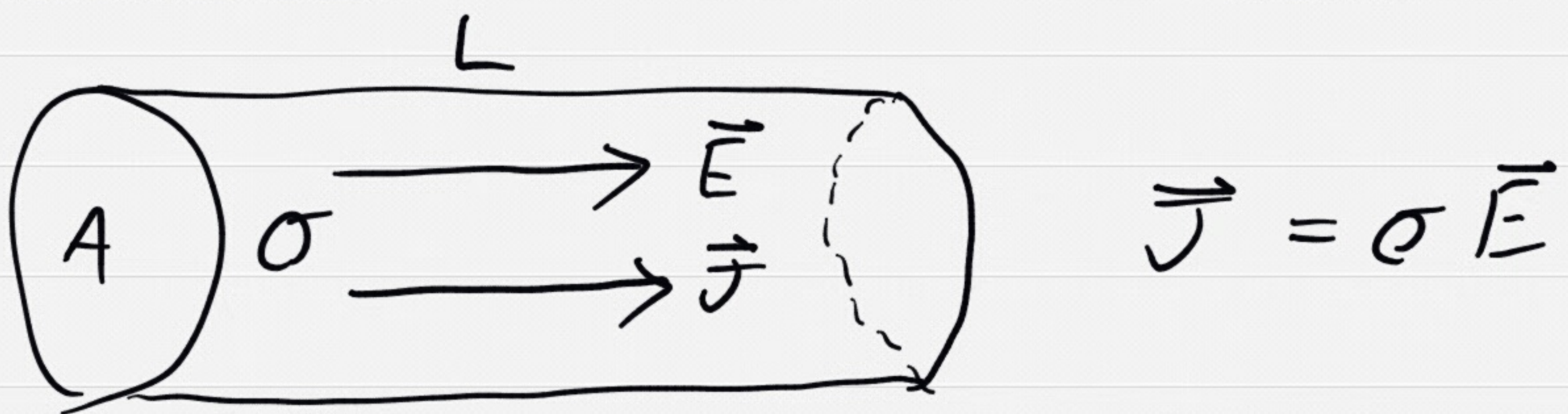


- Ohm's law implies dissipation of energy

$$\vec{J} = \sigma \vec{E}$$

$$\begin{aligned} P &= \frac{dW}{dt} = \frac{d}{dt} \int \vec{F} \cdot d\vec{l} \\ &= \vec{F} \cdot \vec{v} \\ &= Q \cdot \vec{E} \cdot \vec{v} \\ &= \rho \cdot \text{Volume} \cdot \vec{E} \cdot \vec{v} \\ &= (\vec{J} \cdot \vec{E}) \cdot \text{Volume} \\ &\quad \text{if } \vec{J} = \rho \vec{v} \end{aligned}$$

Microscopic  $\rightarrow$  Macroscopic



$$I = J \cdot A, \quad \Delta V = - \int \vec{E} \cdot d\vec{l} = -EL$$

$$\Rightarrow \frac{I}{A} = -\sigma \cdot \frac{\Delta V}{L}$$

$$\Rightarrow \Delta V = -\frac{L}{\sigma A} \cdot I = -IR$$

$$\text{w/ } R = \frac{L}{\sigma A}$$



$$P = (\vec{J} \cdot \vec{E}) \cdot \text{Volume}$$

$$= \frac{I}{A} \cdot \frac{\Delta V}{L} \cdot LA$$

$$= I \cdot \Delta V \quad \text{usually shortened to } P = VI$$

$$P = VI = I^2 R = \frac{V^2}{R}$$

(since  $V = IR$ )

- This power dissipation is called "Joule heating" and is due to collisions

- In a perfect conductor

$$\sigma \rightarrow \infty$$

$$\Rightarrow R, \vec{E}, \Delta V \rightarrow 0$$

$$\Rightarrow P \rightarrow 0$$



# EMF

$$\mathcal{E} = \oint \vec{F} \cdot d\vec{l}$$

= "Electromotive Force"

- Loop integral of force per charge

- Many sources of EMF

- Mechanical (generator)

- Chemical (battery)

- Electric field

- Electrostatic EMF  
is zero

$$\oint \vec{E} \cdot d\vec{l} = \int (\nabla \times \vec{E}) \cdot d\vec{a}$$

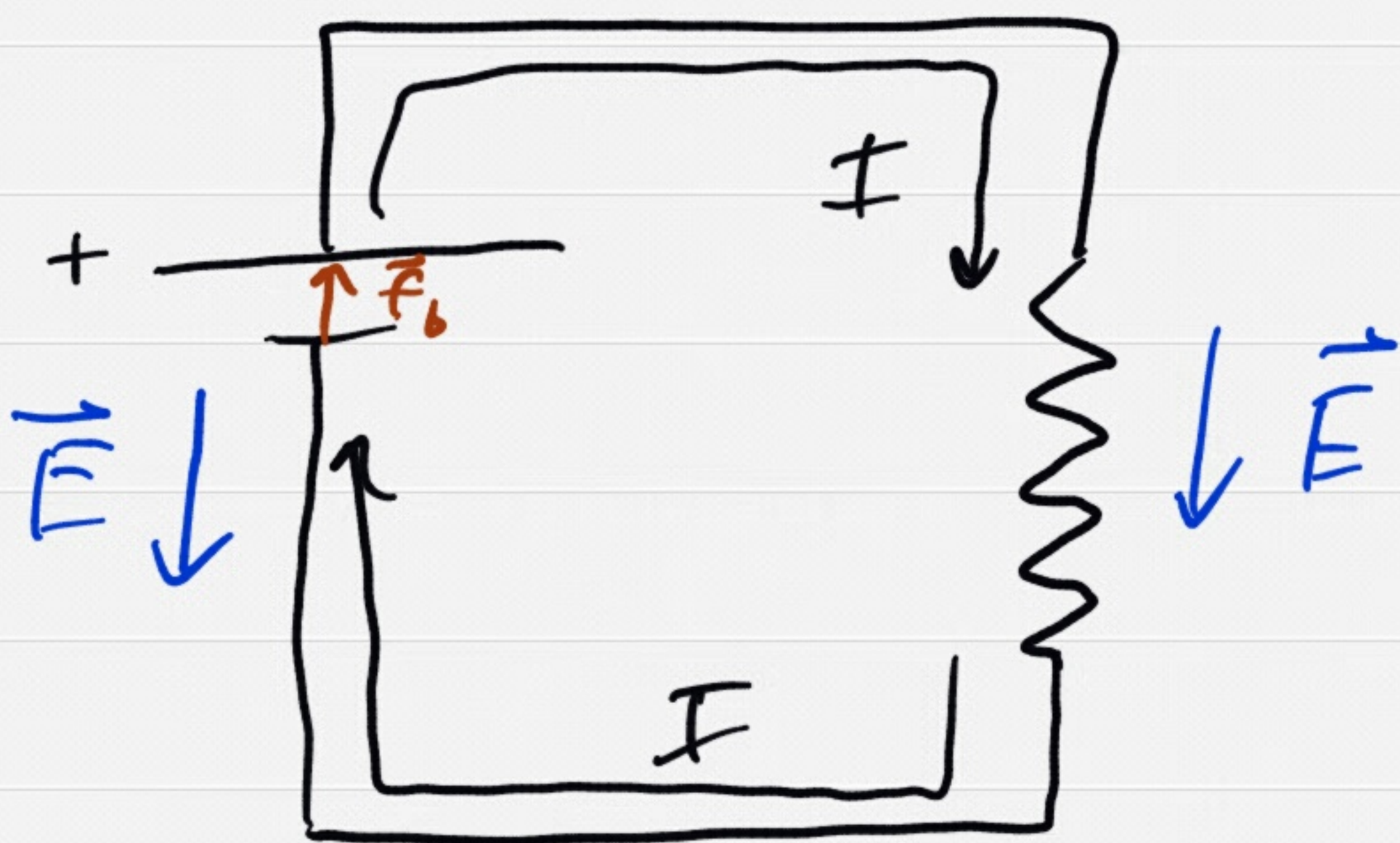
zero if  $\nabla \times \vec{E} = 0$

⇒ Kirchhoff's Law:

Sum of voltage drops  
around circuit equal  
to zero



# Simple Circuit



$$\Delta V_{\text{battery}} = +V_0$$

$$\Delta V_{\text{resistor}} = -V_0$$

$$\Delta V_{\text{total}} = V_0 - V_0 = 0$$

$$= -\oint \vec{E} \cdot d\vec{\ell}$$

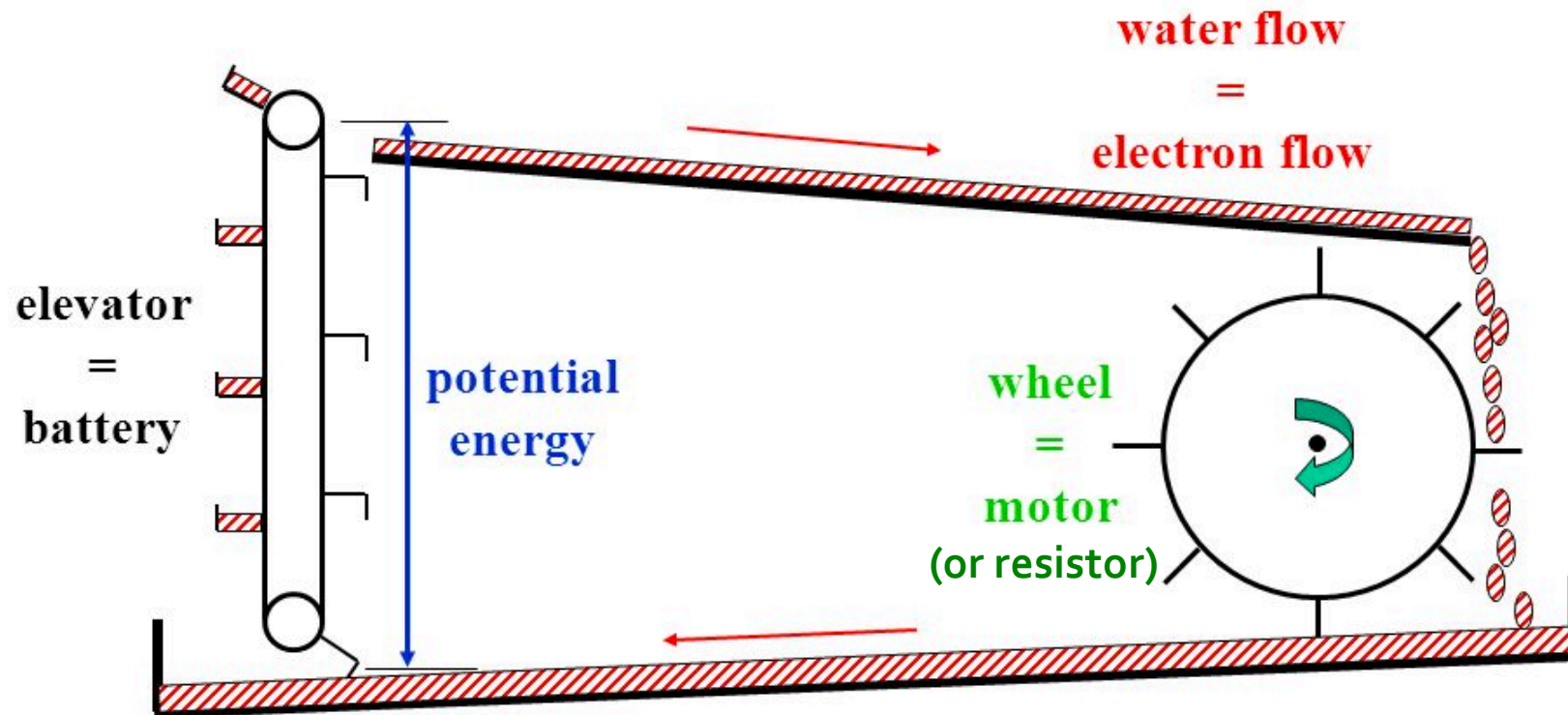
$$\mathcal{E} = \oint \vec{F} \cdot d\vec{\ell} = \oint \vec{f}_0 \cdot d\vec{\ell}$$

- In battery  $\vec{f}_0 = -\vec{E}$
- $\vec{f}_{\text{total}} = \vec{f}_0 + \vec{E} = 0$  in battery
- ohm's law not satisfied in battery
- $\mathcal{E} = V_0$

- $\vec{f}_0$  allows charges to overcome  $\vec{E}$  in battery
- then they flow in response to  $\vec{E}$  through circuit

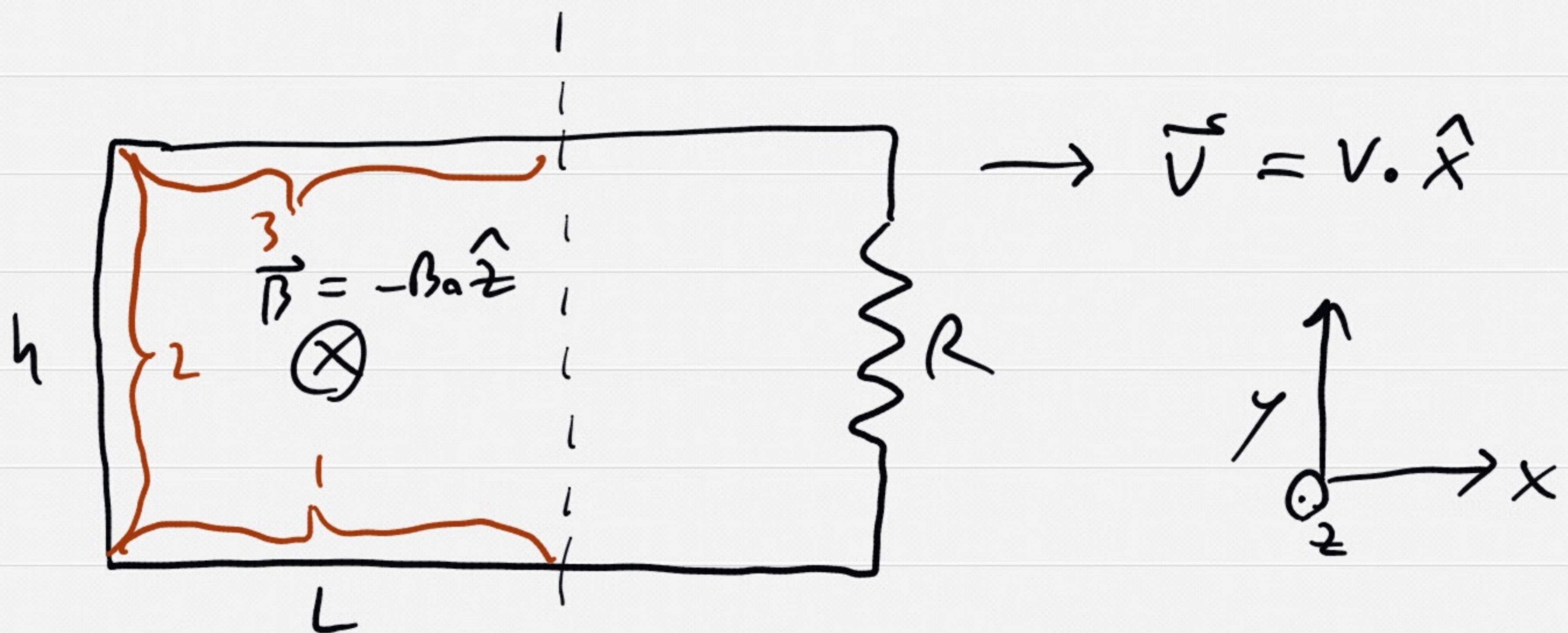


# Electric Circuit Analogy





# Motional EMF



$$\begin{aligned}\vec{F}_1 &= q_1 \vec{v}_1 \times \vec{B} \\ &= q_1 v \cdot B_0 \hat{y}\end{aligned}$$

$$\vec{F}_2 = q_2 v \cdot B_0 \hat{y}$$

$$\vec{F}_3 = q_3 v \cdot B_0 \hat{y}$$

$$\vec{f} = \vec{F}/q = v \cdot B_0 \hat{y} \text{ for all three}$$

$$\mathcal{E} = \oint \vec{f} \cdot d\vec{\ell} \quad \text{w/ } d\vec{\ell} \text{ clockwise}$$

$$= v \cdot B_0 h \quad (\text{only from segment 2})$$

Drives current  $I = \frac{\mathcal{E}}{R}$

$$= \frac{v \cdot B_0 h}{R}$$



Who does work?

w/ current flowing  
 $I = \mathcal{E}/R = q u$

$$\begin{aligned}\vec{F}_{1u} &= q_1 u \cdot -\hat{x} \times \vec{B} \\ &= -q_1 u B_0 \hat{y} \\ \vec{F}_{2u} &= q_2 u \hat{y} \times \vec{B} \\ &= -q_2 u B_0 \hat{x} \\ \vec{F}_{3u} &= q_3 u \hat{x} \times \vec{B} \\ &= q_3 u B_0 \hat{y}\end{aligned}$$

$$q_1 = q_3 \rightarrow \vec{F}_u = -q_2 u B_0 \hat{x}$$

$$\vec{F}_{\text{pull}} = -\vec{F}_u = q_2 u B_0 \hat{x}$$

$$\vec{F}_{\text{pull}} = \vec{F}_{\text{pull}}/q_2 = u B_0 \hat{x}$$

$$\begin{aligned}w/q &= \int \vec{F}_{\text{pull}} \cdot d\vec{x} = u B_0 \cdot \Delta x \\ &= u B_0 \cdot v_0 \Delta t\end{aligned}$$

$\Delta t$  is time for charge to traverse segment 2

$$\Rightarrow w/q = B_0 v_0 \cdot l = \mathcal{E}$$